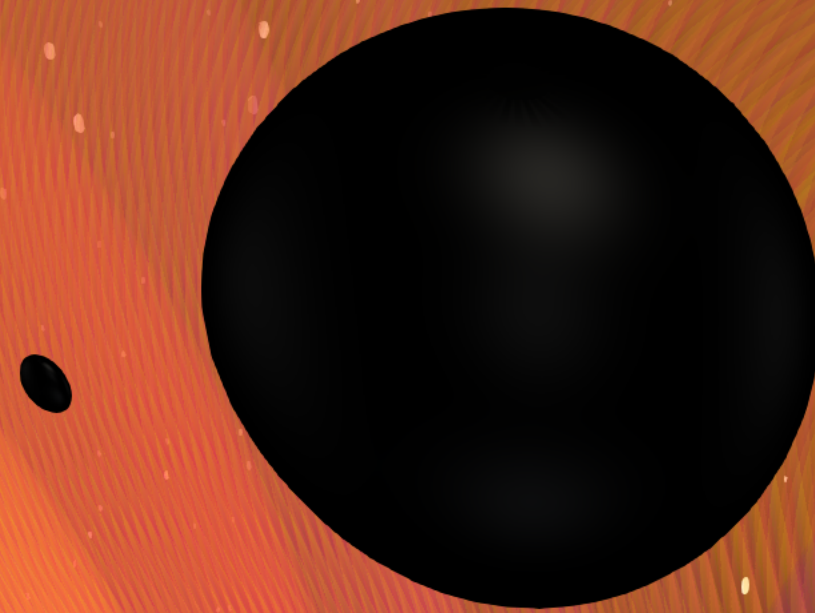


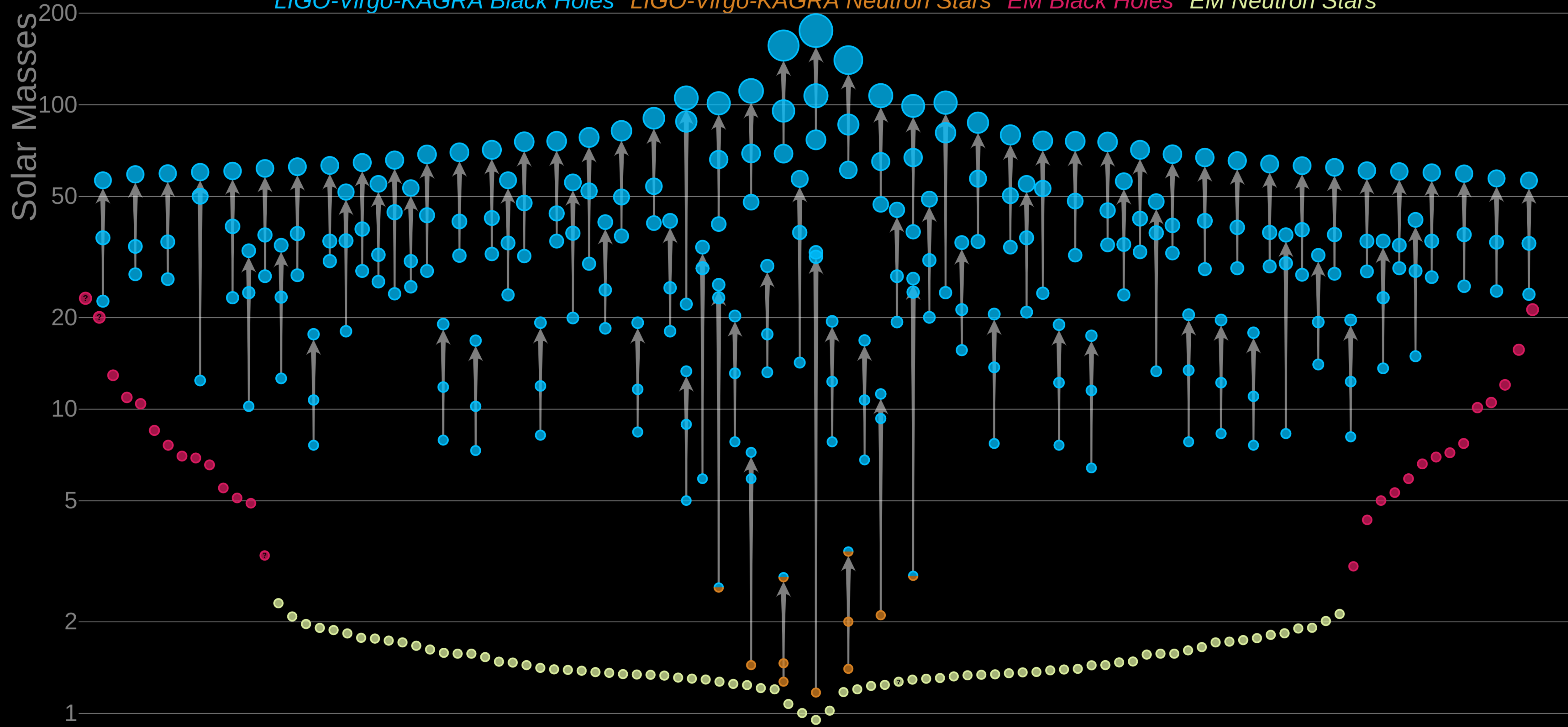
Frontiers in Numerical Relativity: Binary black holes with high mass- ratio and eccentricity



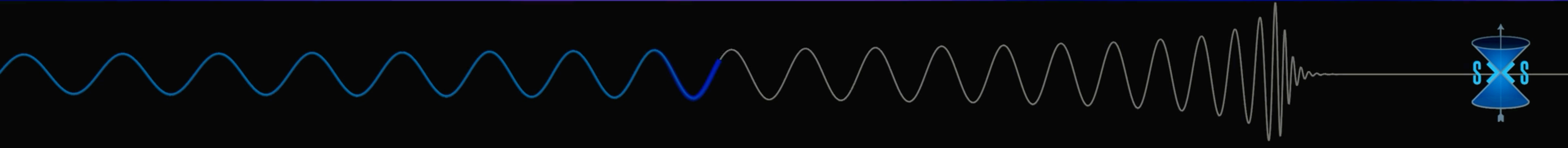
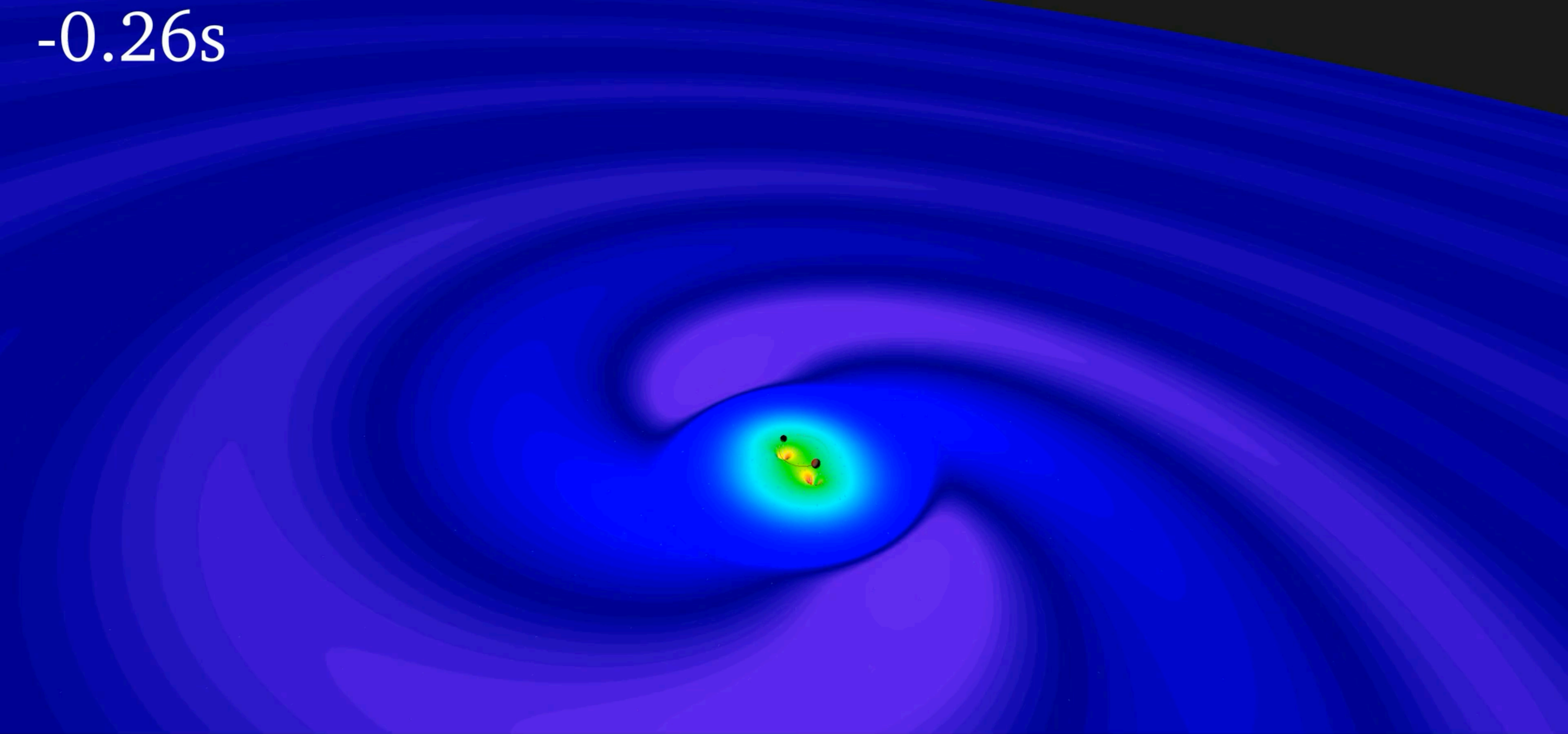
Harald Pfeiffer
MPI for Gravitational Physics
DESY Theory Seminar
Zeuthen
Dec 14, 2023

LIGO-Virgo-KAGRA Gravitational Wave Observations

LIGO-Virgo-KAGRA Black Holes LIGO-Virgo-KAGRA Neutron Stars EM Black Holes EM Neutron Stars



-0.26s



$$L_{\text{GW peak}} \sim 10^{23} L_{\odot} \sim 50 L_{\text{universe}}$$

$$P = 1/20s$$

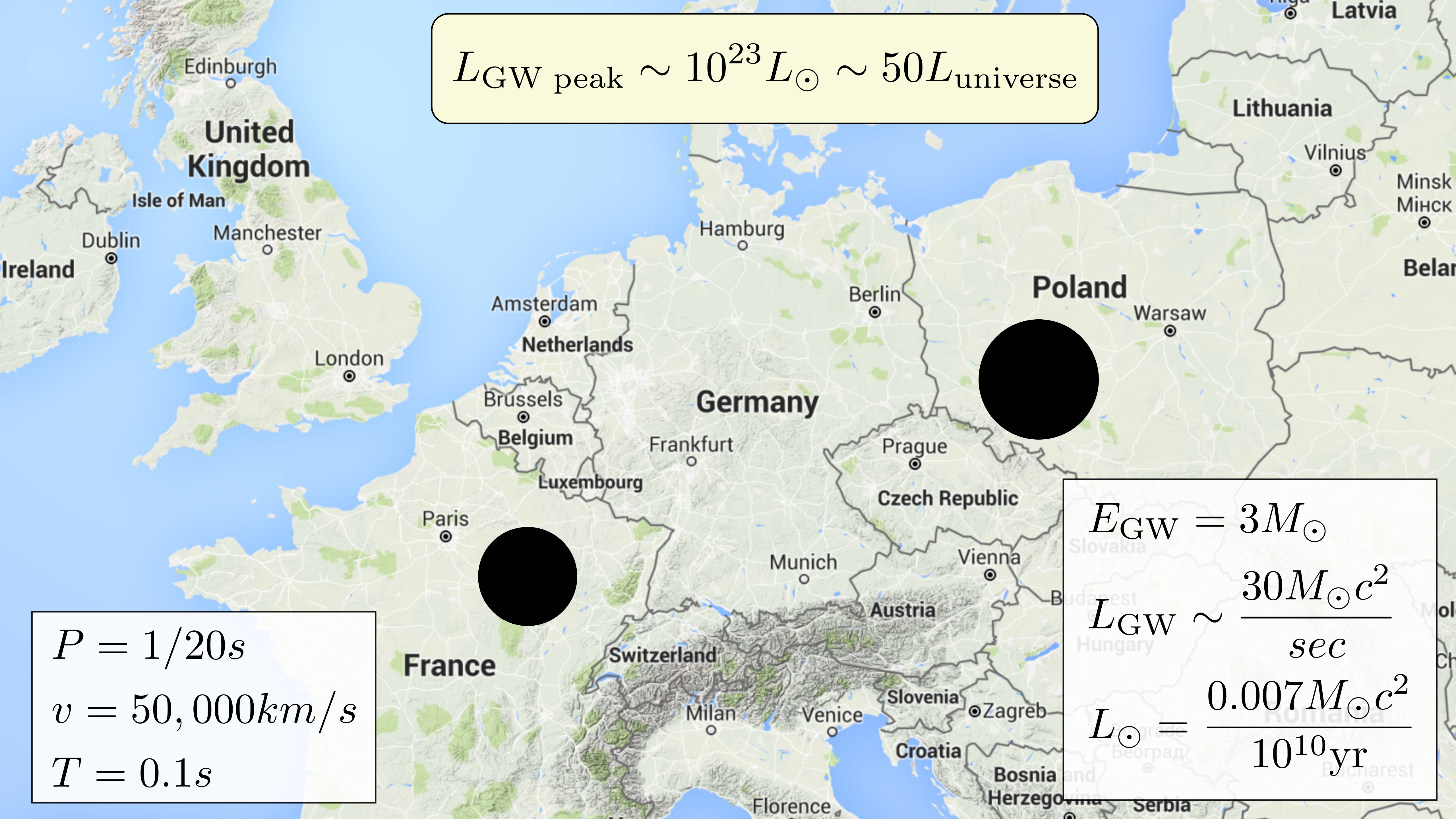
$$v = 50,000 \text{ km/s}$$

$$T = 0.1s$$

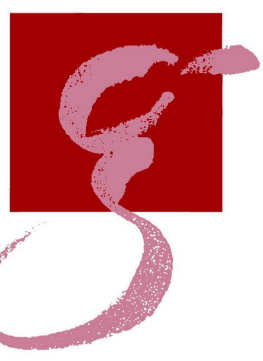
$$E_{\text{GW}} = 3M_{\odot}$$

$$L_{\text{GW}} \sim \frac{30M_{\odot}c^2}{\text{sec}}$$

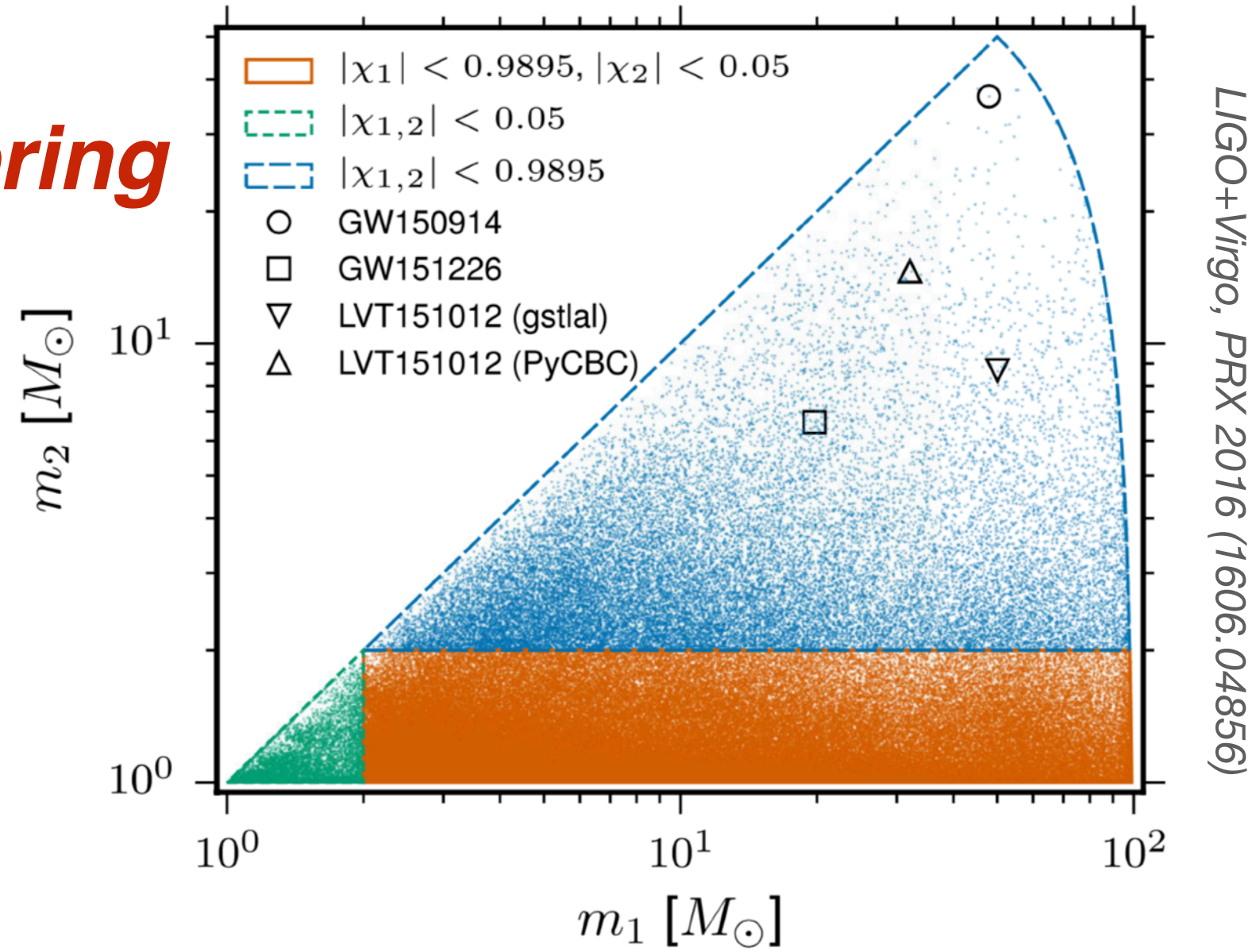
$$L_{\odot} = \frac{0.007M_{\odot}c^2}{10^{10}\text{yr}}$$



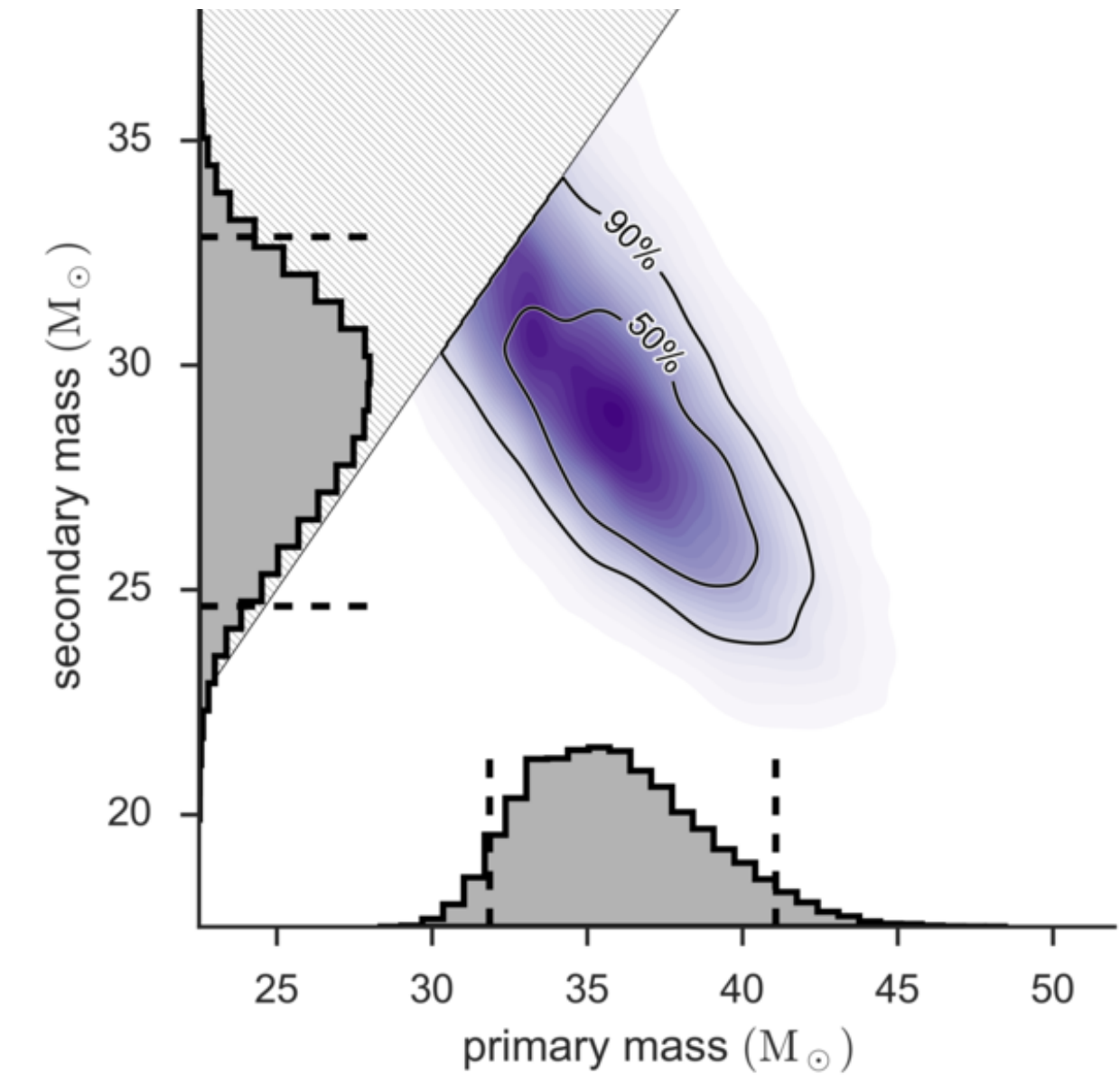
Waveform knowledge *essential for GW astronomy*



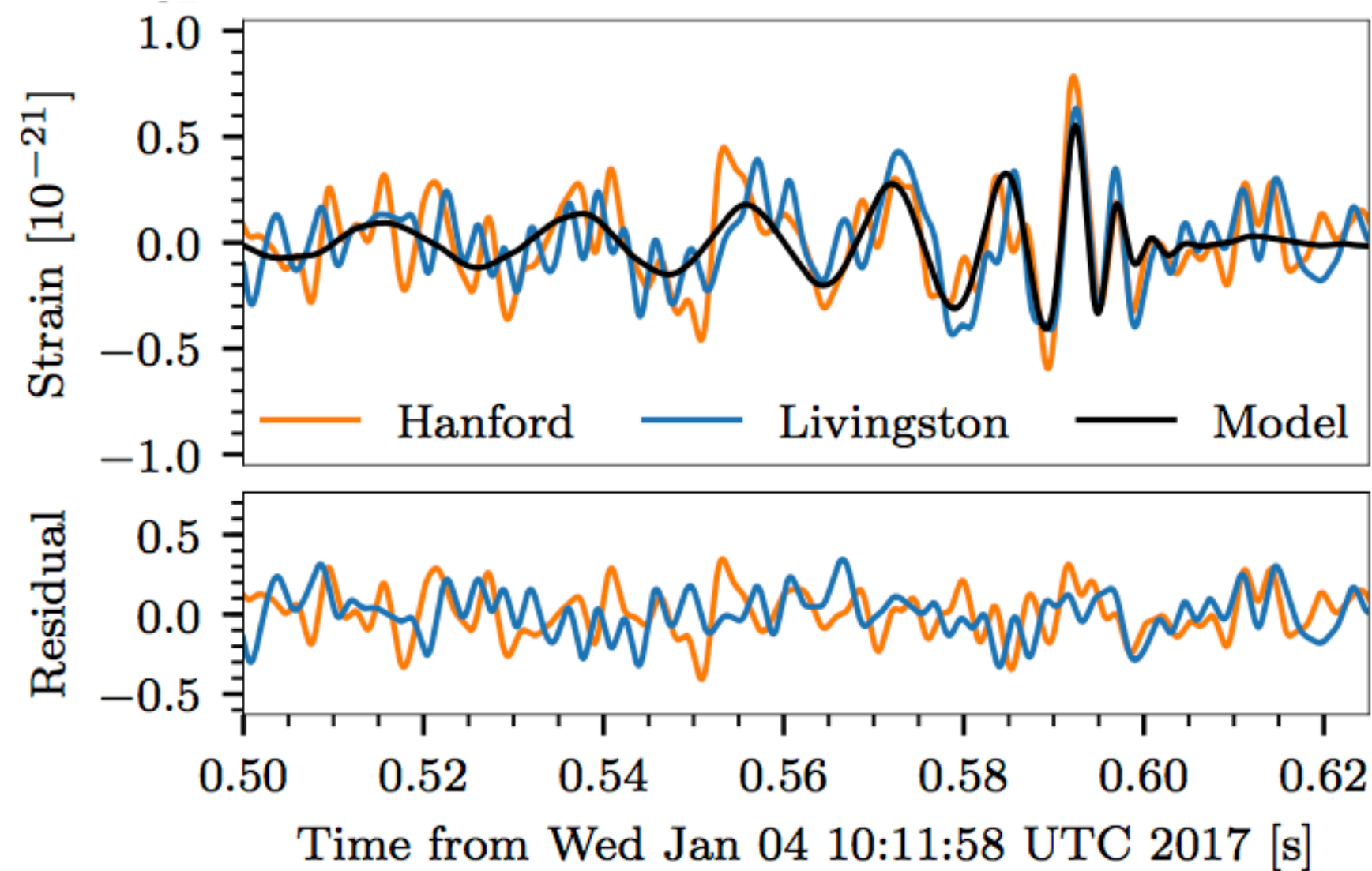
Detection by matched filtering



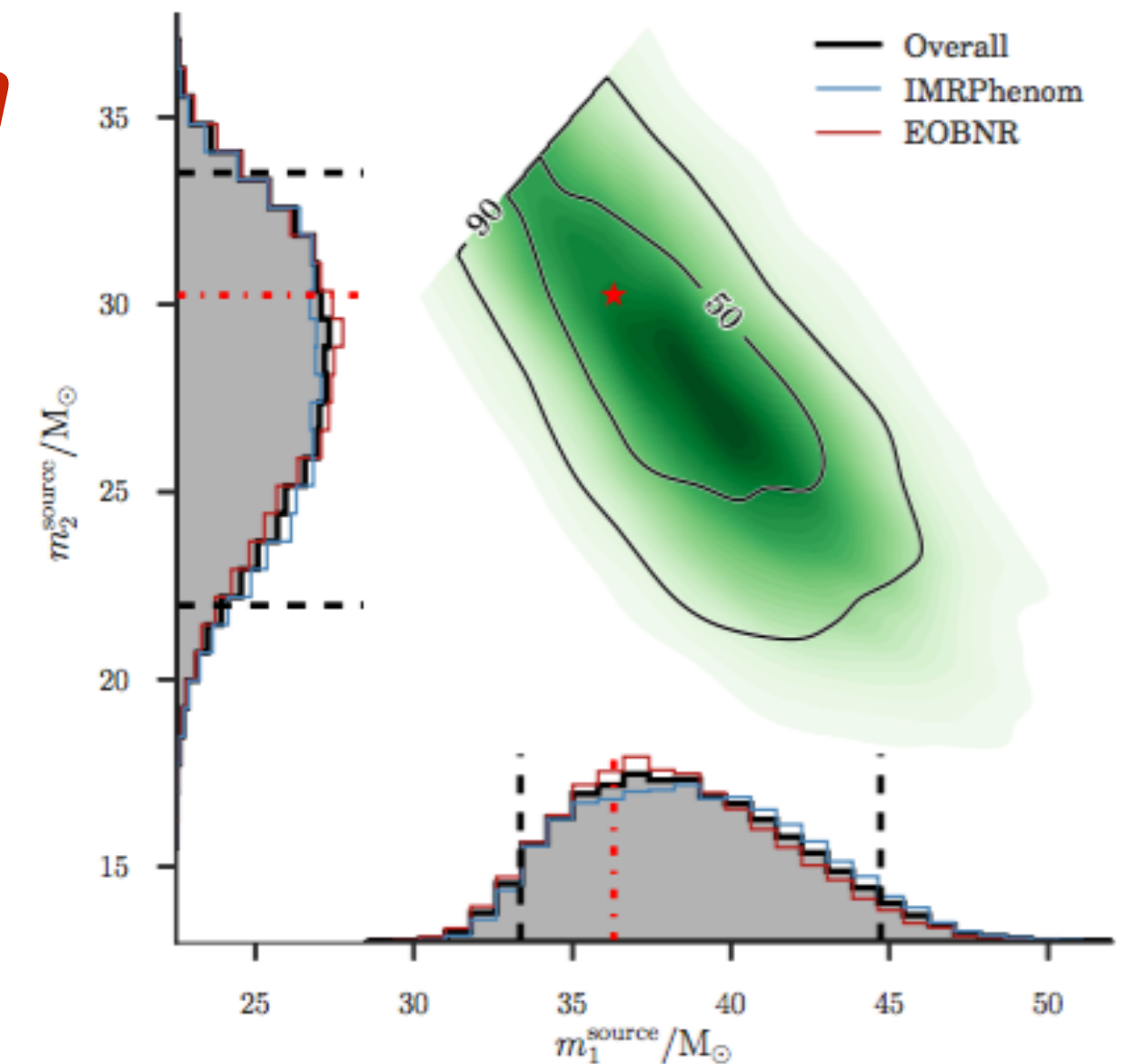
Parameter estimation



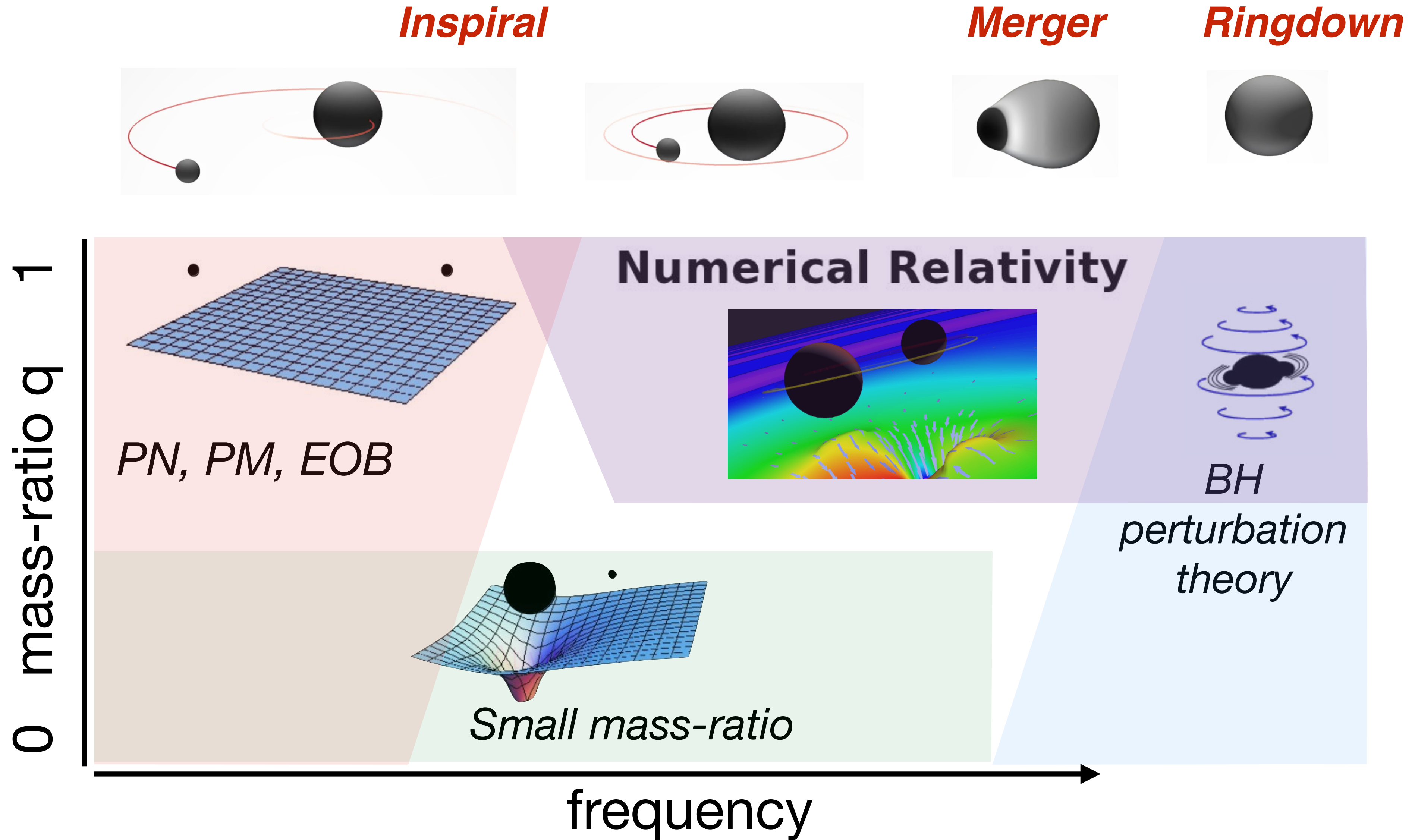
Testing GR



Validation



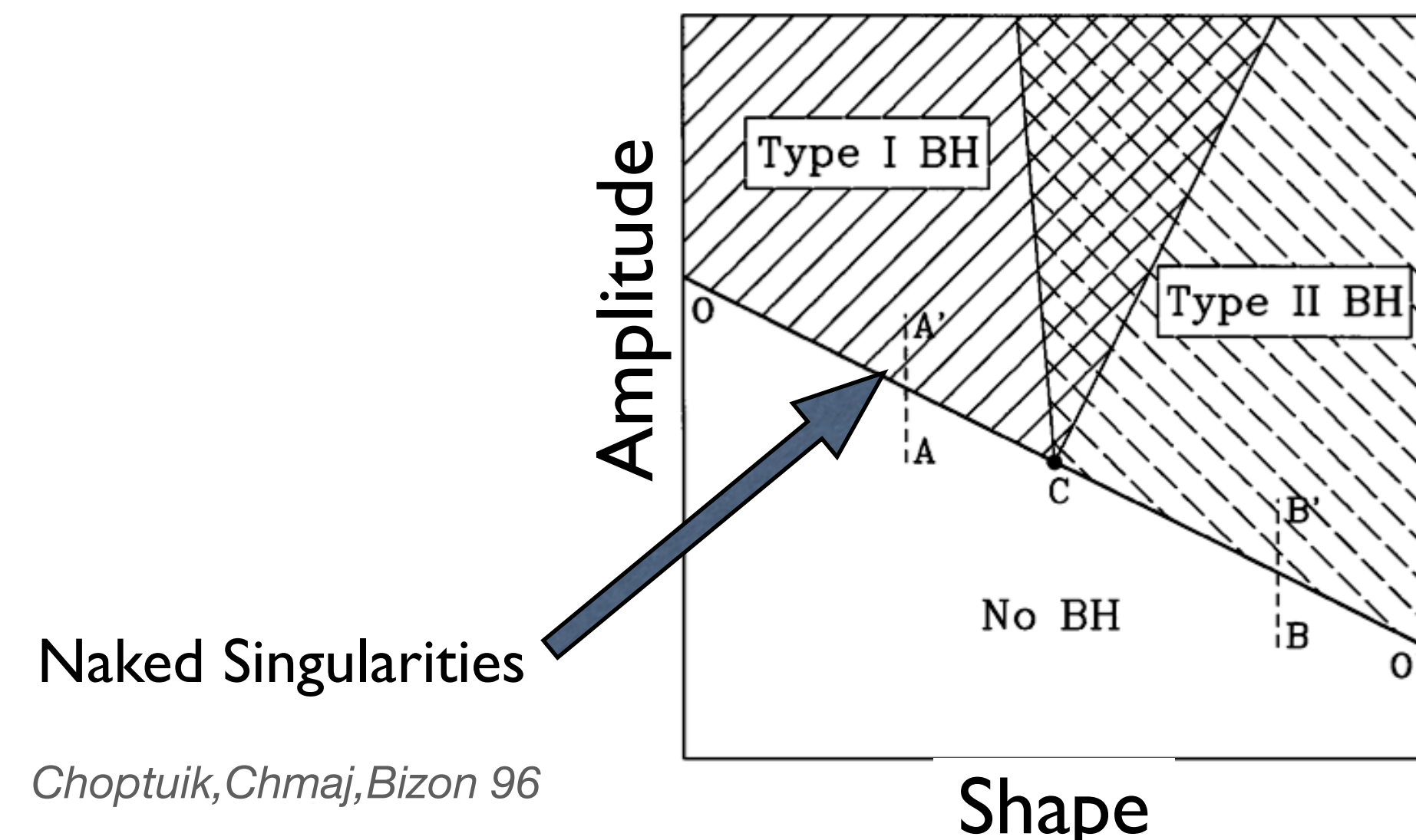
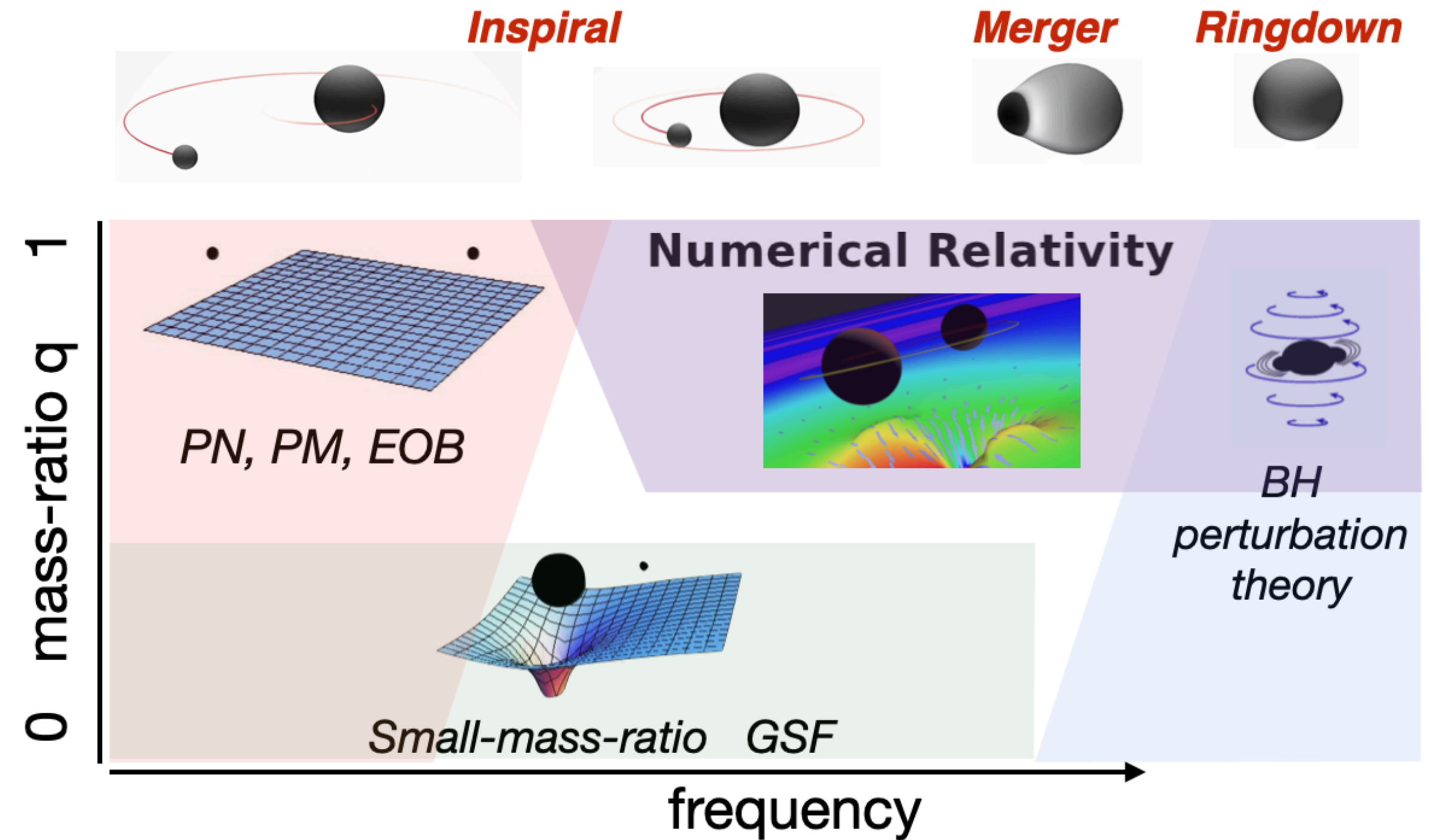
Methods for modeling BBH



Role of NR

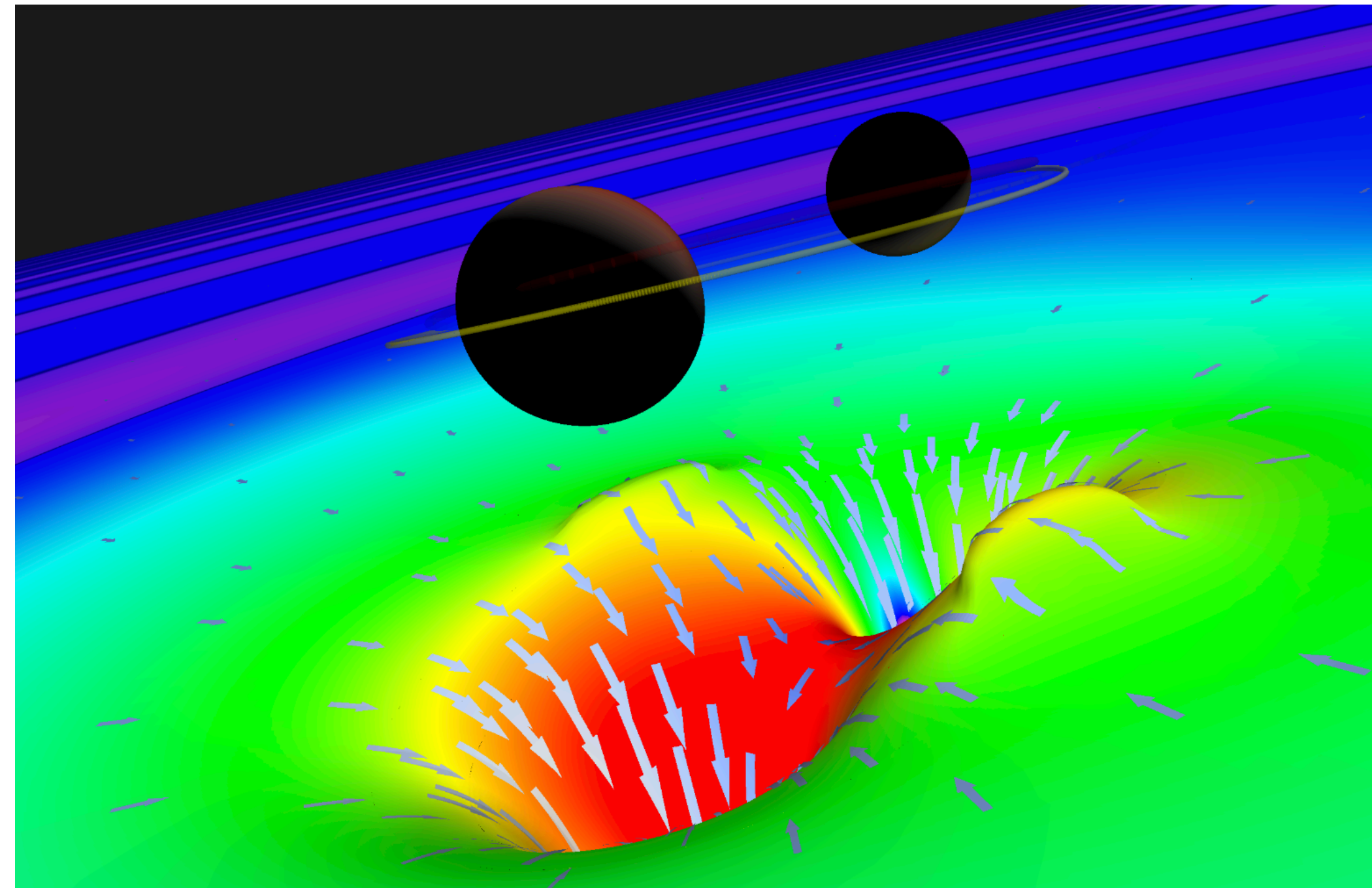


- **Solutions of GR** for GW-Astro for late inspiral + merger
 - cover parameter space
 - **error estimates**
- **Regions of validity** of perturbative methods
 - all available perturbation orders needed for science
- **Properties of GR** in nonlinear & dynamic regime

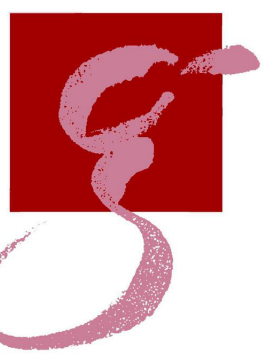


Choptuik, Chmaj, Bizon 96

Numerical Relativity



Solving Einstein Equations - Basic idea



- Goal: Space-time metric $g_{ab}(x^i, t)$ satisfying

$$R_{ab}[g_{ab}] = 0$$

- Split spacetime into space and time
- Evolution equations

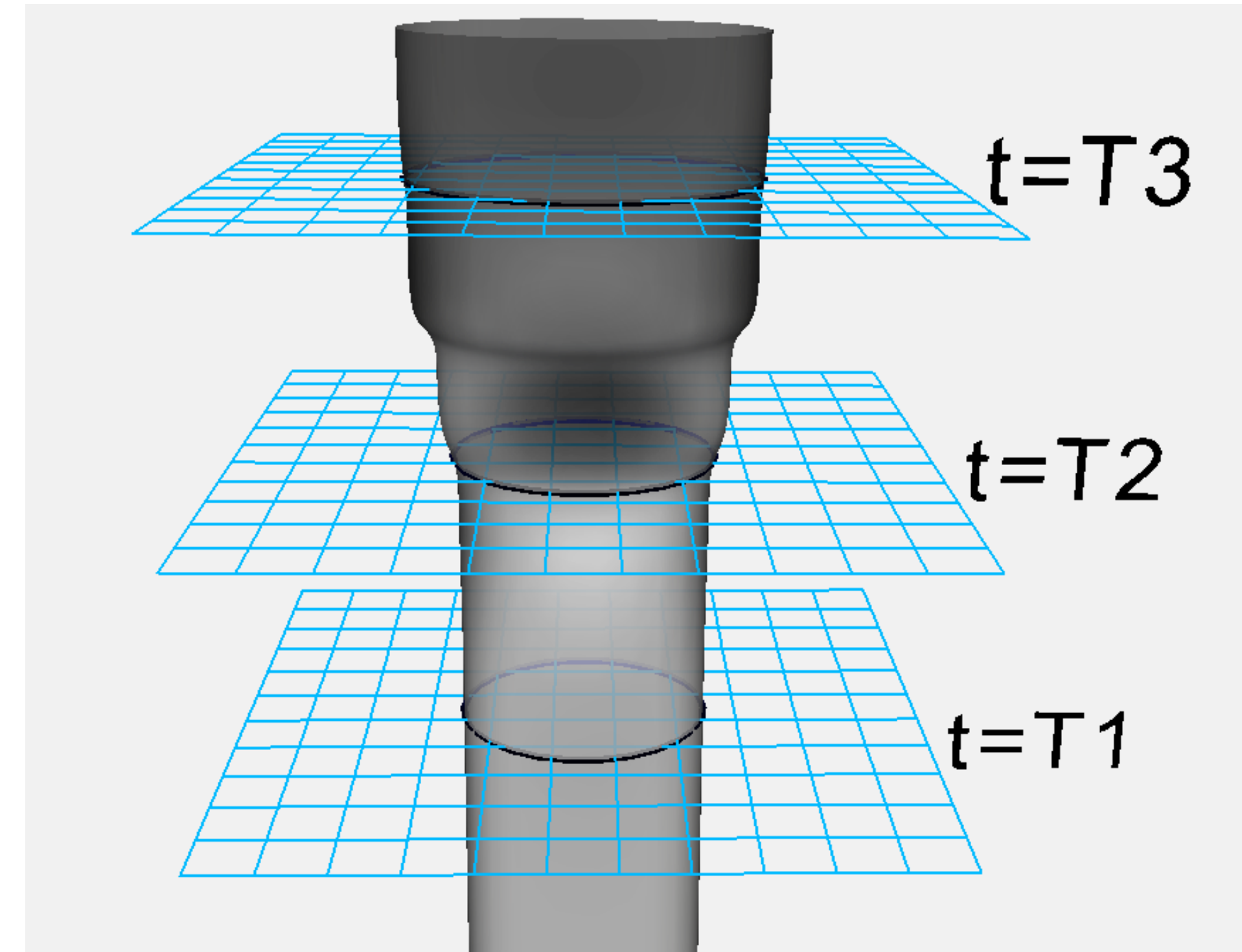
$$\partial_t g_{ij} = \dots$$

$$\partial_t K_{ij} = \dots$$

- Constraints

$$R[g_{ij}] + K^2 - K_{ij}K^{ij} = 0$$

$$\nabla_j (K^{ij} - g^{ij}K) = 0$$



cf. Maxwell's equations

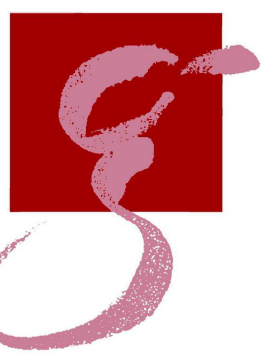
$$\partial_t \vec{E} = \nabla \times \vec{B}$$

$$\partial_t \vec{B} = -\nabla \times \vec{E}$$

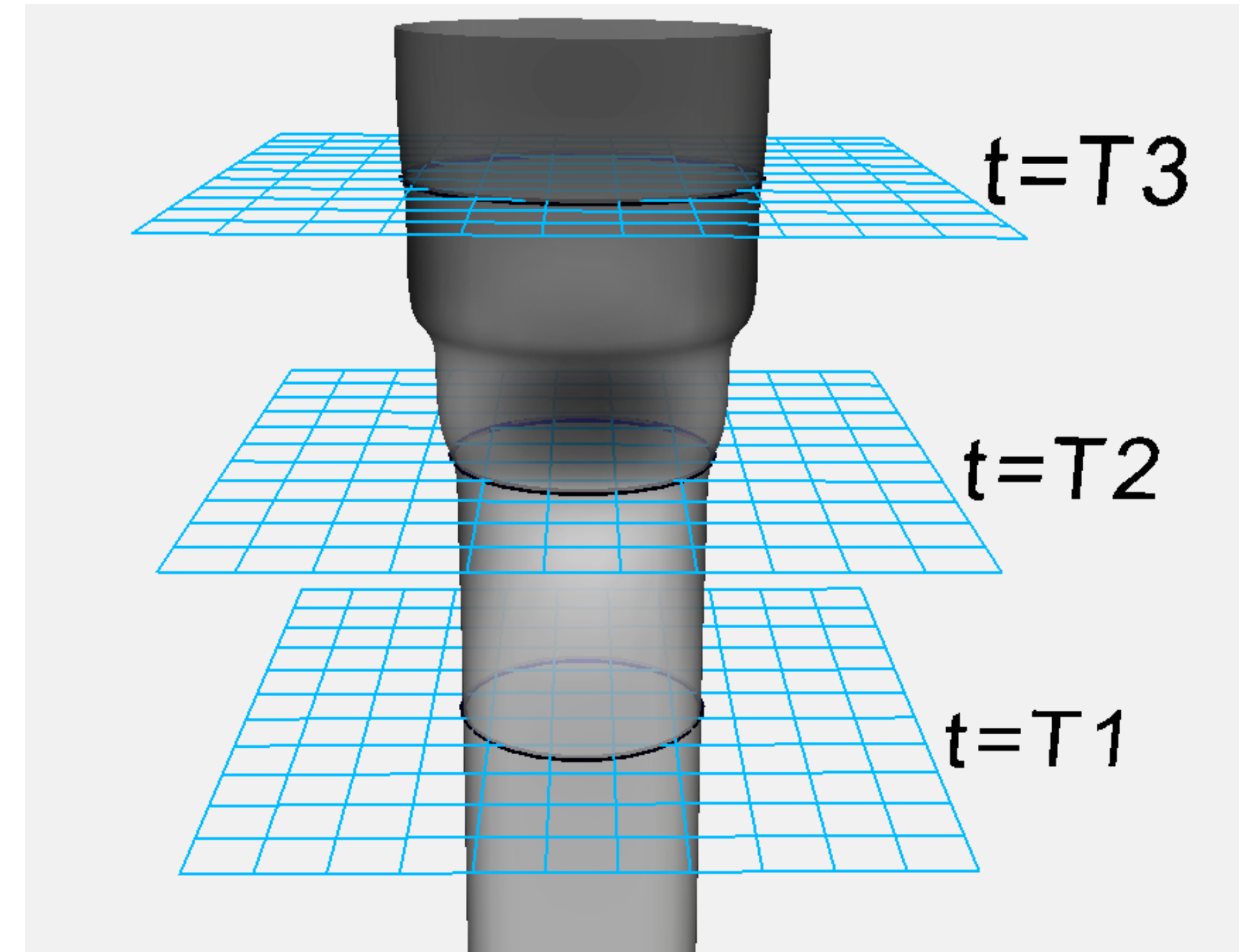
$$\nabla \cdot \vec{E} = 0$$

$$\nabla \cdot \vec{B} = 0$$

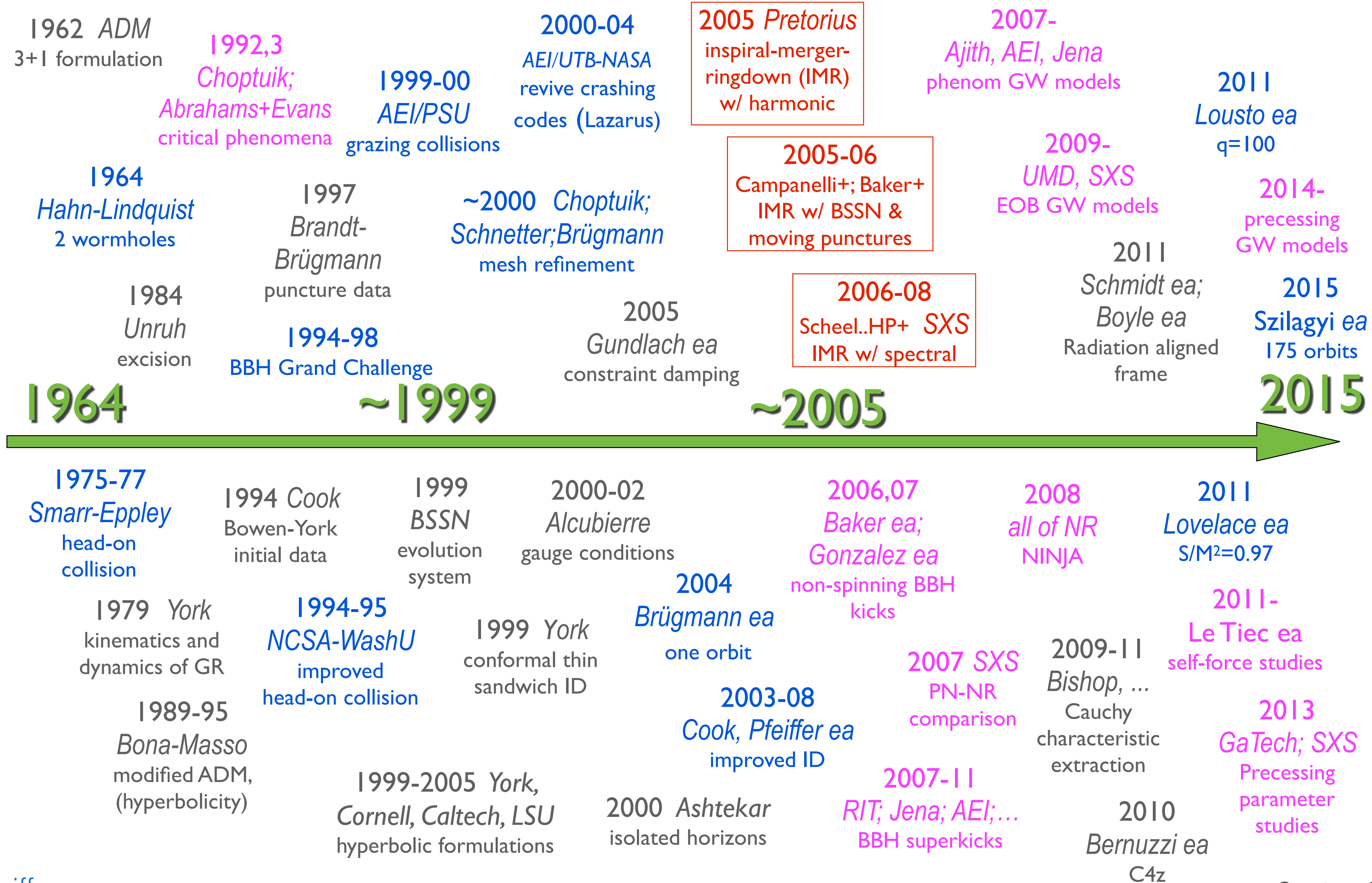
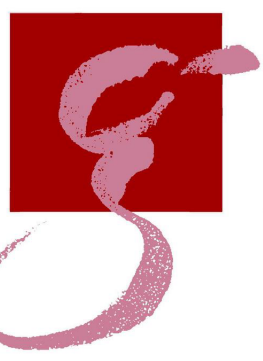
Why is this hard?



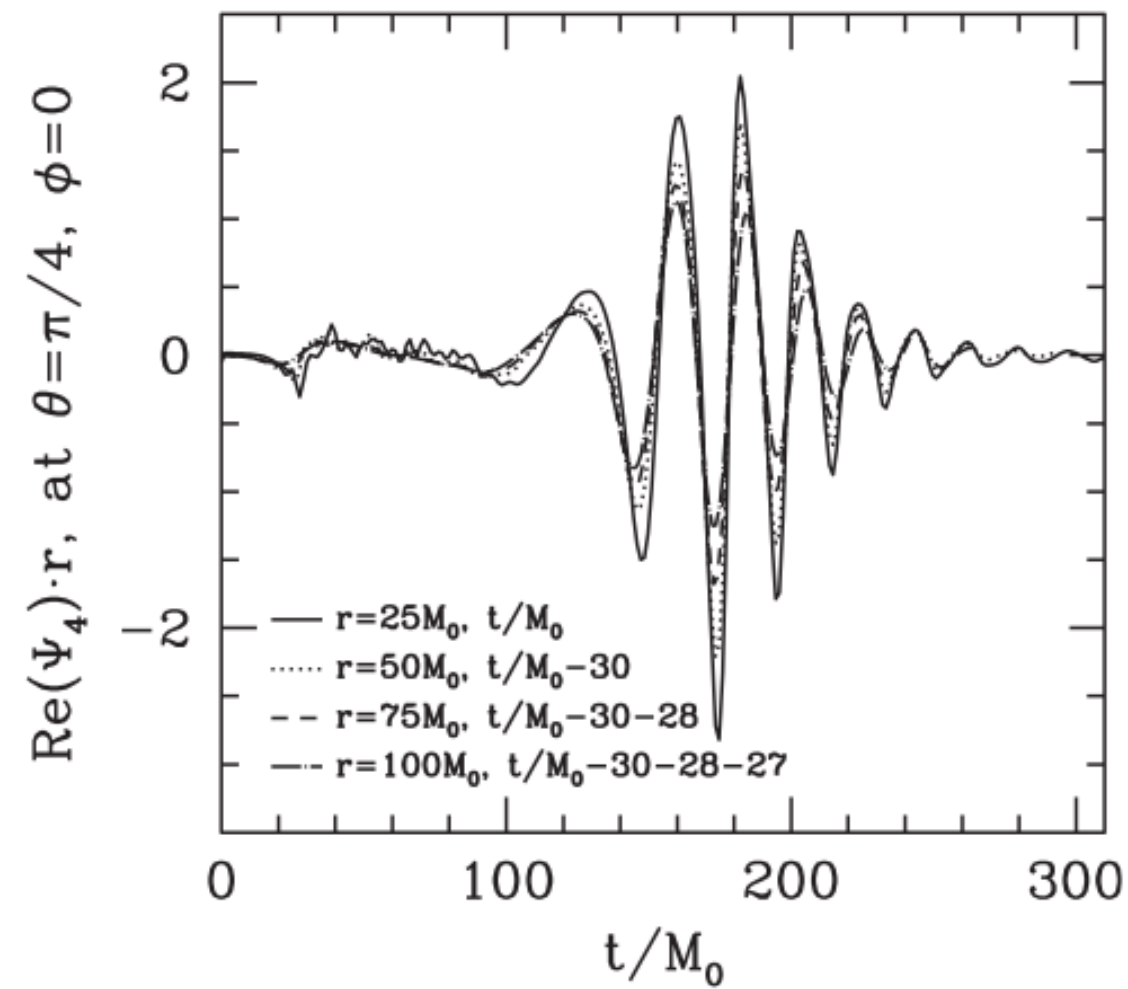
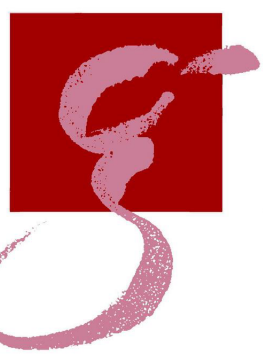
- ADM equations ill-posed; rewrite as hyperbolic system
- **Singularities** inside black holes
- Constraints difficult to preserve
- **Coordinate freedom**
 - How to choose coordinates for a space-time one does not know yet?
- Many common numerical challenges
 - 20-50 variables
 - 10,000 FLOP / grid-point / time-step
 - Different length scales
 - **High accuracy requirements**



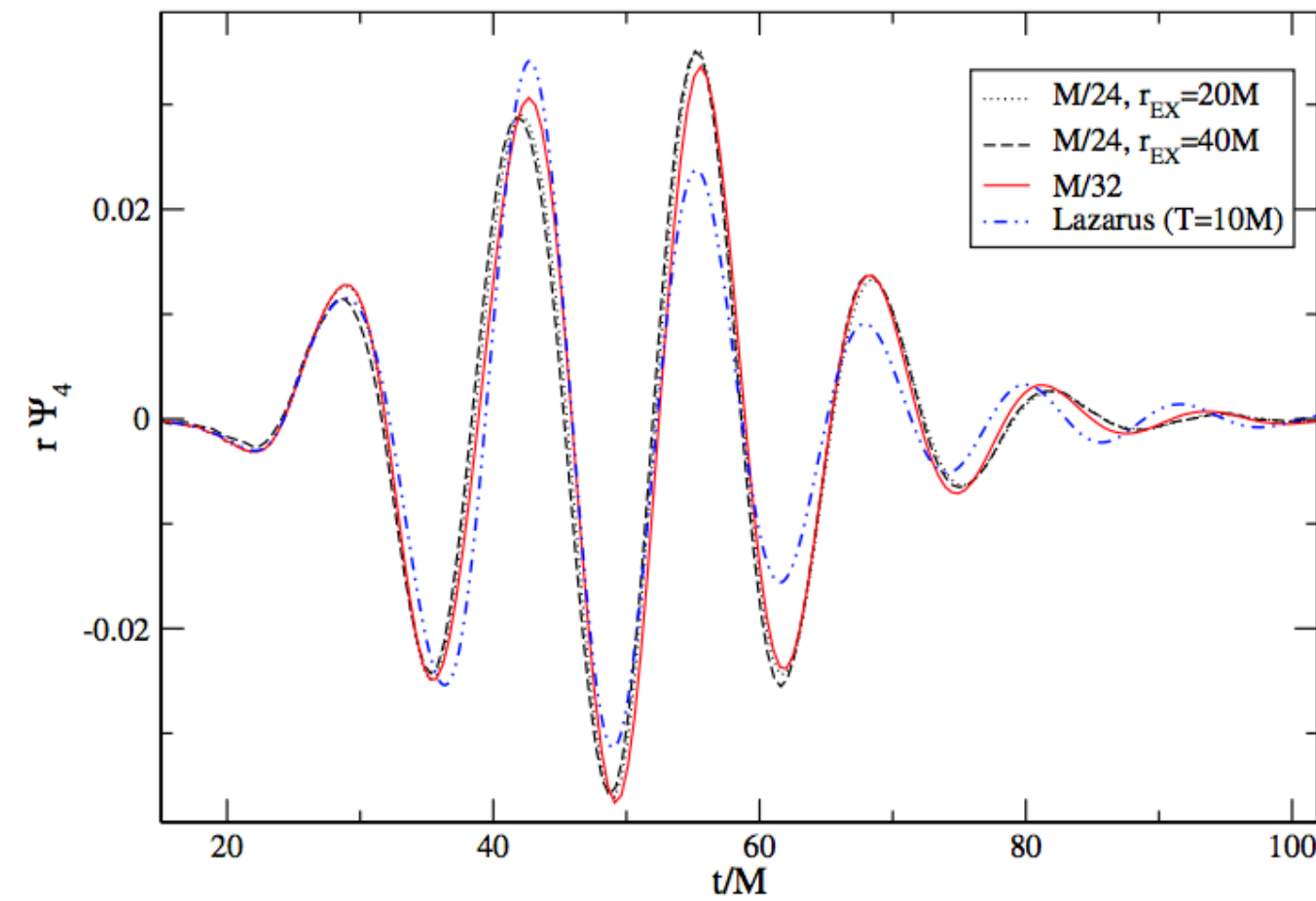
The first 50 Years of numerical relativity for BBH



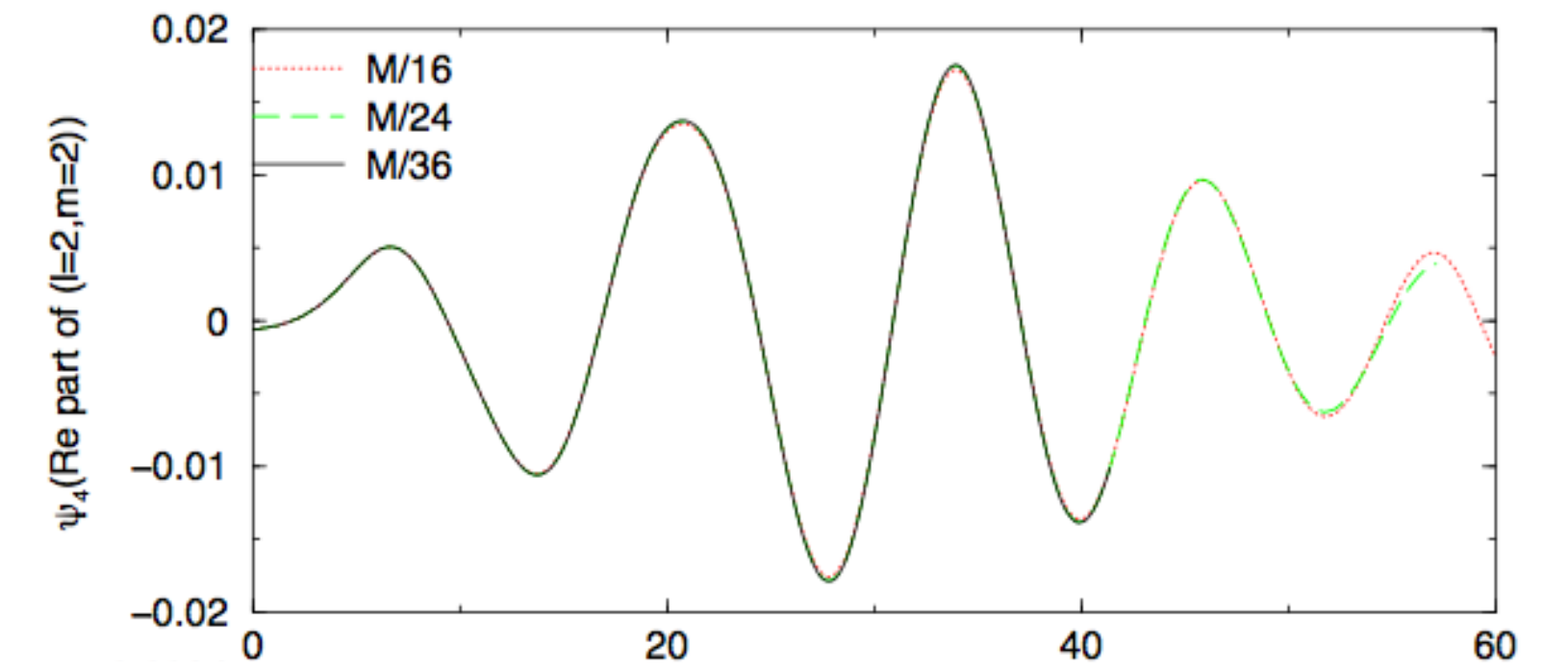
2005: First working BBH inspirals



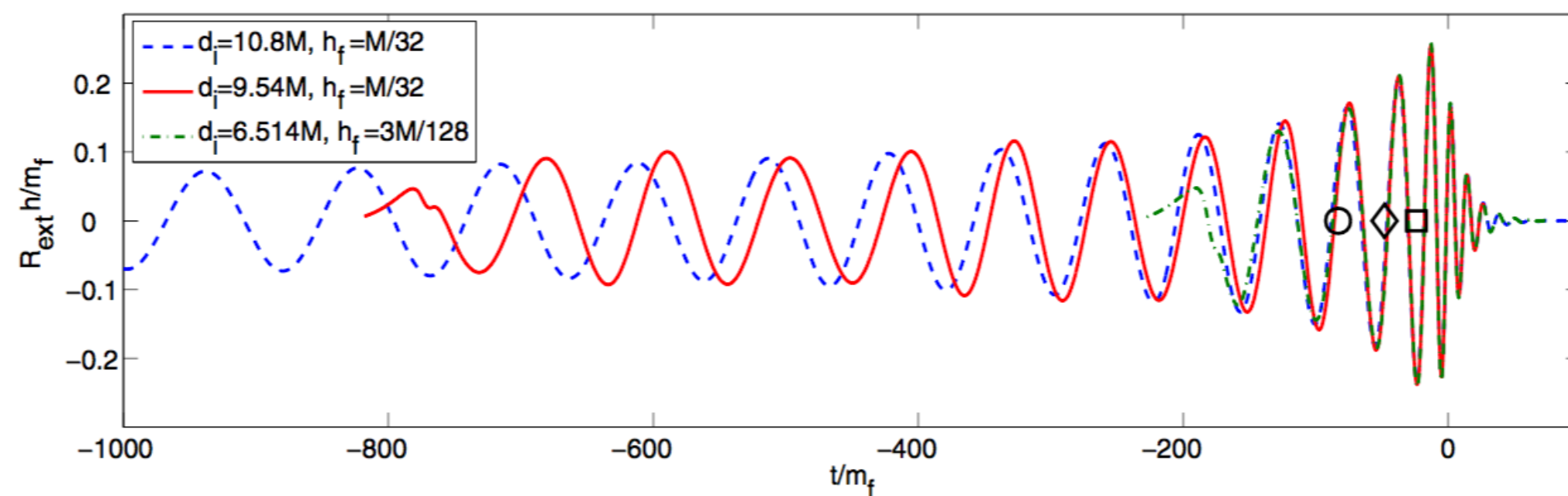
Pretorius 05



Baker+06



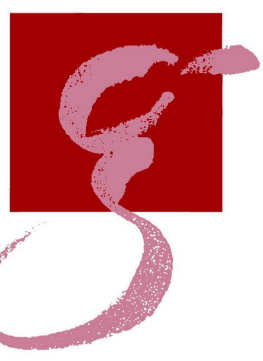
Campanelli+06



Baker+07

Important early result:
Simplicity of merger
 Continuous transition
 inspiral → ringdown

Major approaches towards BBH simulations



“BSSN & Moving punctures”

Puncture initial-data

$\chi \lesssim 0.9$ (but see Zlochower+ 17)

BSSN or CC4z

Moving puncture

mergers “easy”

Sommerfeld outer BC

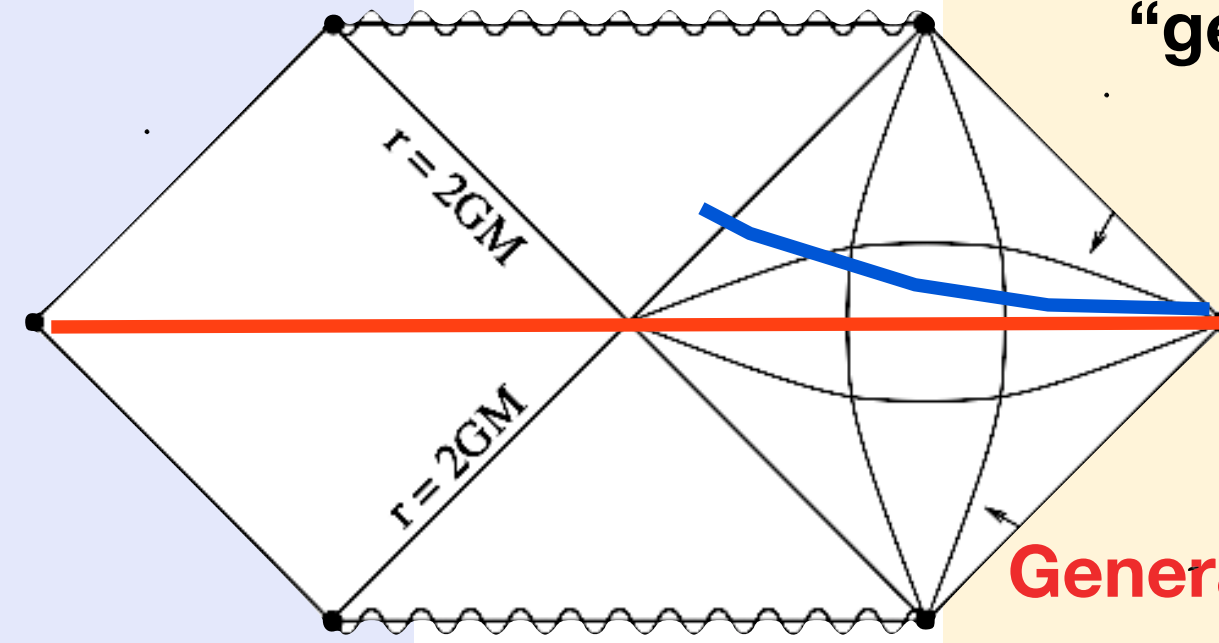
4th to 8th order finite-difference

BHs advect through static grids

GW extrapolation

(Healy,Lousto '20 for LazEv COM correction)

LazEv, Maya, ETK, BAM,
Goddard, GRchombo, ...



“generalized harmonic & spectral”

Quasi-equilibrium excision data

$\chi \lesssim 0.999$

Generalized-Harmonic Evolution System

BH excision

mergers difficult

Constraint preserving, minimally reflective outer BC

Spectral methods

Moving grid

long, phase-accurate inspirals

GW extrapolation & center-of-mass correction

Cauchy-characteristic extraction

accurate $m=0$ modes, GW memory

SpEC (SXS collaboration)

Spectral Einstein Code (SpEC)

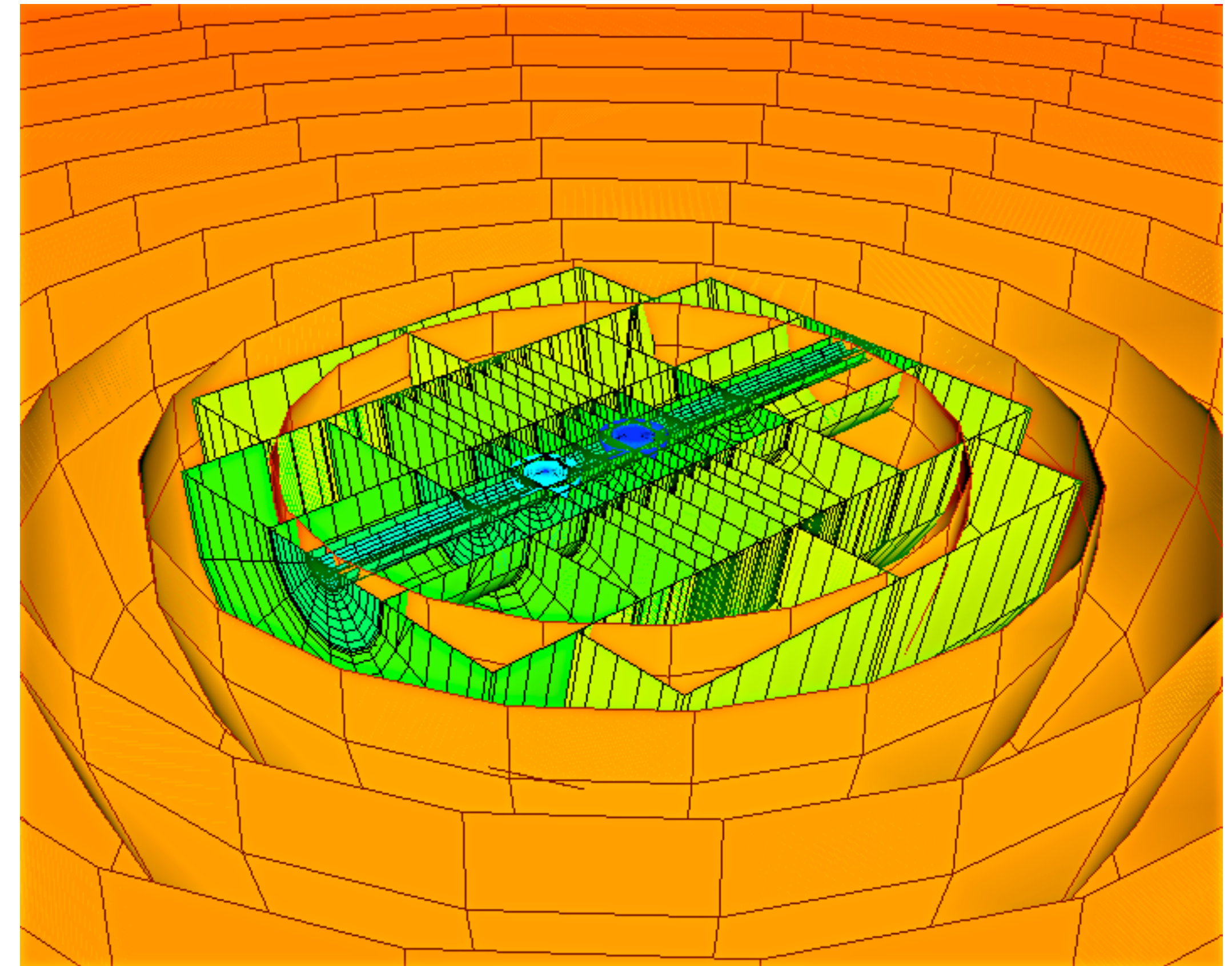


Simulating eXtreme Spacetimes collaboration



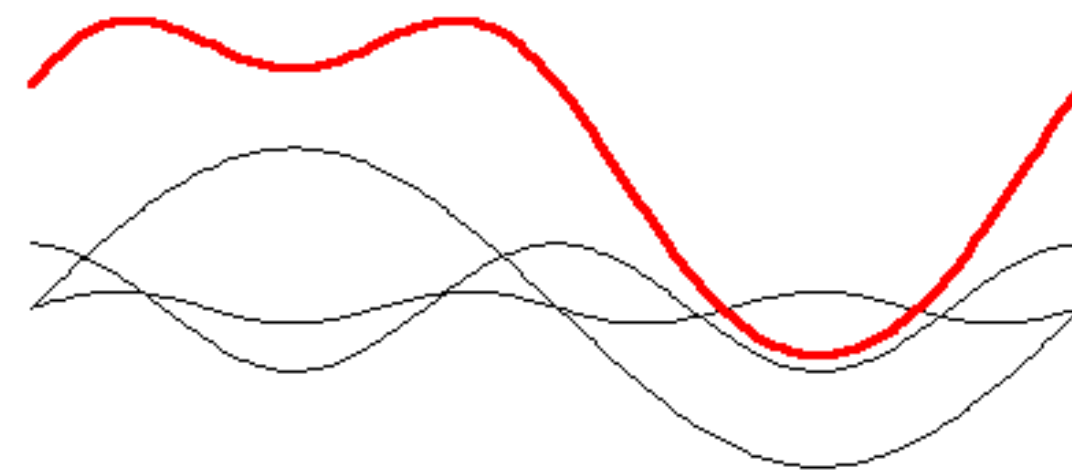
<http://www.black-holes.org/SpEC.html>

Combine w/ domain-decomposition



- **Expand in basis-functions**

$$u(x, t) = \sum_{k=1}^N \tilde{u}_k(t) \Phi_k(x)$$



- **Compute derivatives analytically**

$$u'(x, t) = \sum_{k=1}^N \tilde{u}_k(t) \Phi'_k(x)$$

- **Exponentially fast convergence**
 - for smooth problems

Einstein constraints: Formalism



$$R + (\text{tr}K)^2 - K^2 = 0$$

$$\nabla \cdot (K - g \text{tr}K) = 0$$

$$g = \psi^4 \tilde{g}$$

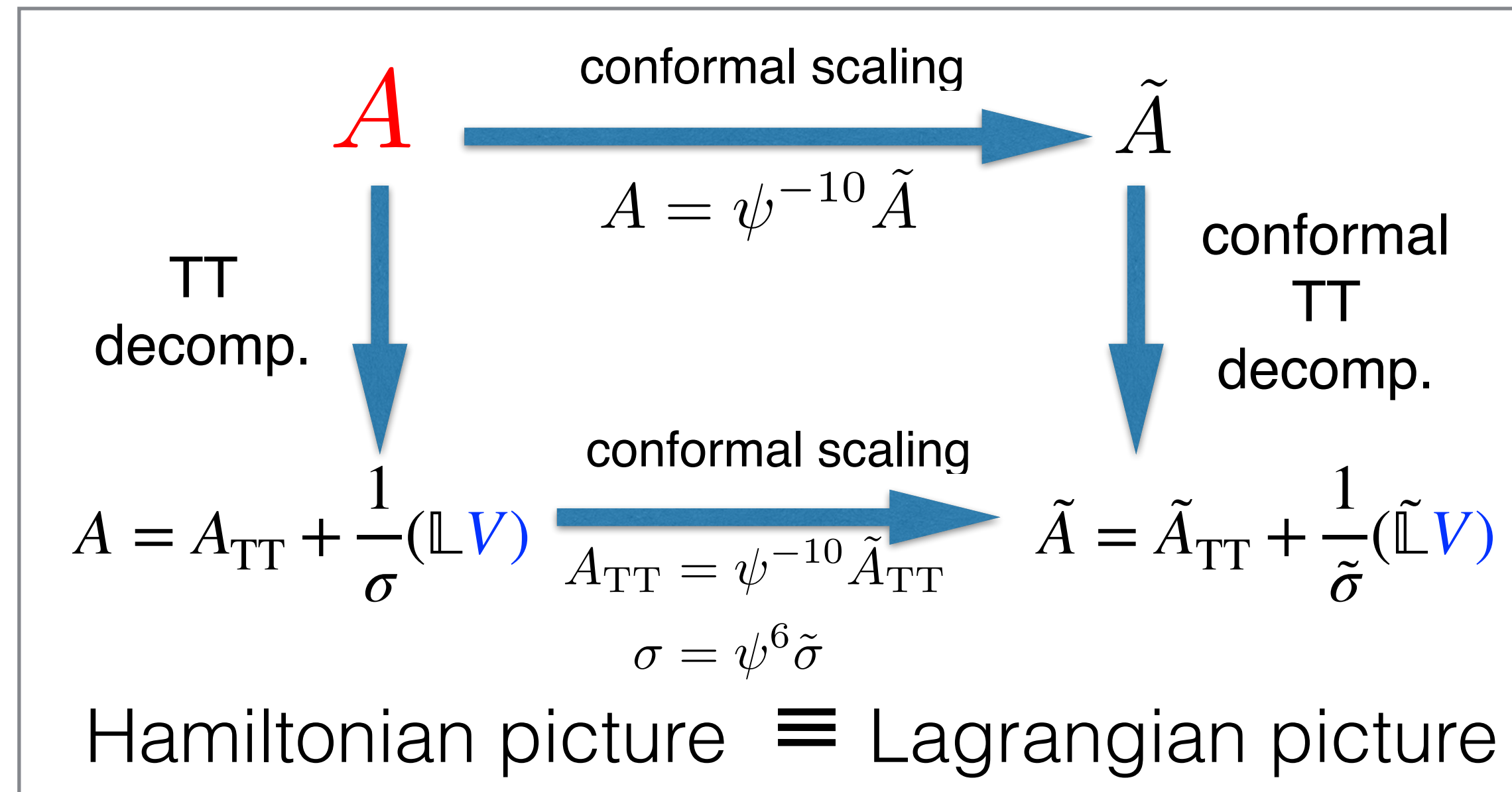
Lichnerowicz 44

$$\tilde{\nabla}^2 \psi = \dots$$

$$\tilde{\nabla} \cdot \left(\frac{1}{\tilde{\sigma}} \tilde{\mathbb{L}} V \right) = \dots$$

$$K = \frac{1}{3} \text{tr}K g + A$$

coupled nonlinear elliptic PDEs in 3D



York(+) 72; 74; 99, HP, York 03

Applied to binary black holes



- **Asymptotics/boundary conditions**

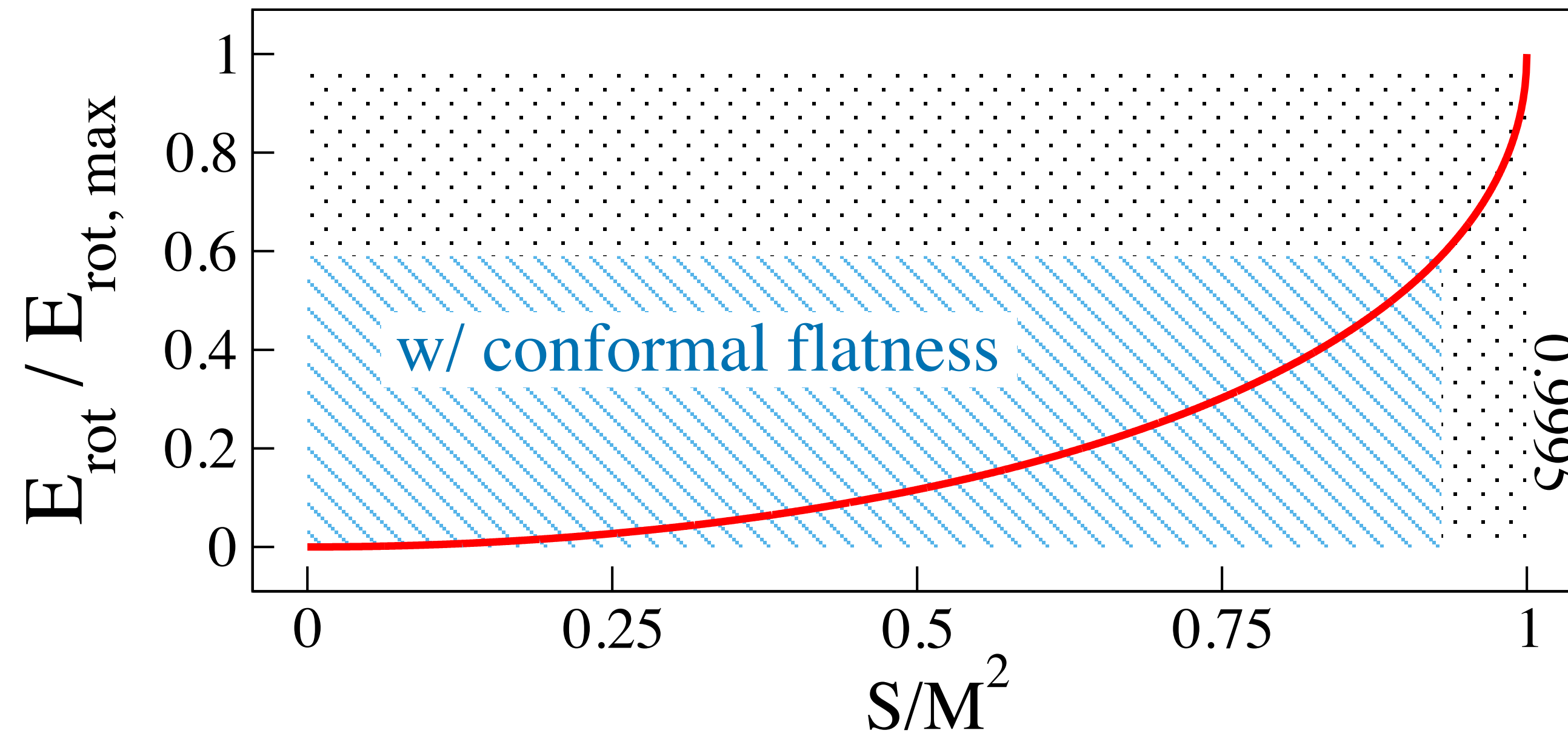
Brandt, Brügmann 97; Cook, HP 04

- **Elliptic solver**

HP+ 02; Ansorg 04; Vu, HP+ 21a,b

- **Spins > 0.9**

Lovelace..HP+ 08

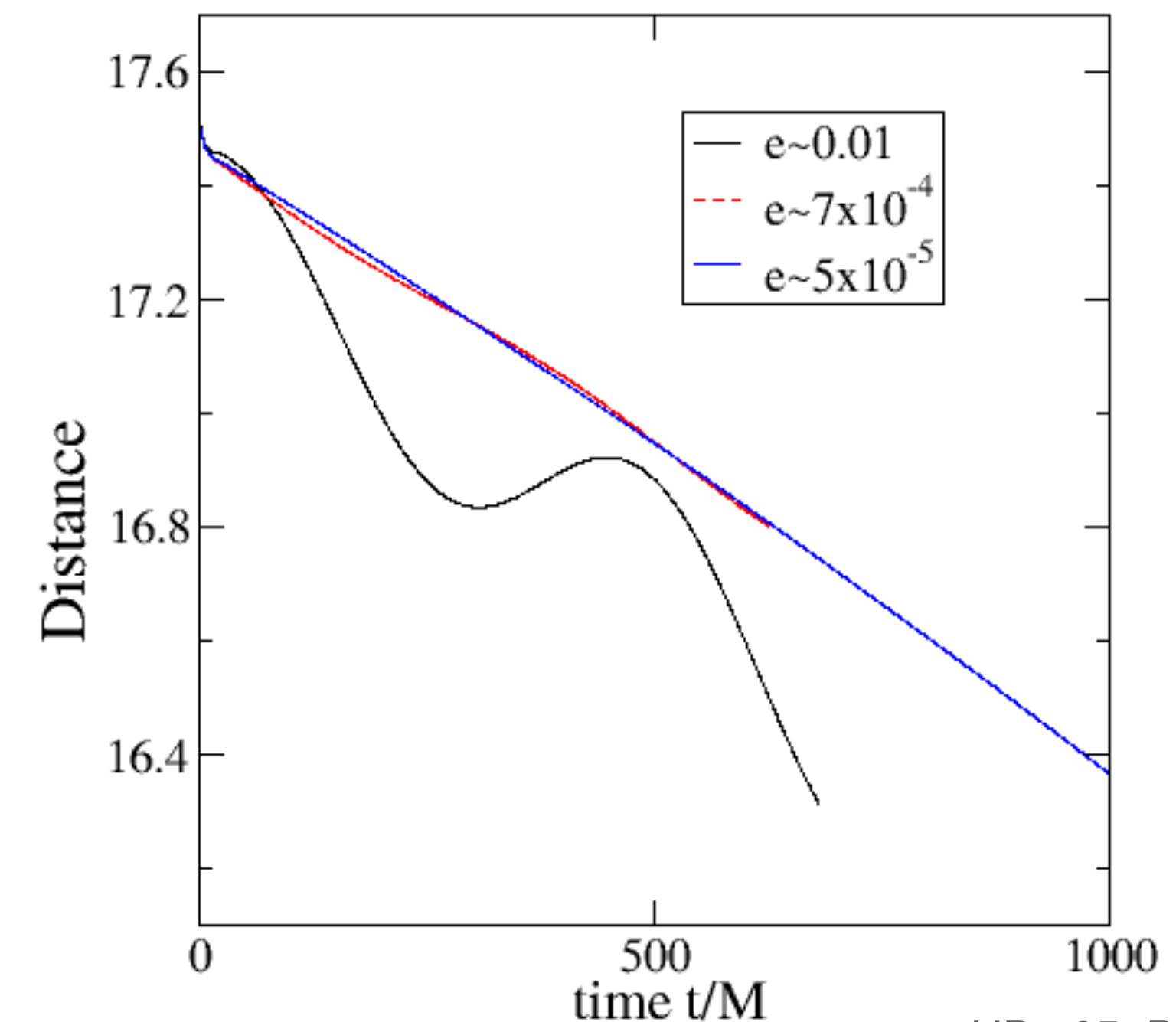


$$\tilde{\nabla}^2 \psi = \dots$$

$$\tilde{\nabla} \cdot \left(\frac{1}{\tilde{N}} \tilde{\mathbb{L}} \beta \right) = \dots$$

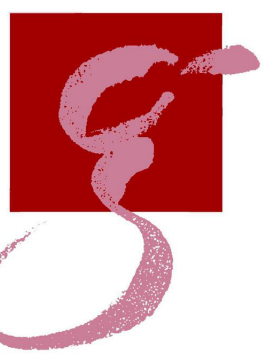
$$\tilde{\nabla}^2 \tilde{N} = \dots$$

- **Control eccentricity**



*HP+ 05; Buonanno..HP+ 08
Nee, HP+ (in prep)*

Generalized Harmonic Evolution System



- Einstein's equations

$$0 = R_{ab}[g_{ab}] = -\frac{1}{2}\square g_{ab} + \nabla_{(a}\Gamma_{b)} + \text{lower order terms}, \quad \Gamma_a = -g_{ab}\square x^b.$$

- Generalized harmonic coordinates $g_{ab}\square x^b \equiv H_a(x^a, g_{ab})$
(Friedrich 1985, Pretorius 2005; $H = 0$ used since 1920's)

$$\square g_{ab} = \text{lower order terms.}$$

$$\Rightarrow \text{Constraint } C_a \equiv H_a - g_{ab}\square x^b = 0$$

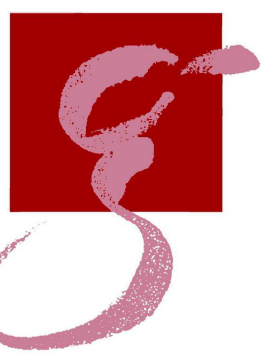
- **Constraint damping** (Gundlach, et al., Pretorius, 2005)

$$\square g_{ab} = \gamma \left[t_{(a} C_{b)} - \frac{1}{2} g_{ab} t^c C_c \right] + \text{lower order terms}$$

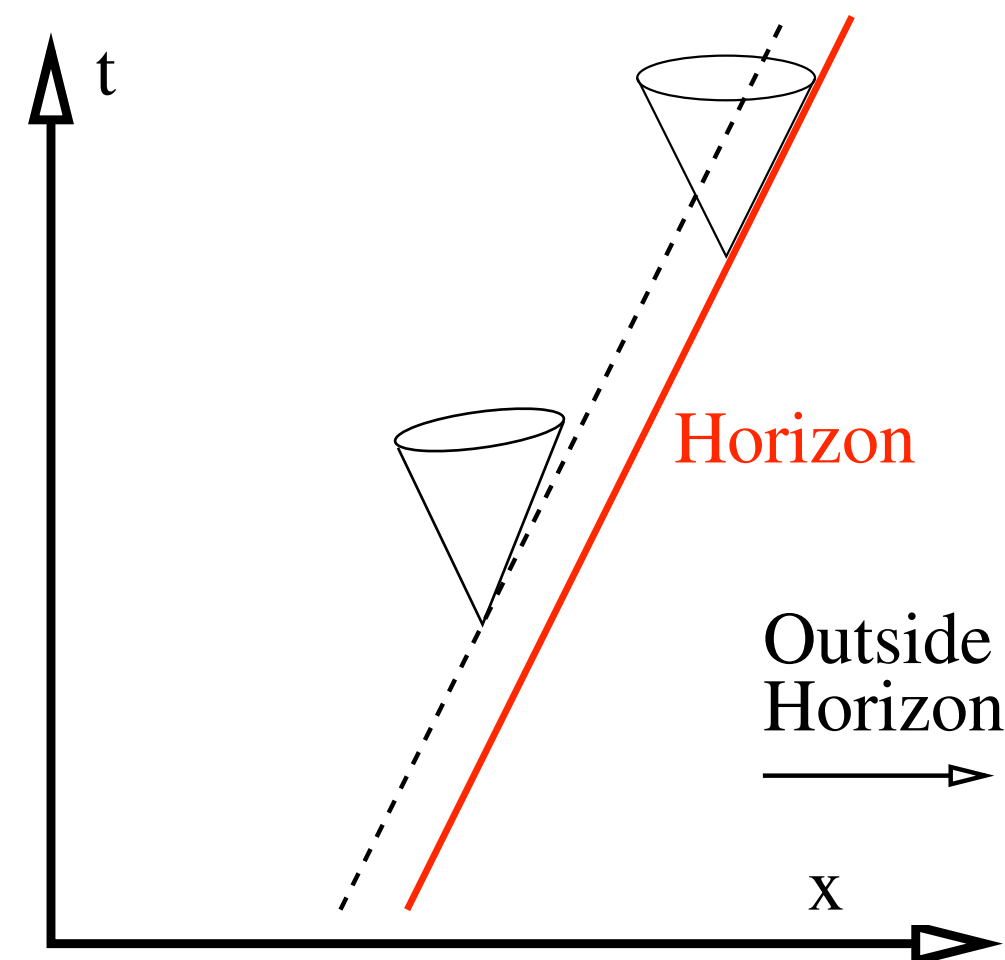
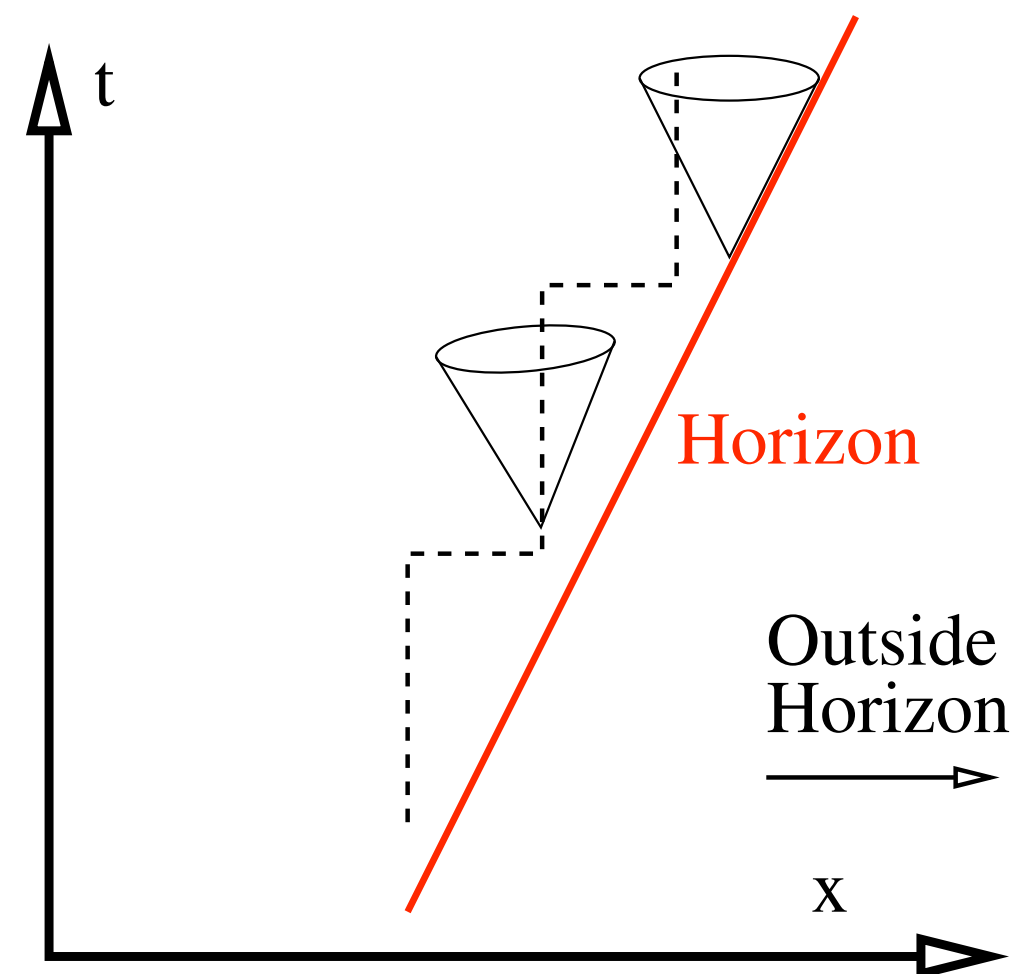
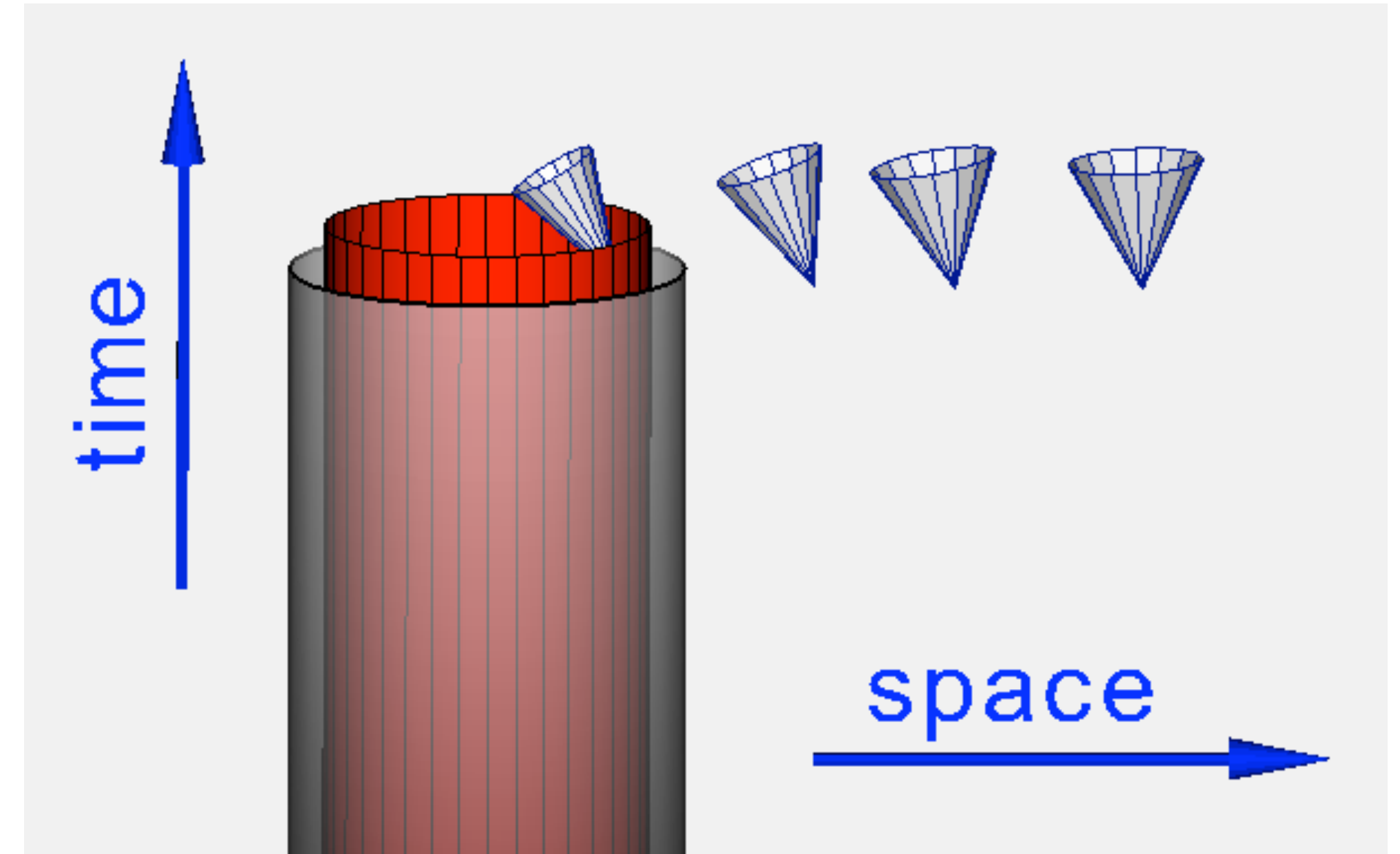
$$\partial_t C_a \sim -\gamma C_a.$$

*Excellent GH exposition:
Lindblom et al 2006*

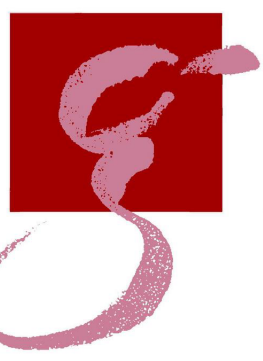
BH Excision



- Excise inside BH horizons
- Excision boundaries:
 - **follow BHs continuously**
 - conform to shape of AH



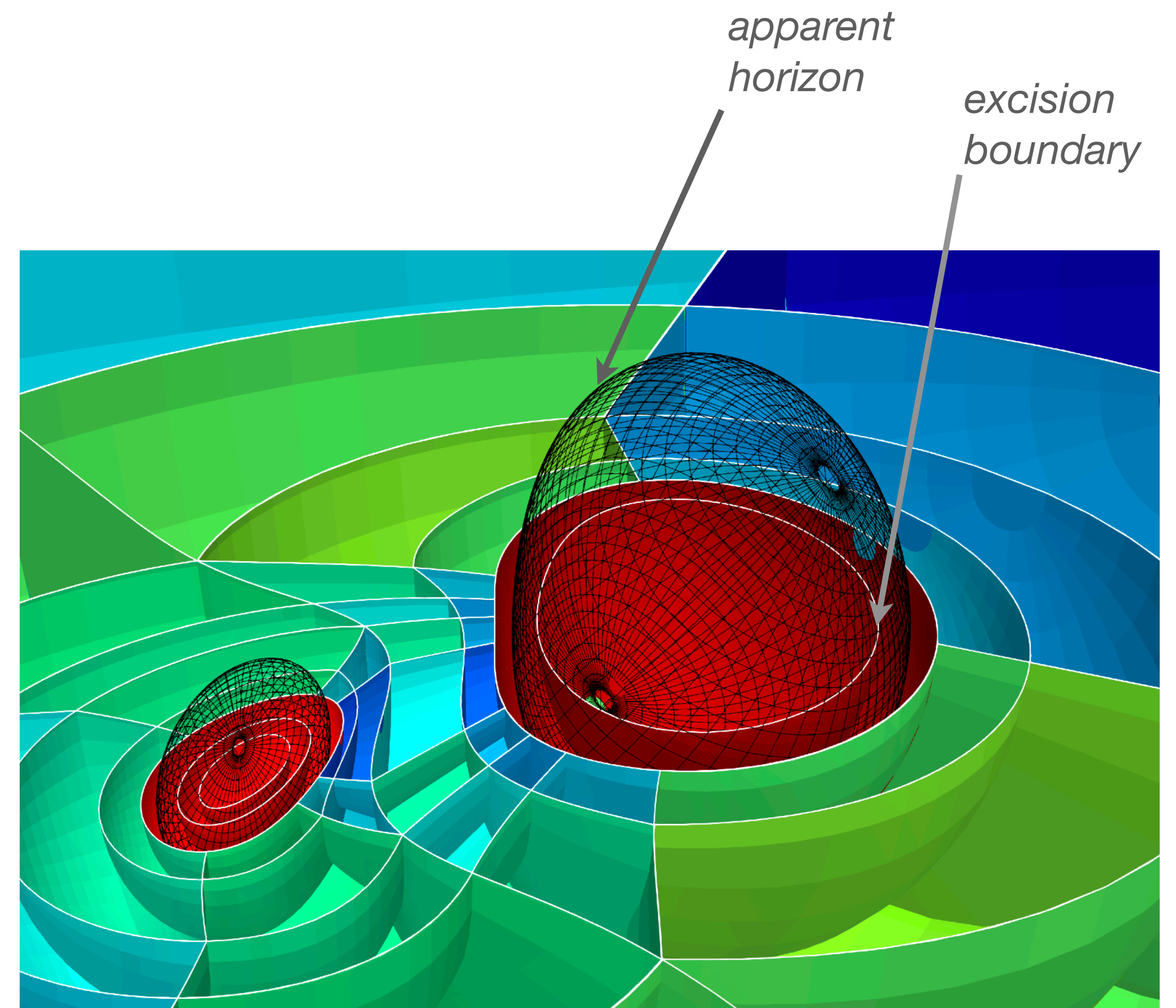
*Scheel, HP+ 08, Szilagyi+ 08,
Hemberger+ 13*



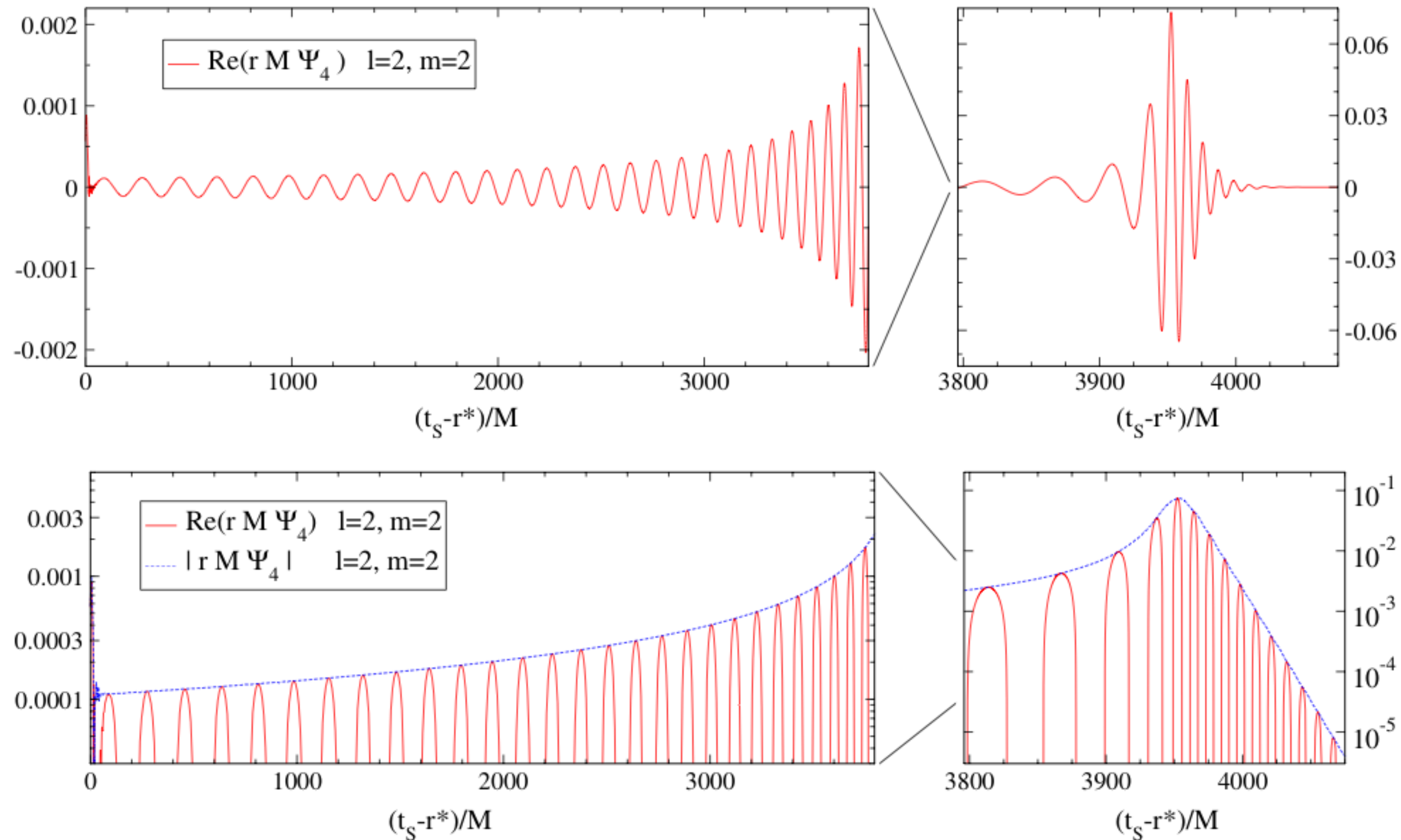
- Excise inside BH horizons
- Excision boundaries:
 - **follow BHs continuously**
 - conform to shape of AH
- **Horizon tracking & shape-control**

$$\vec{x}_{\text{inertial}} = a(t)\mathbf{R}(t)\vec{\xi}_{\text{grid}} + \text{deformations}$$

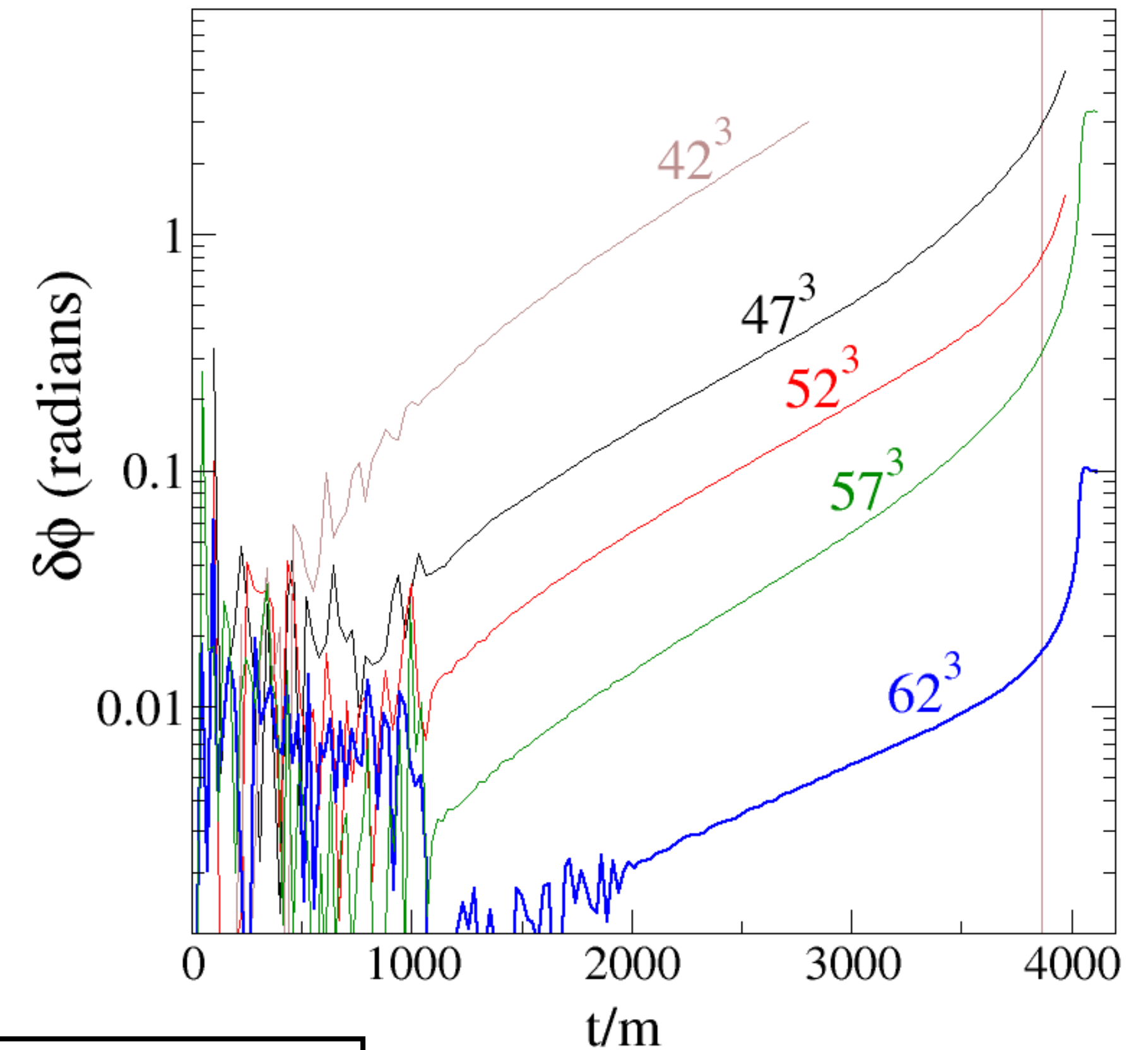
- $\{a(t), \mathbf{R}(t), \dots\}$ determined by feedback-loop
 - find AH(t) in $\vec{x}_{\text{inertial}}$
 - adjust $\{a(t), \mathbf{R}(t), \dots\}$ to keep excision boundaries inside AH



Accuracy of SpEC (circular inspiral)

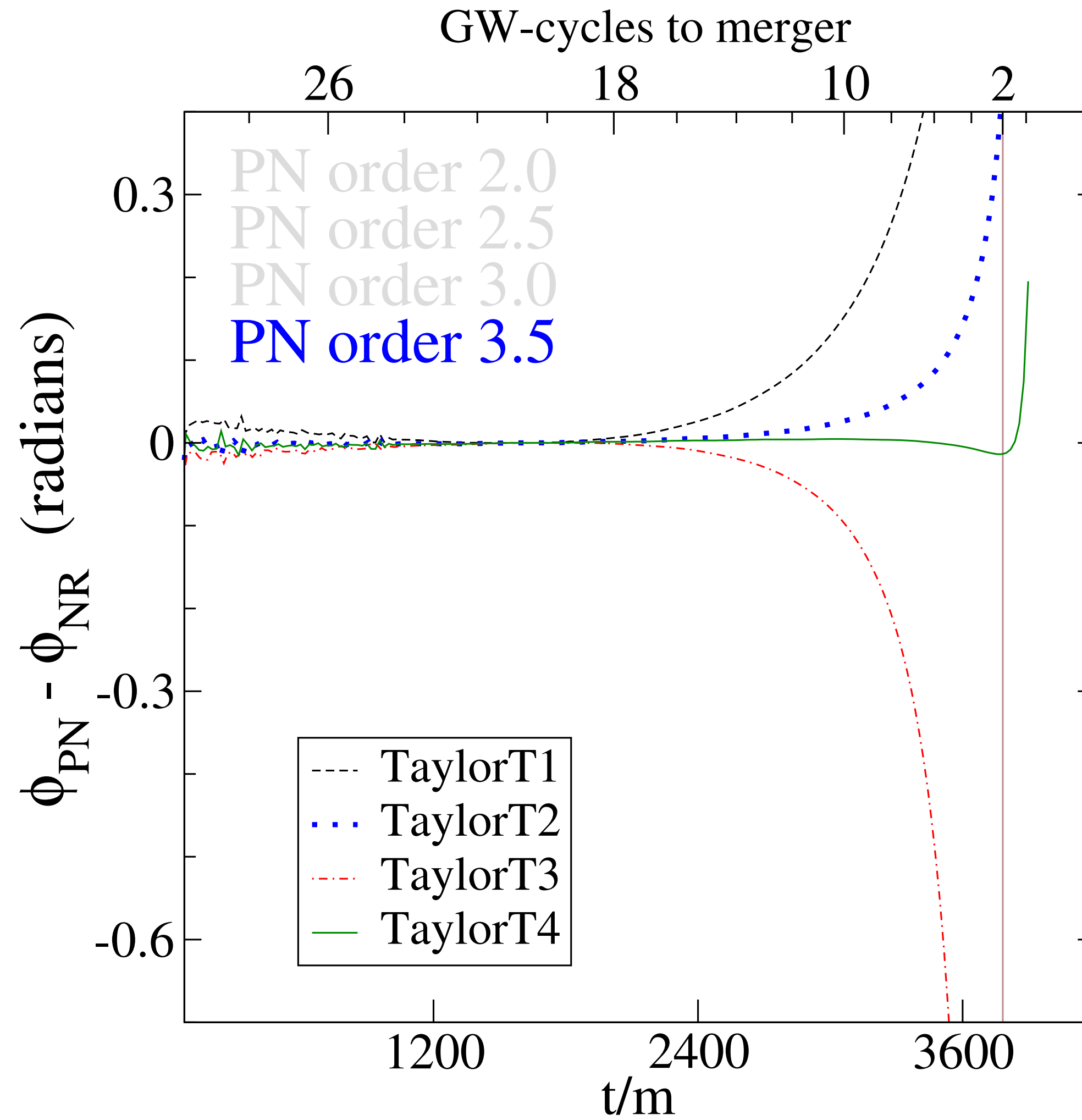
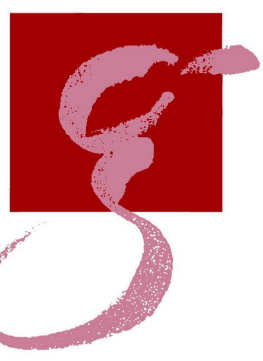


Phase error at different resolutions



- **Rapid convergence** due to spectral methods
- **Small errors** due to moving grid
- **Best code for long inspirals** (but mergers hard)

post-Newtonian vs. NR



TaylorT1...T4
Different choices to truncate
energy balance equation

$$\frac{dE}{dt} = -F_{\text{GW}}$$

Boyle..HP+ 07

NR Parameter space exploration

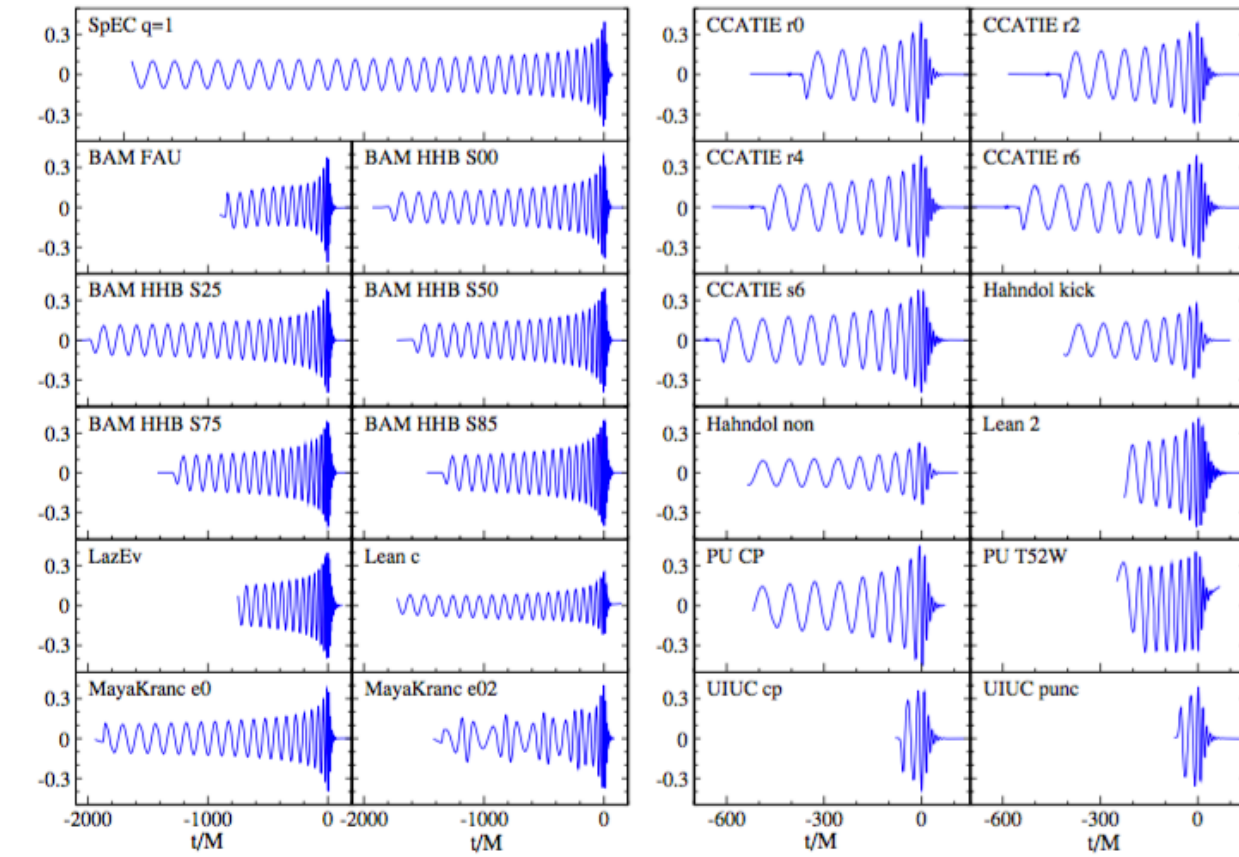
2nd SXS Catalog



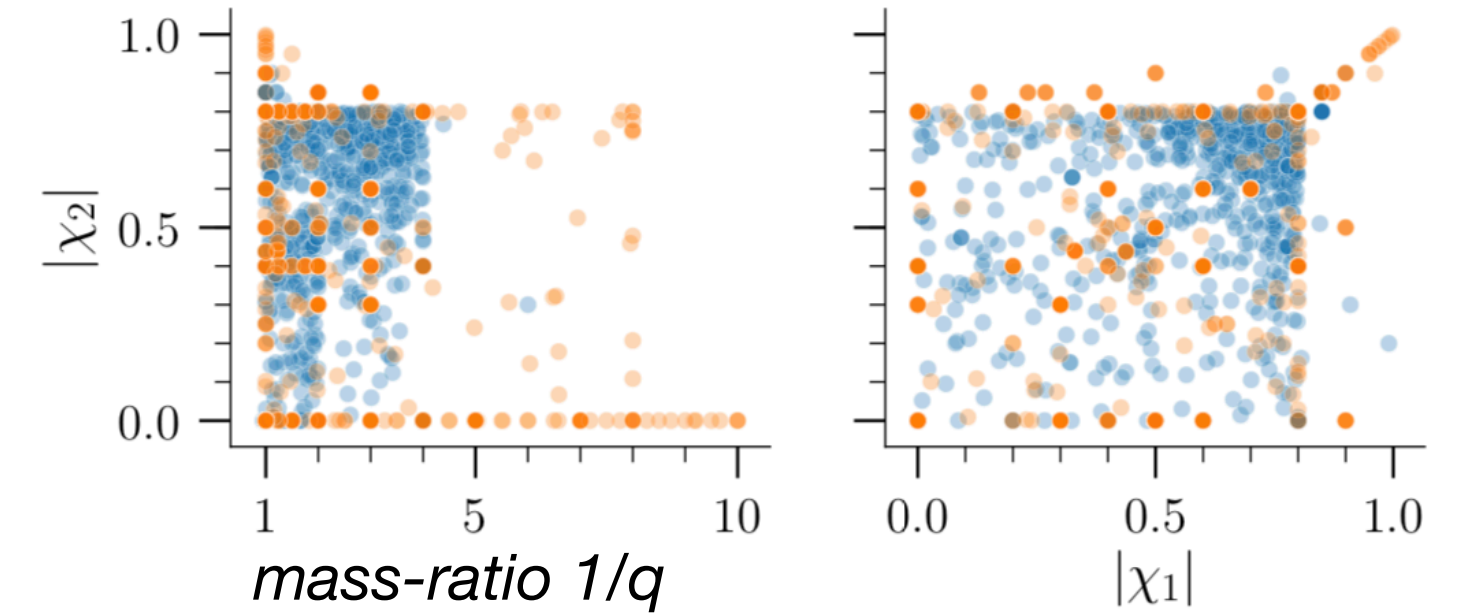
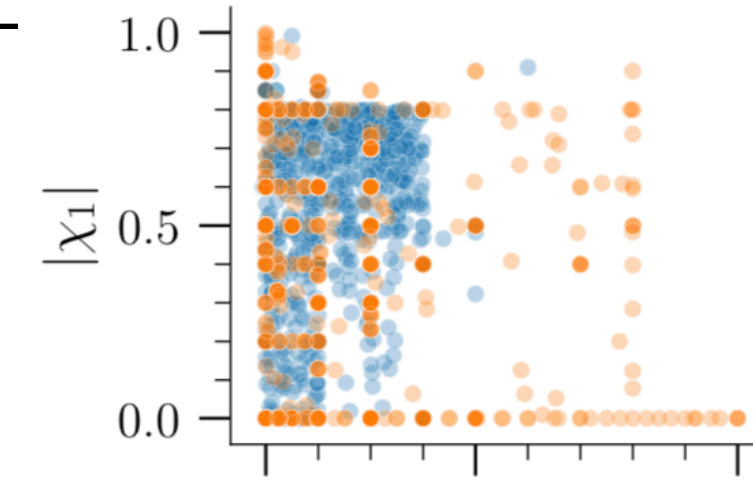
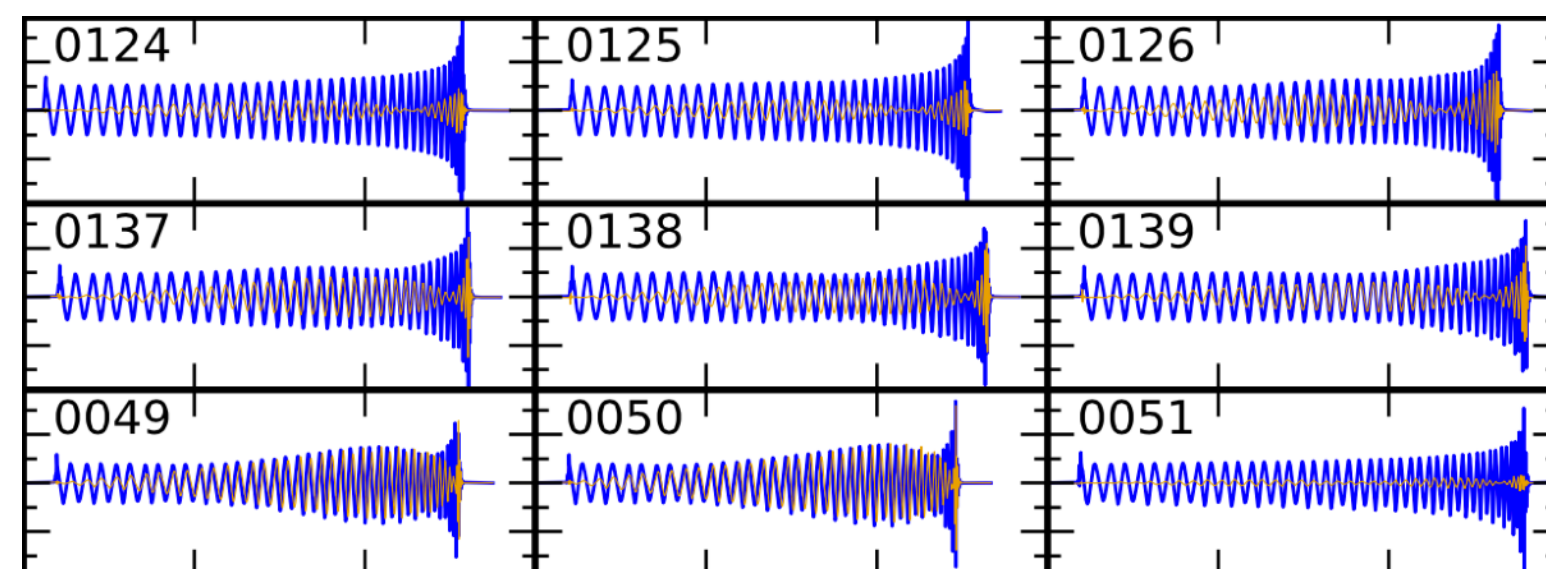
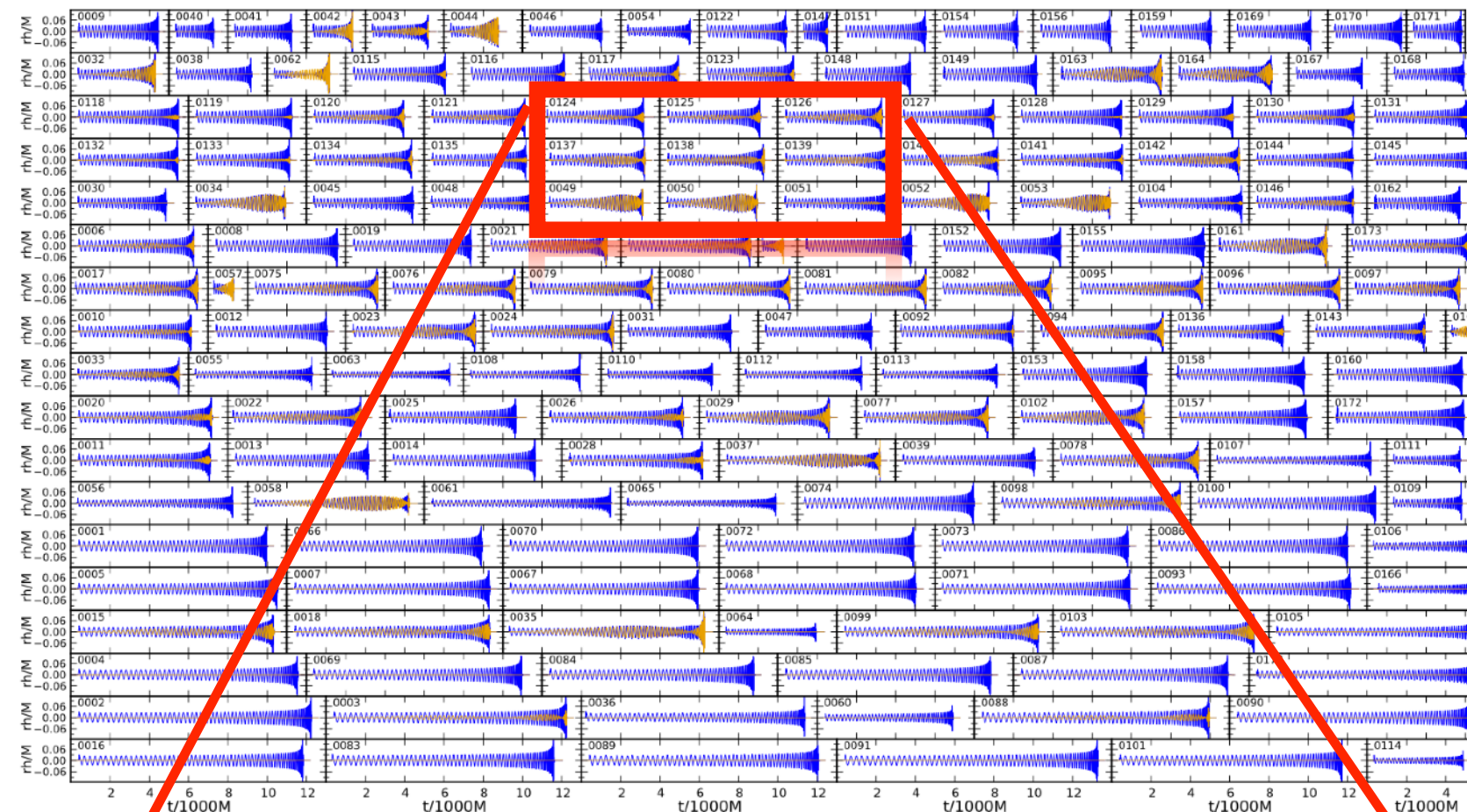
Boyle..HP+ 19

NINJA

Aylott .. HP+ 09

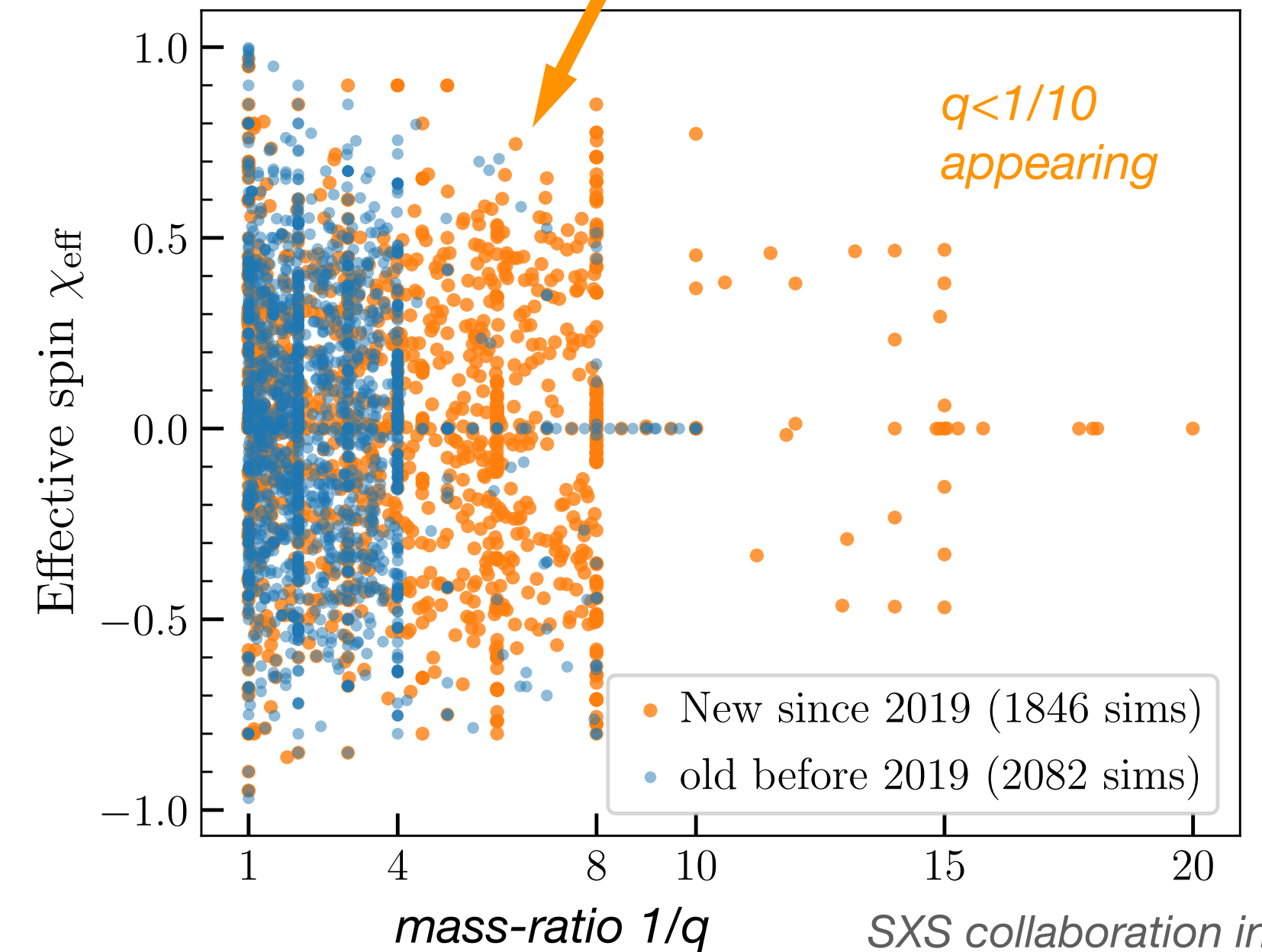


1st SXS Catalog Mroue .. HP+ 13



SXS now

$1/4 > q \geq 1/8$ populated



$q < 1/10$ appearing

- New since 2019 (1846 sims)
- old before 2019 (2082 sims)

SXS collaboration in prep

More parameter space exploration efforts



$q \geq 1/15$

Catalog	Started	Updating?	Simulations	m_1/m_2 range	$ \chi_1 $ range	$ \chi_2 $ range	Precessing?	Median N_{cyc}	Public?	
NINJA [98, 115]	2008	✗	63	1–10	0–0.95	0–0.95	✗	15	✗	
NRAR [120]	2013	✗	25	1–10	0–0.8	0–0.6	✓	24	✗	
Georgia Tech [122]	2016	✓	452	1–15	0–0.8	0–0.8	✓	4	✓	
RIT (2017) [123]	2017	✓	126	1–6	0–0.85	0–0.85	✓	16	✓	
RIT (2020) [124]	2017	✓	777	1–15	0–0.95	0–0.95	✓	19	✓	
NCSA (2019) [125]	2019	✗	89	1–10	0	0	✗	20	✗	<i>eccentric</i>
SXS (2018)	2013	✓	337	1–10	0–0.995	0–0.995	✓	23	✓	
SXS (2019)	2013	✓	2018	1–10	0–0.998	0–0.998	✓	39	✓	<i>longest sims</i>

highest spins

Table from Boyle et al 2019
(1904.04831)

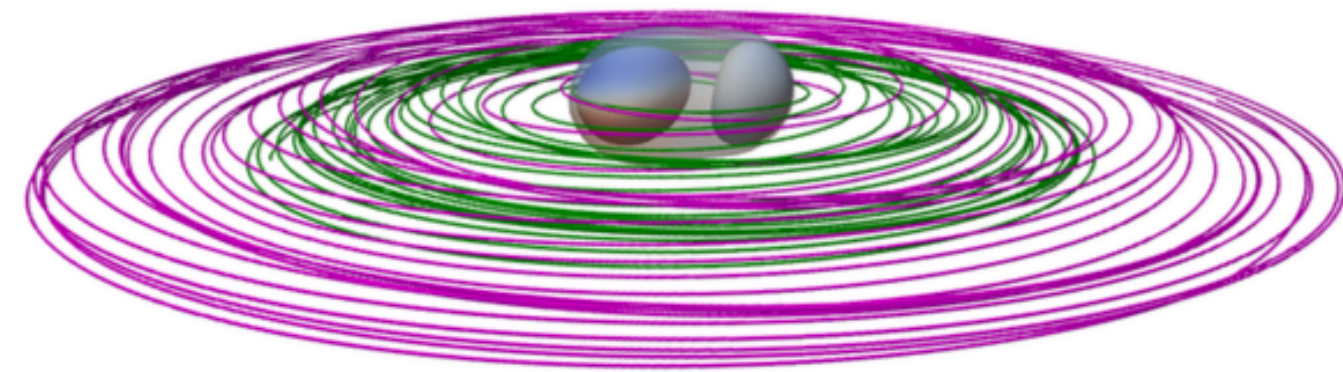
AND...

Palma group (Husa+), Cardiff group (Hannam+)

Eccentric catalog (Healy+Lousto (PRD 2022), 2202.00018)

Maya 2nd catalog (Ferguson+ 2309.00262)

Parameter space: NR records

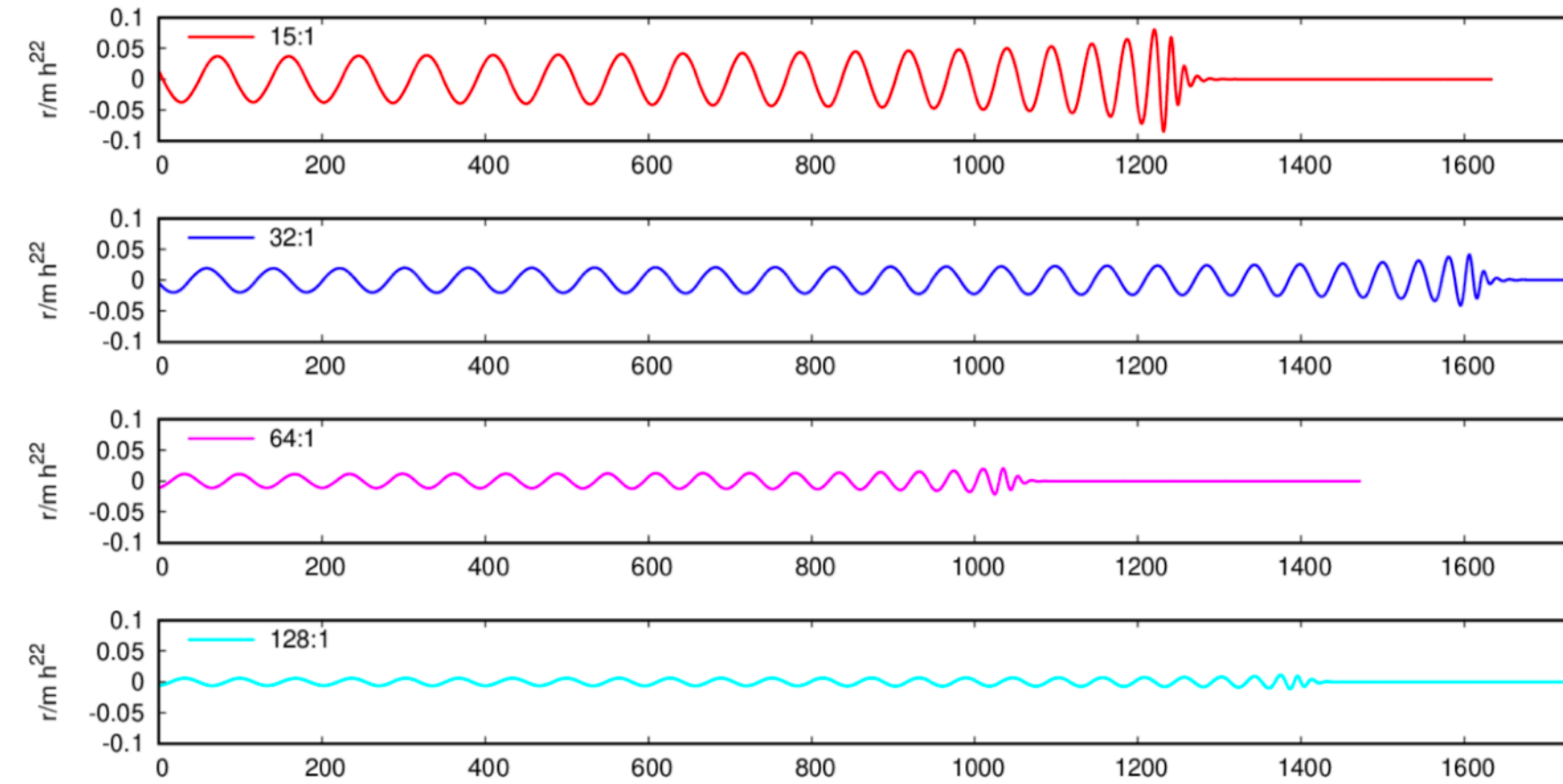


$q=1: S/M^2 = 0.998$
 $q<1: S/M^2 = 0.95$

Scheel+ 14
 Lovelace..HP+15
 Boyle..HP+ 19

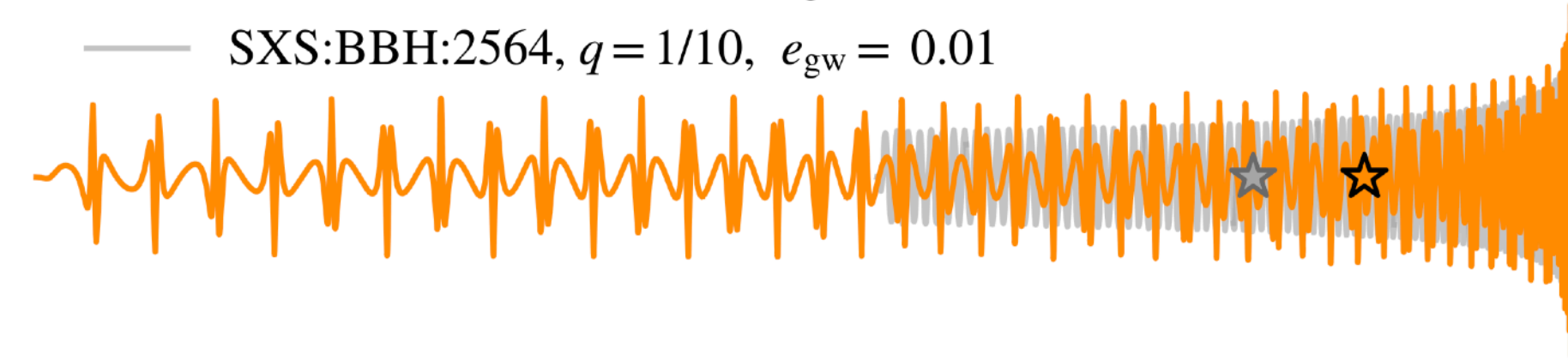
$q = 1/128$

Lousto, Healy 20



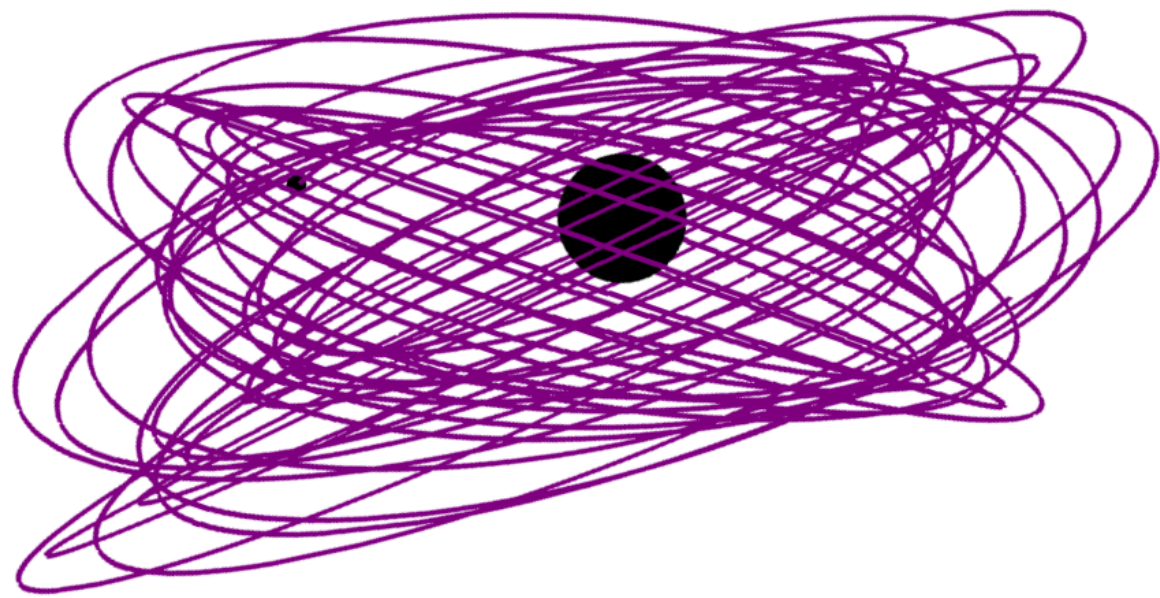
eccentric, high q

— SXS:BBH:2567, $q=1/10, e_{\text{gw}} = 0.34$
 — SXS:BBH:2564, $q=1/10, e_{\text{gw}} = 0.01$



Ramos-Buades, .. HP+ 22

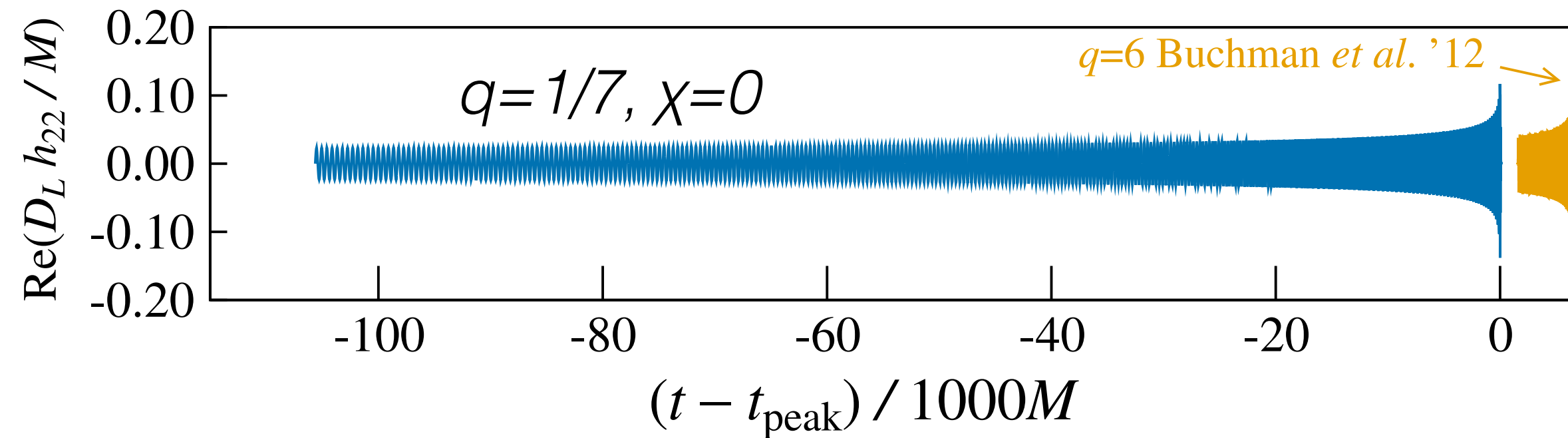
eccentric, precessing $q = 1/7$



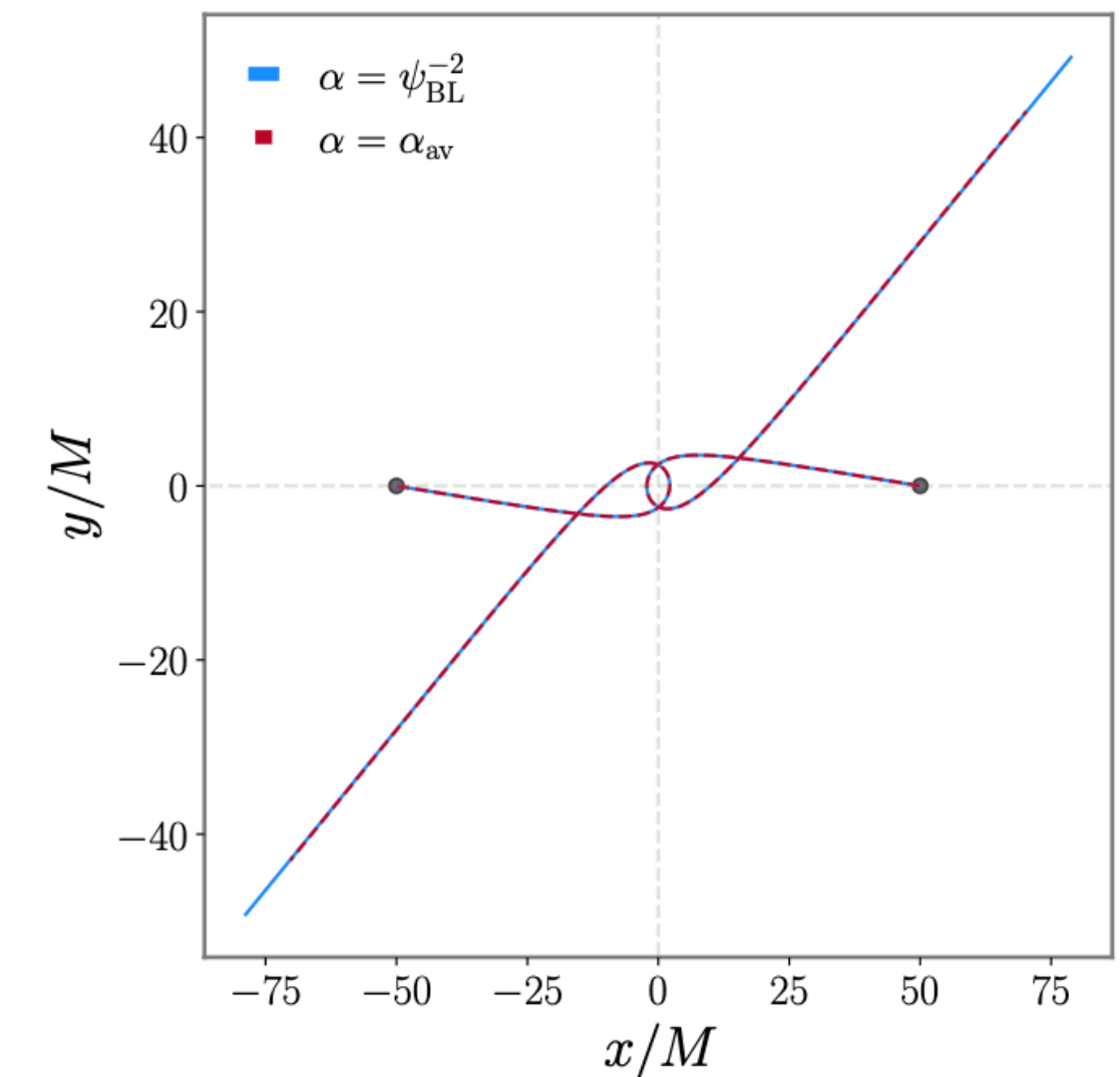
Lewis, Zimmerman, HP 17

350 GW cycles

Szilagyi..HP+15



hyperbolic scatter w/ spin



Retegno, Pratten+ 23

$M = 50 M_{\odot}$

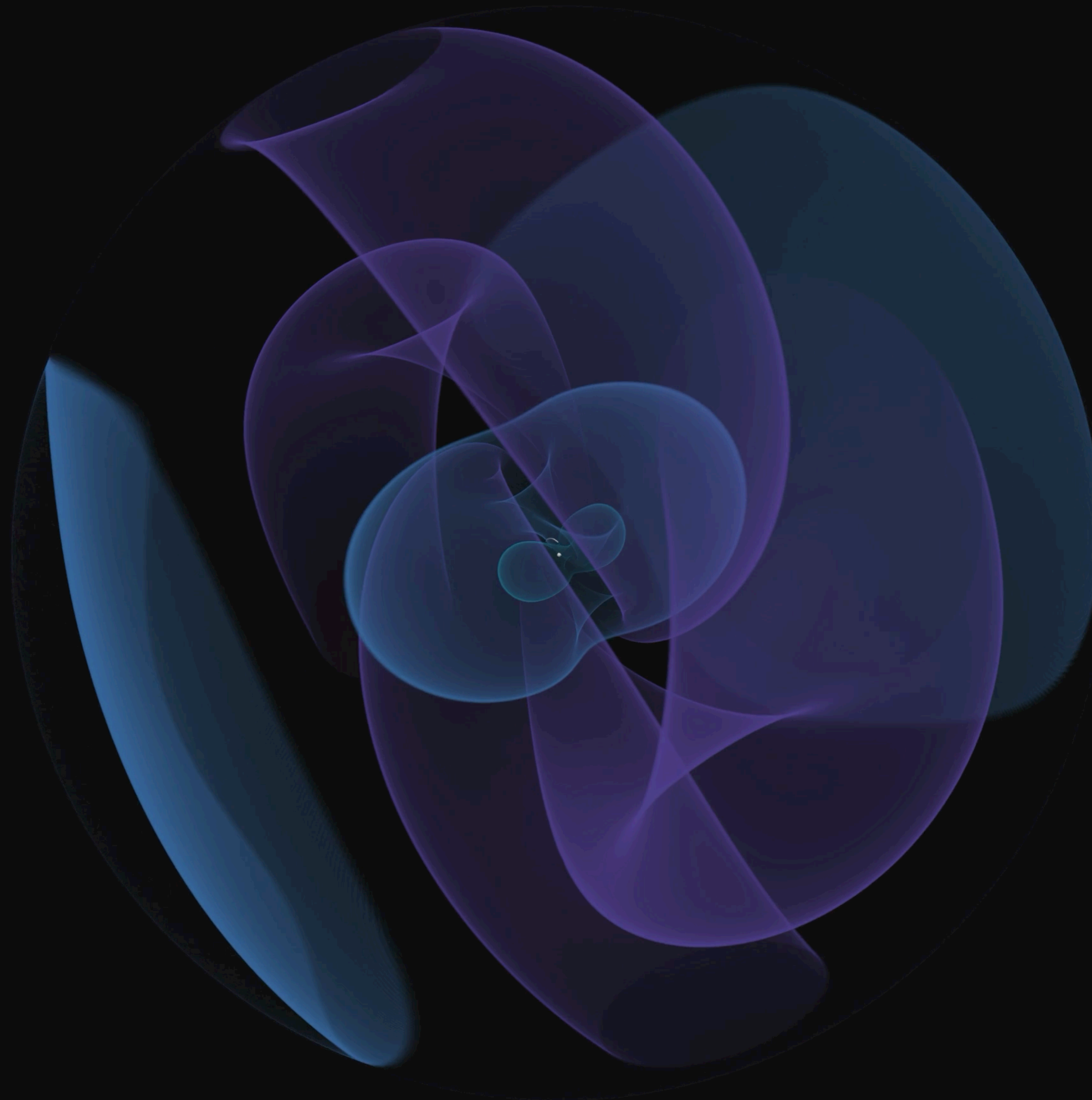
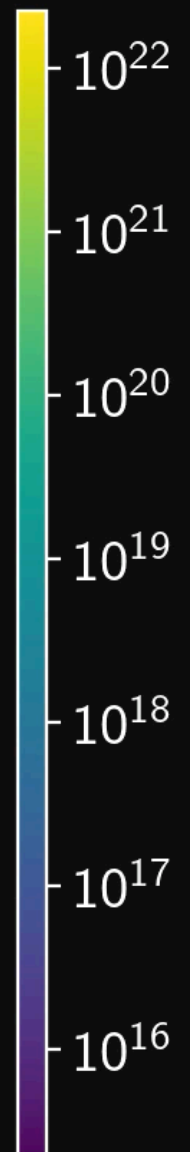
$q = 1 : 6$

$\vec{\chi}_1 = (0.17, -0.38, 0.68)$

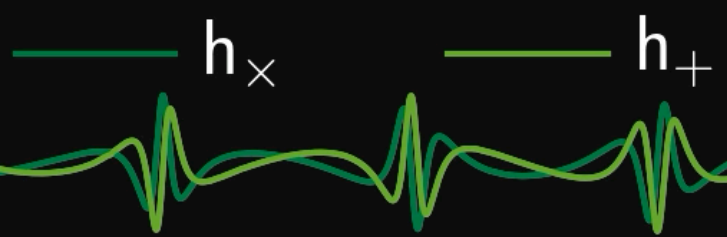
$\vec{\chi}_2 = (0.00, 0.00, 0.00)$

$e = 0.40$

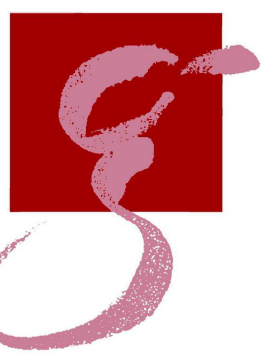
$F [L_{\odot} \text{sr}^{-1}]$



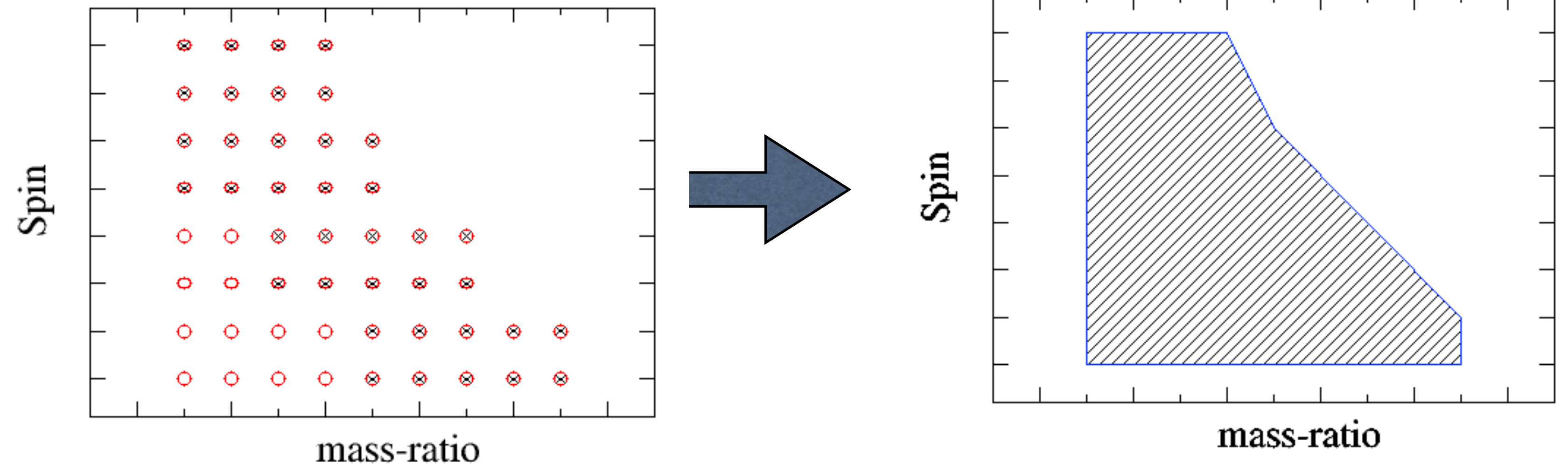
-1340 ms



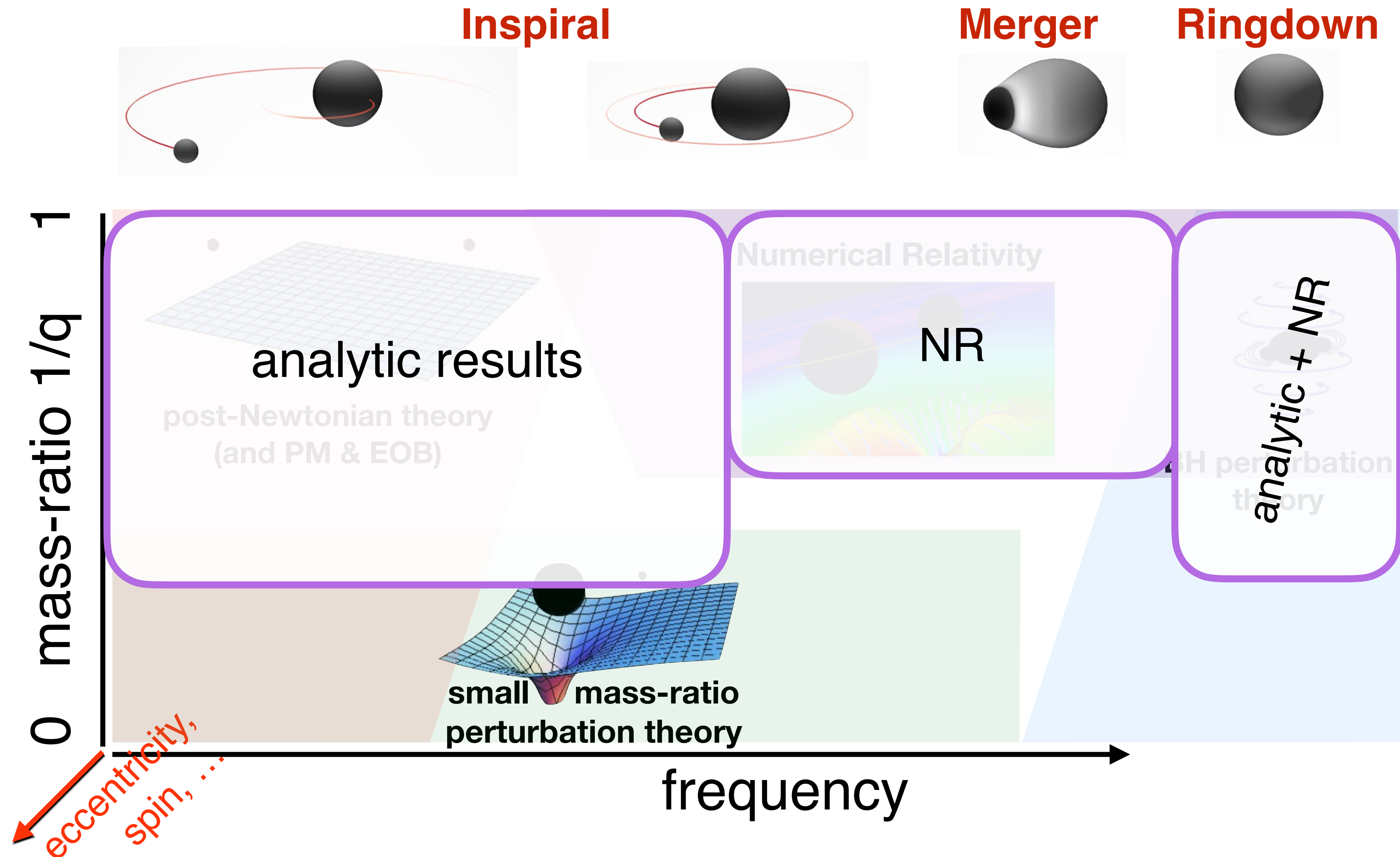
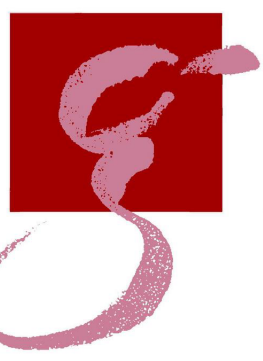
Waveform models



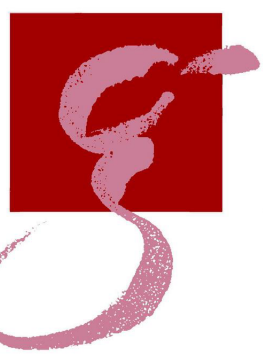
- **Continuous** in parameters $\theta = \{m_1, m_2, \vec{S}_1, \vec{S}_2; e, l; \iota, \phi, \psi; \text{RA}, \text{dec}, D_L, T_c\}$
- **Fast** evaluation
- Cover parameter-space
- Accurate



LIGO/Virgo IMR waveform models



Inspiral-merger-ringdown BH-BH waveform models



Effective one body (EOB)

$$H = \mu \sqrt{p_r^2 + A(r) \left[1 + \frac{p_r^2}{r^2} + 2(4 - 3\nu)\nu \frac{p_r^4}{r^2} \right]}$$

Hamiltonian dynamics
dynamics $\Rightarrow h_{lm}(\theta; t)$

Buonanno, Damour 99

EOBNR Buonanno+ 09, Pan+11

SEOBNR (aligned spins) Taracchini+ 12

SEOBNv3 (precessing) Pan+ 14, Taracchini+14,

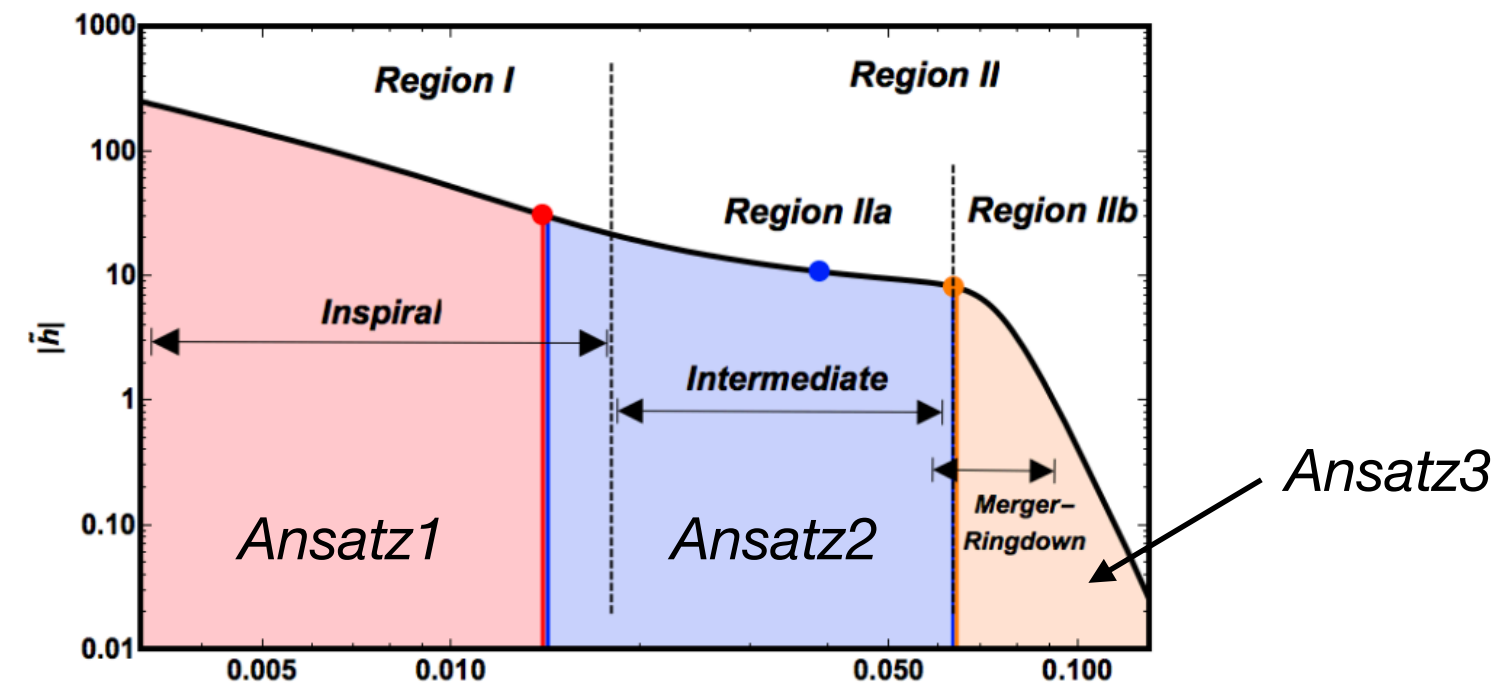
SEOBNRv4 {HM, P, PHM, HM_ROM} Bohe+ 17, Cotesta+ 18,

Cotesta+ 20, Ossokine+ 20

TEOBResumS, TEOBResumSM Nagar+ 18, 20

SEOBNRv5{HM,P,PHM,E,PEHM}

Phenomenological (Phenom)



Power-series for
 $\tilde{h}_{lm}(\theta; f)$

Phenom Ajith+ 08

Phenom{B,C} (aligned spins) Ajith+ 11, Santamaria+ 10

PhenomD Husa+ 15, Khan+ 15

PhenomHM London+ 18

Phenom{P, Pv2} (precessing) Hannam+ 13

Phenom{Pv3, Pv3HM} Khan+ 19, Khan+ 20

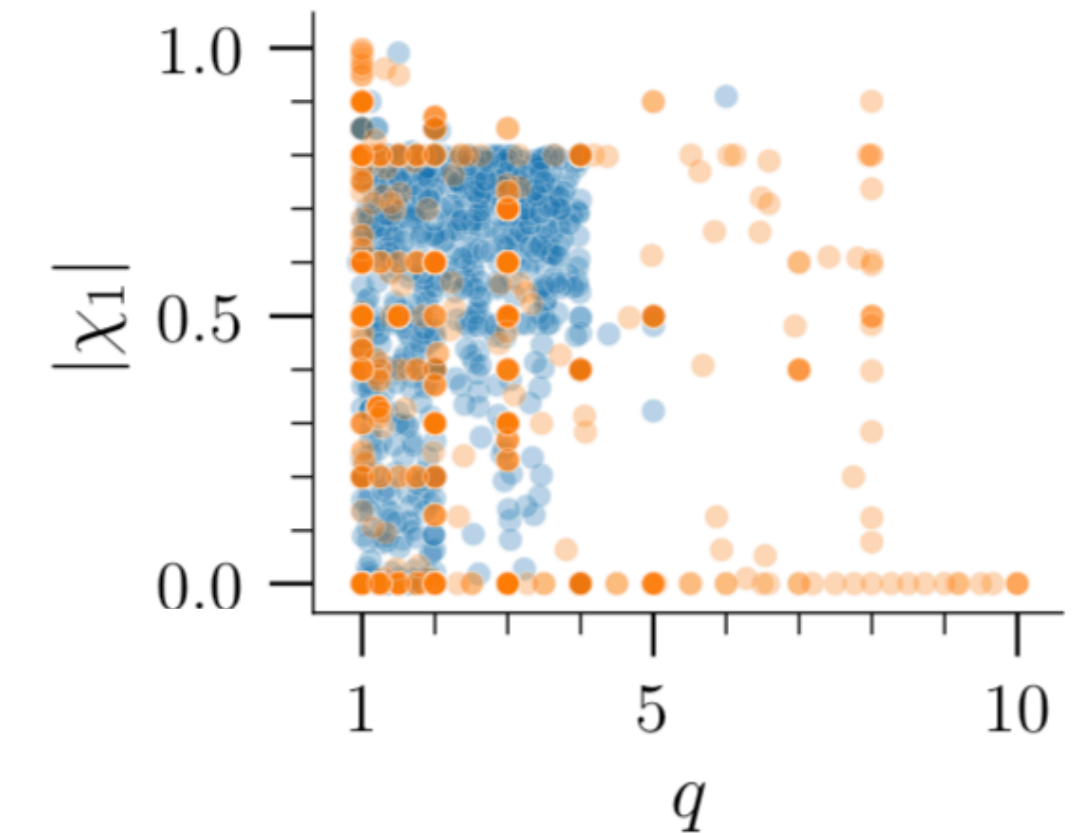
PhenomX{AS, HM, P, PHM} Pratten+ 20,

Garcia-Quiros+ 20ab, Pratten+ 21

PhenomPNR Hamilton+ 21

PhenomX_{O4a, Taylor}

NR surrogate Models



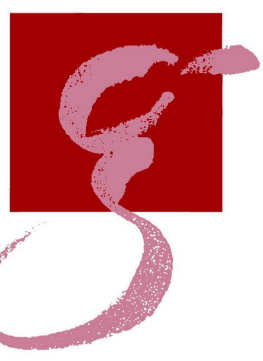
Direct interpolation
in parameters θ

Blackman+ 15,17,18

Varma+ 18,19

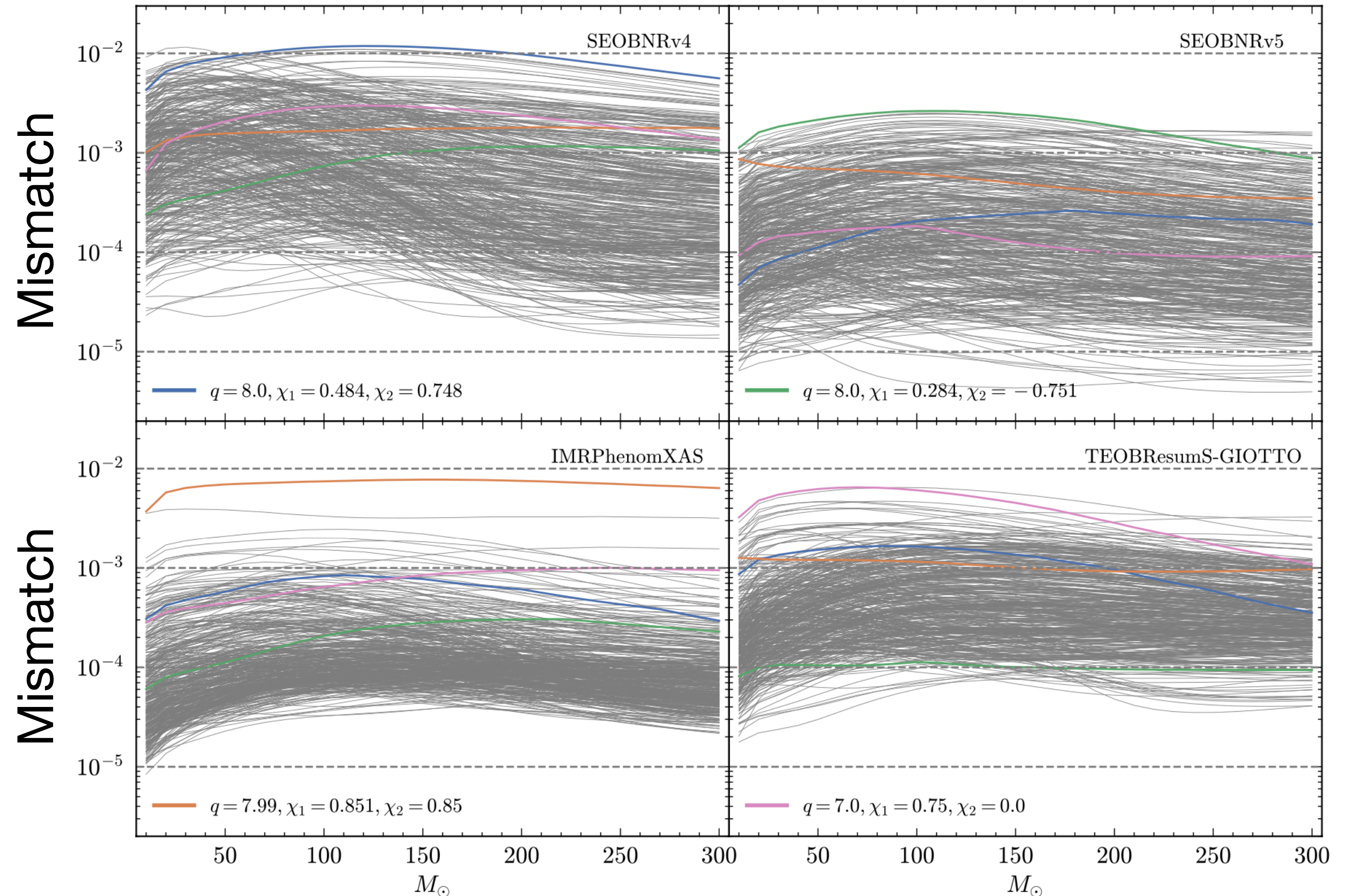
Islam+ 21

Major application of NR: waveform modeling



- **Assess accuracy** of waveform models
- Determine **importance of improvements**
 - higher order
 - additional physics
 - higher modes, precession, memory
- **Calibrate parameters** in model to improve agreement
 - GW modes & fluxes
 - merger dynamics
 - ringdown attachment
 - ...

Errors in 4 waveform models as compared to 442 NR simulations

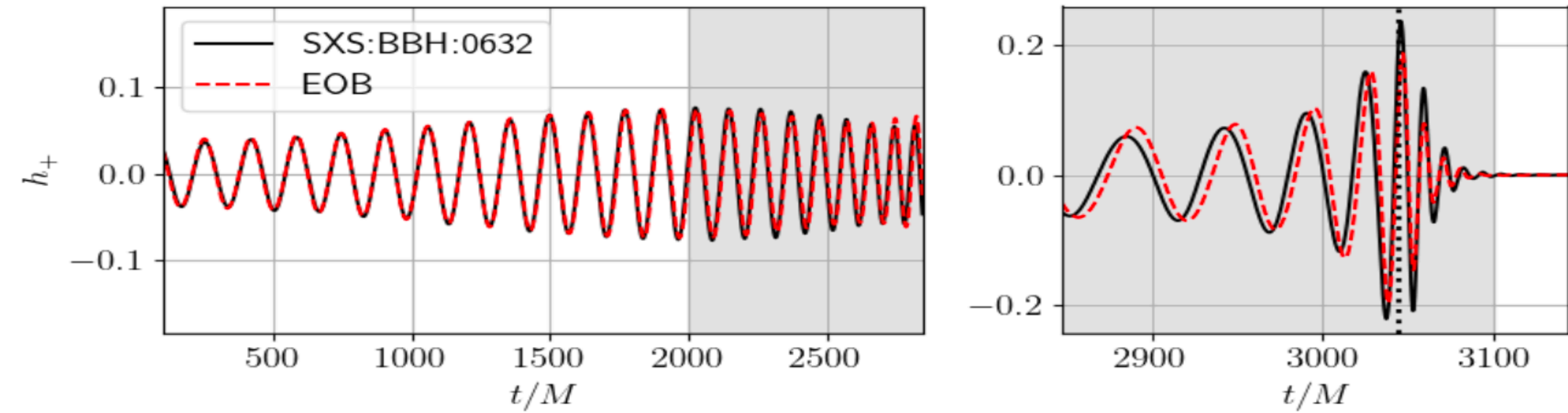


Example plots of waveform models



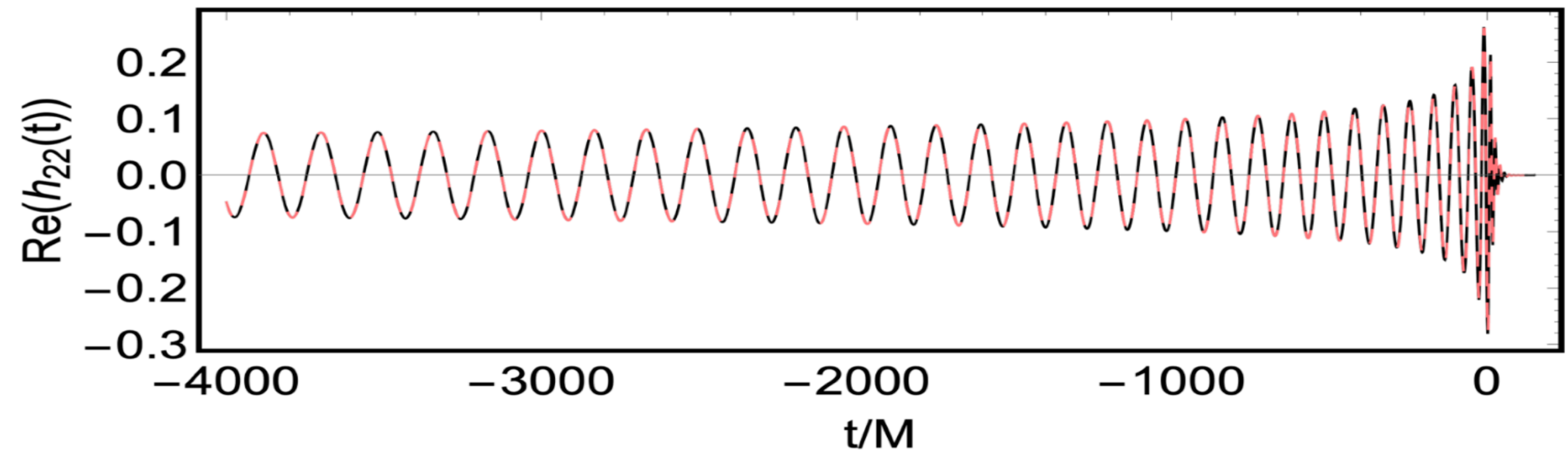
TEOBResumS

Plot from Gamba+ 22



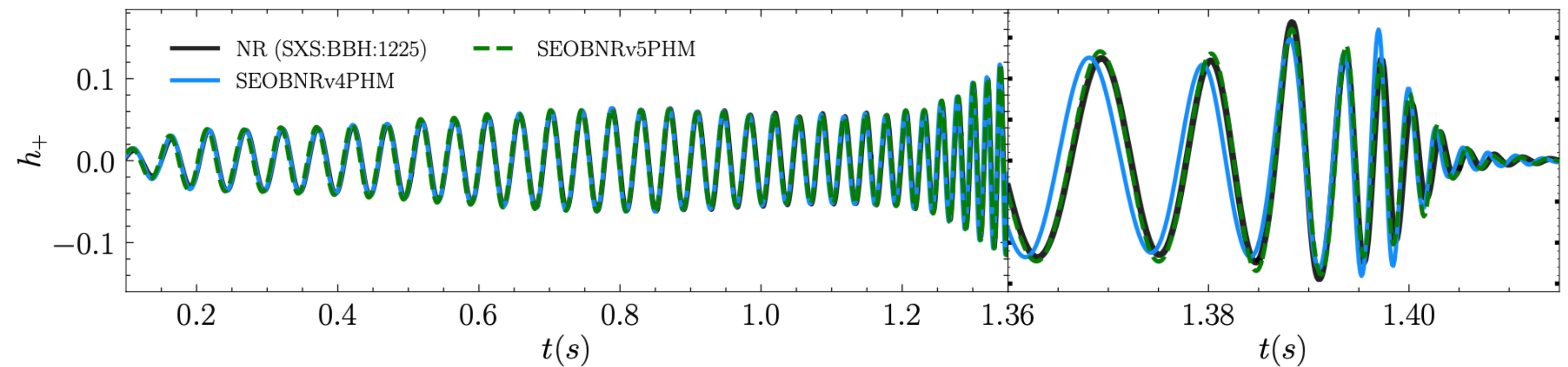
IMRPhenomT

Plot from Estelles+ 20



SEOBNR

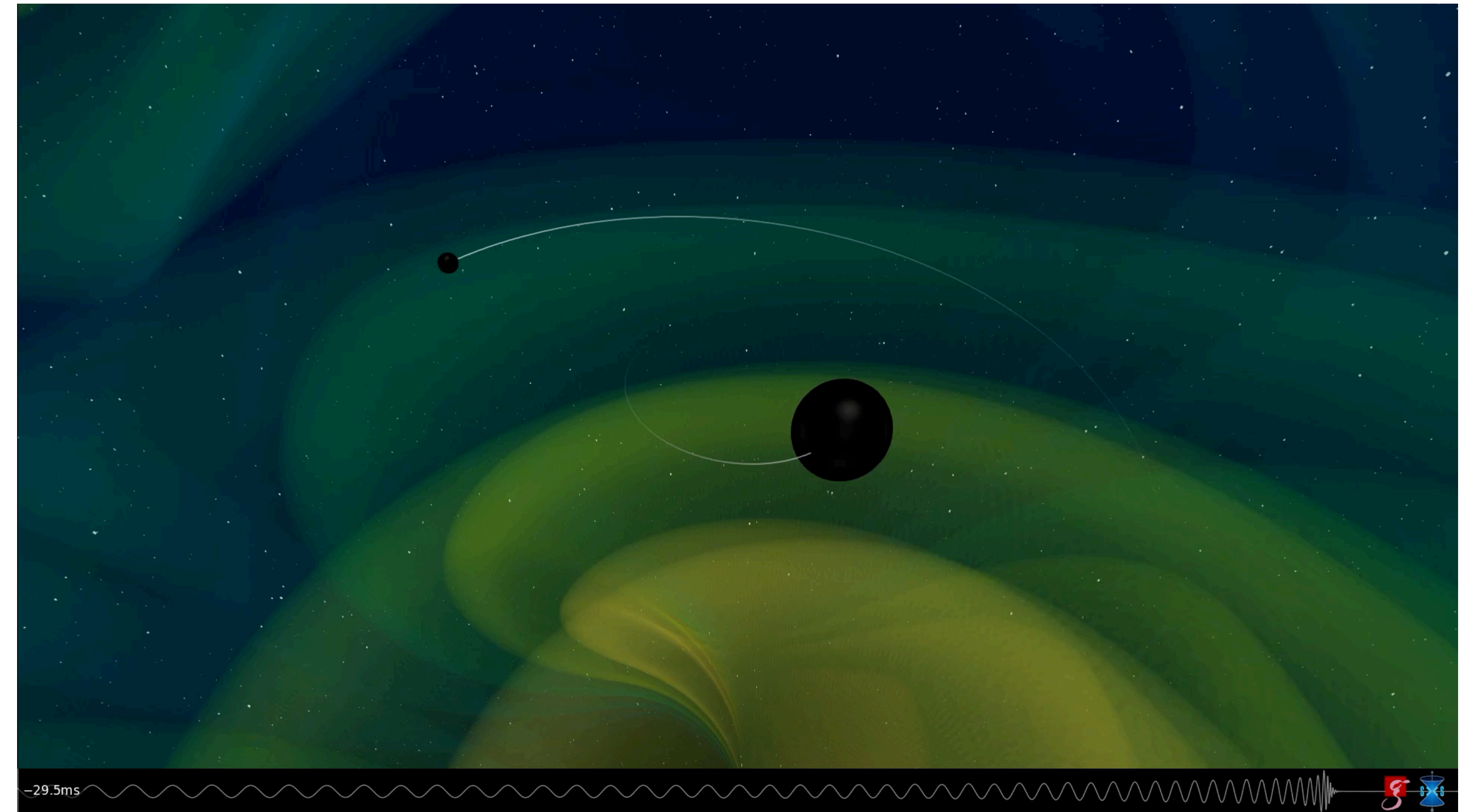
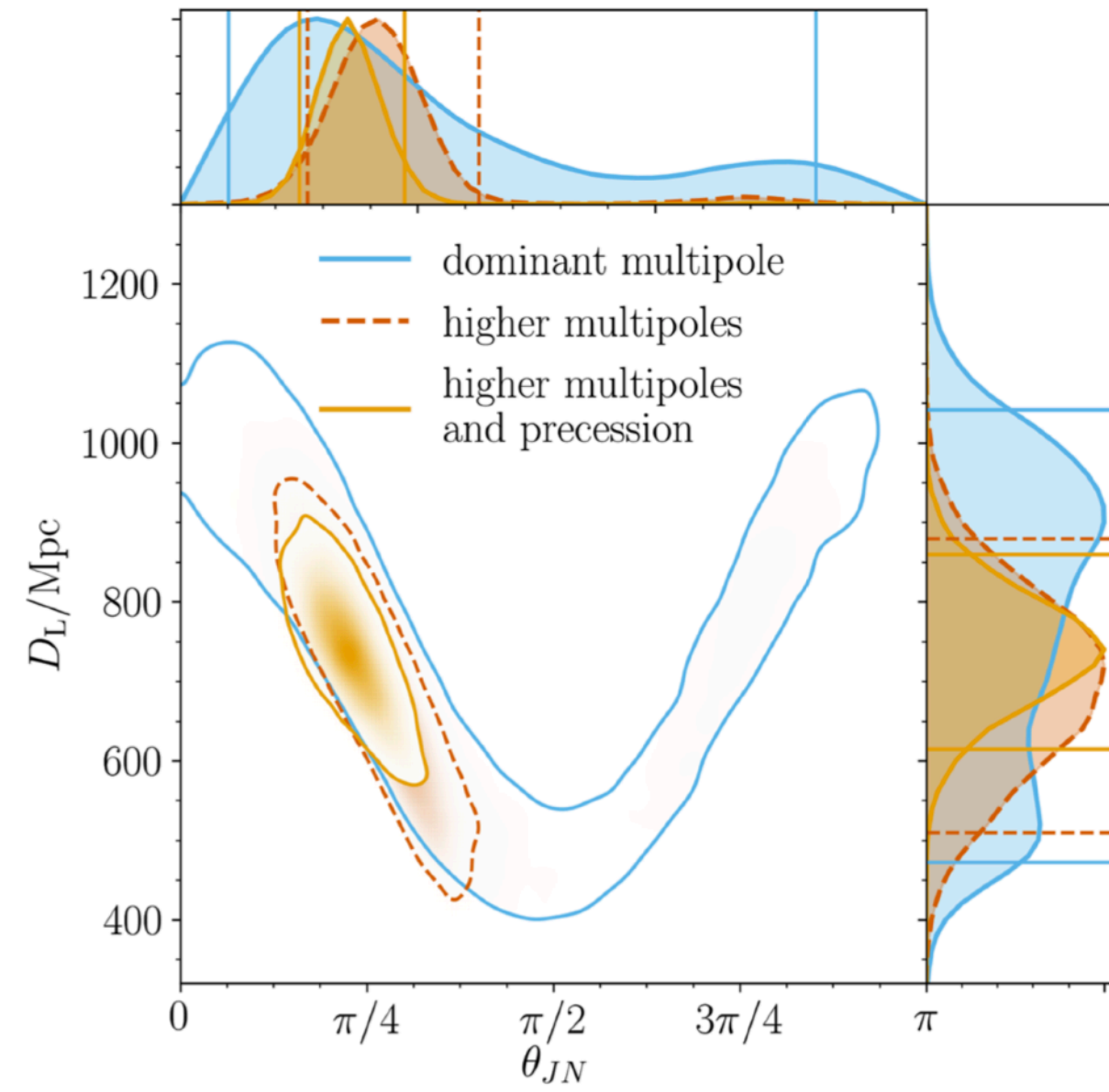
Plot from Ramos-Buades+ 23



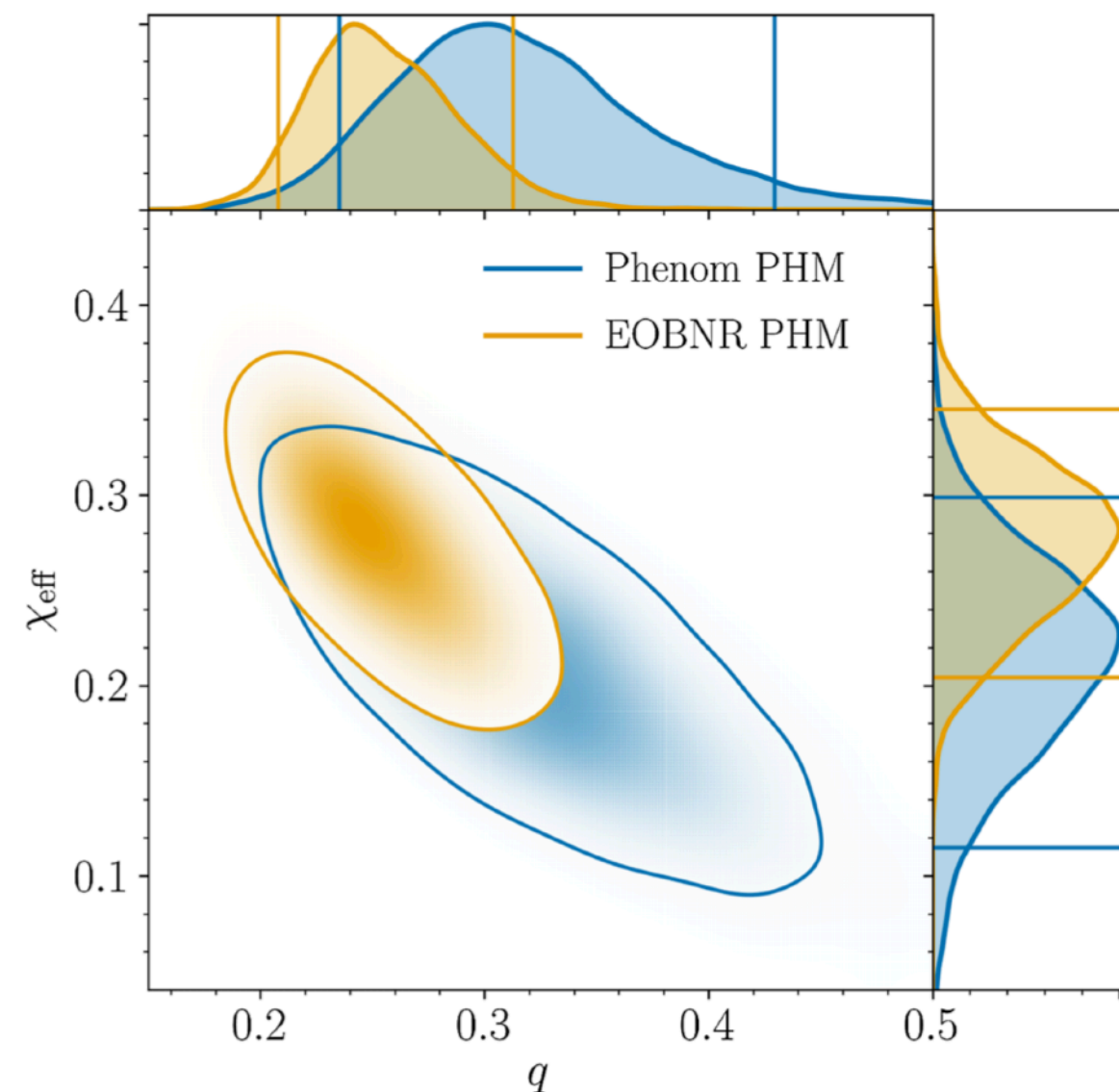
case-study: GW190412 $(30 + 8)M_{\odot}$ at SNR=19



higher-modes break degeneracies



Abbott et al,
PRD 102 043015 (2020)



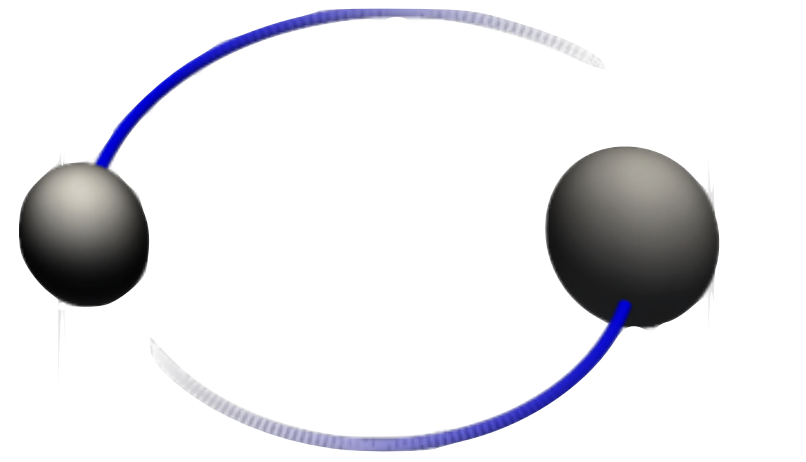
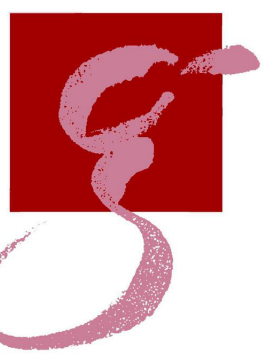
N. Vu, HP

some differences between waveform models remain even at current SNR

IMRPhenomPv3HM Khan et al 19; 20
SEOBNRv4PHM Ossokine+ 20

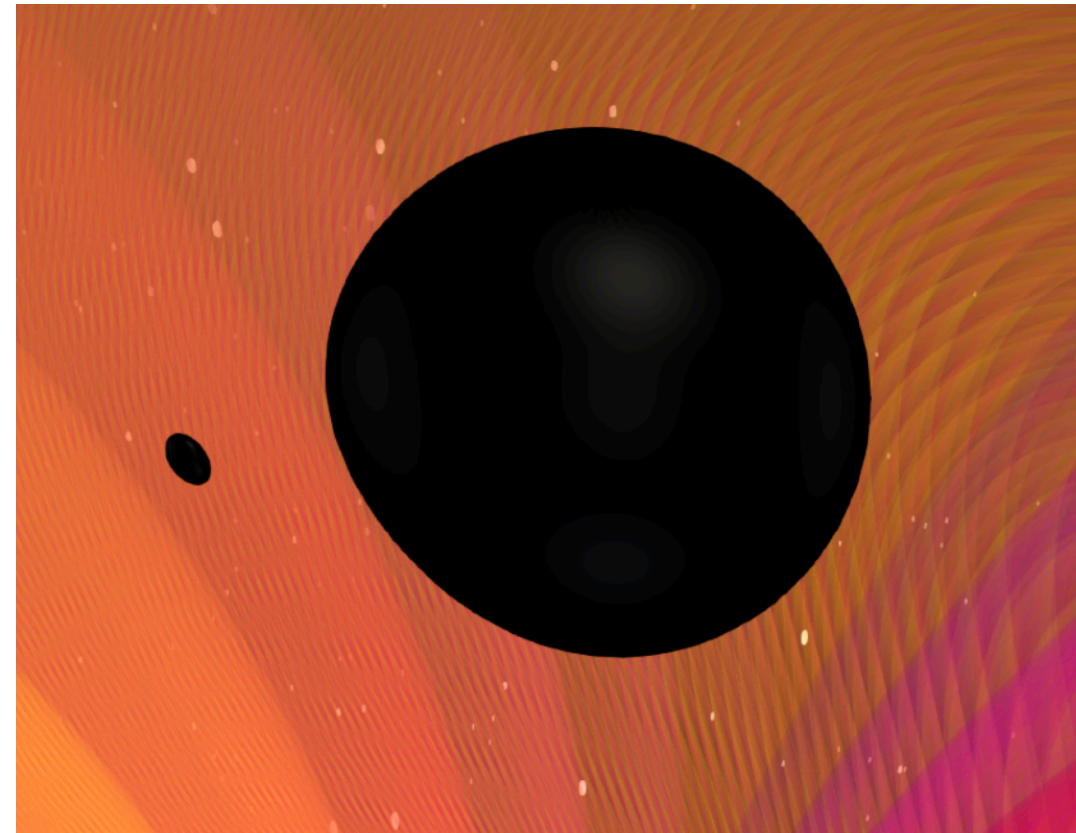
(some)
recent NR developments

Binaries at all mass-ratios



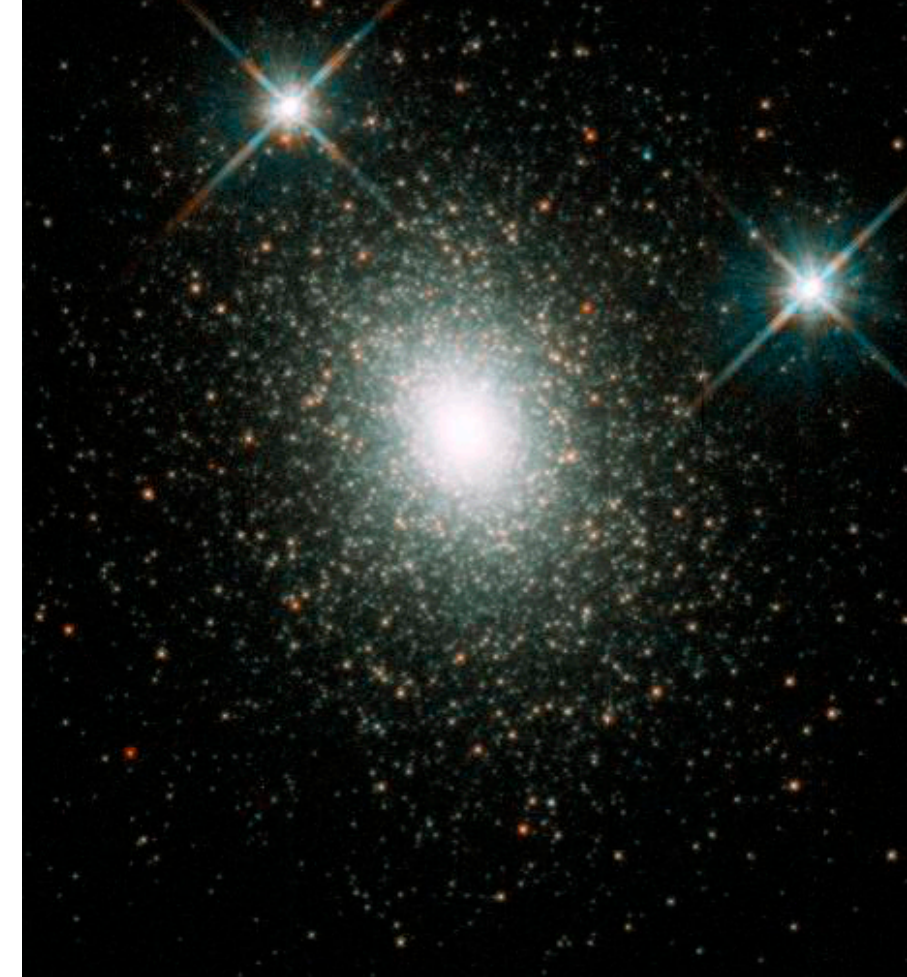
LVC

GW150914
 $q \sim 1$

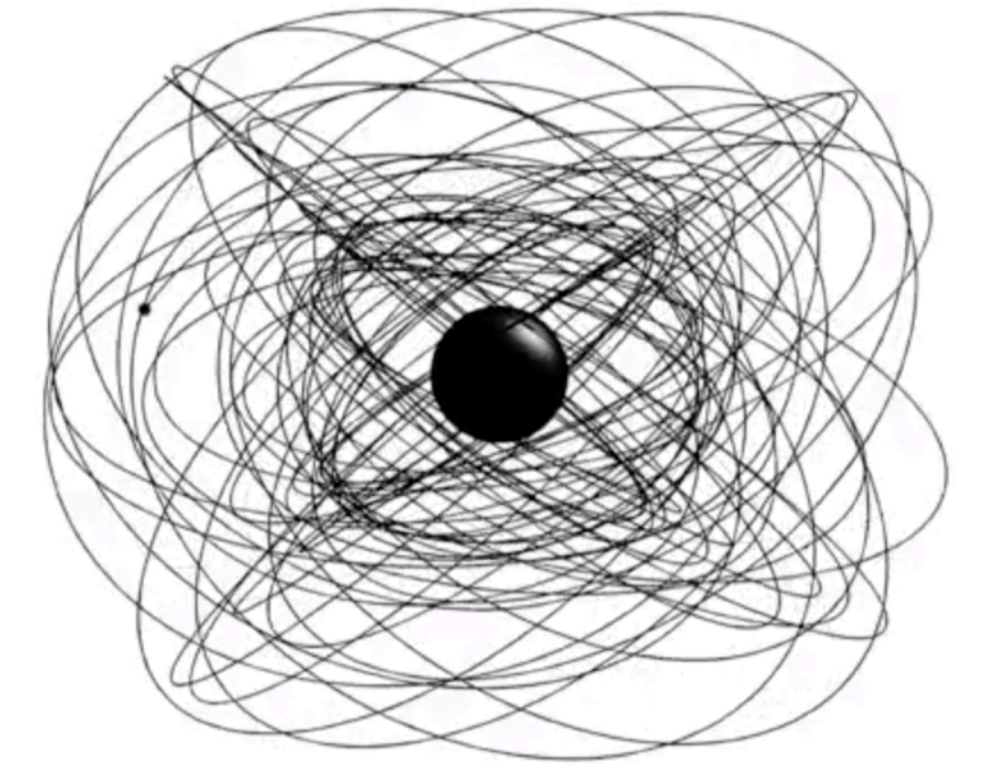


GW190814
 $q \sim 0.1$

Vu, HP



Intermediate mass BH NASA
 $(10 + 1000)M_{\odot}$
 $(10^3 + 10^6)M_{\odot}$



S. Drasco

EMRI
 $(10 + 10^6)M_{\odot}$

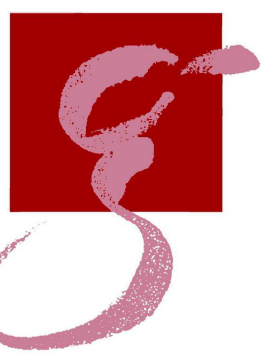


NR
 $q \gtrsim 1/20$

Small-mass-ratio approximation (GSF)

expansion in symmetric mass-ratio $\nu = q/(1+q)^2$

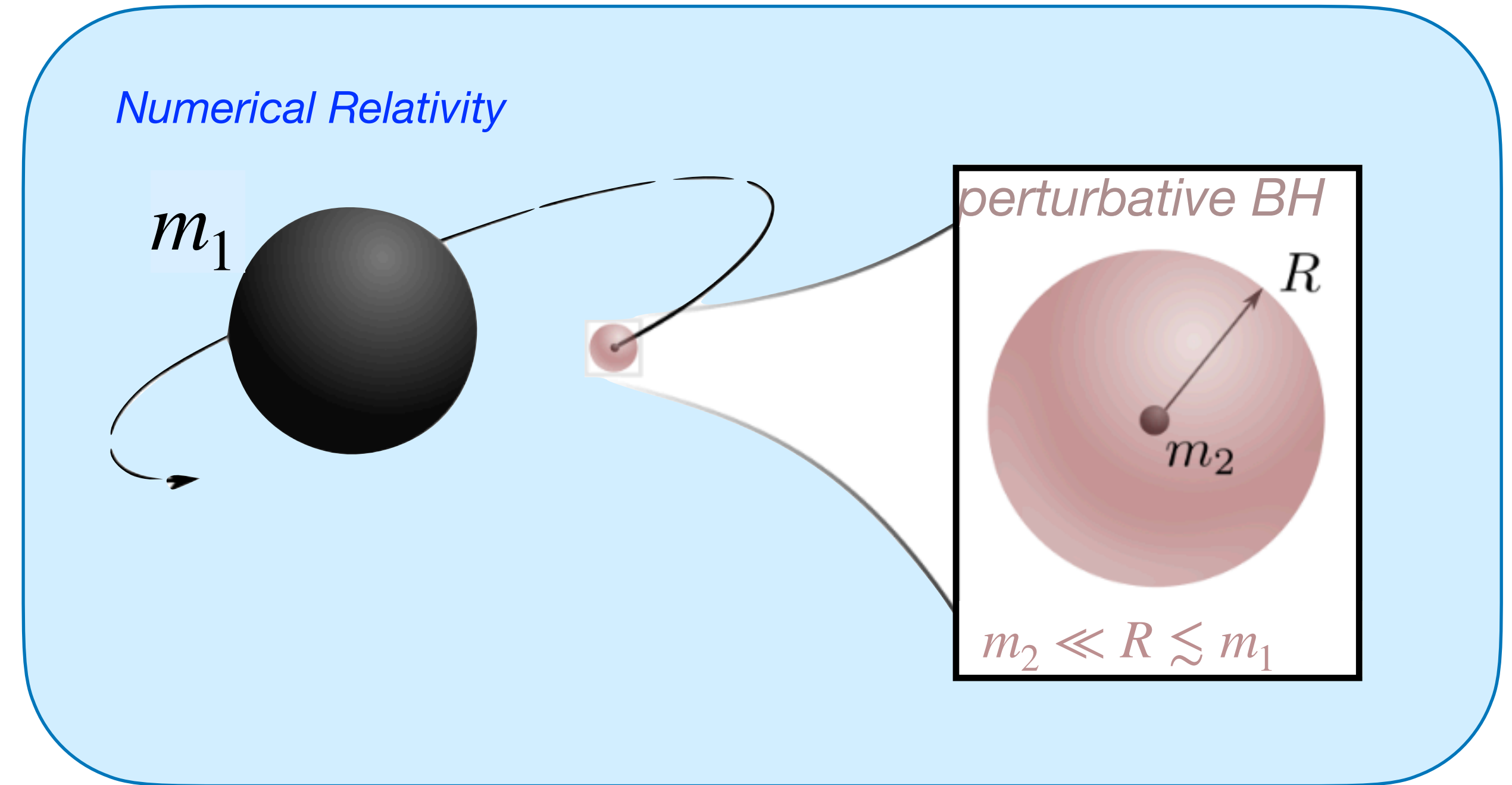
$$\text{Cost} \sim \frac{T}{\Delta t} \sim \frac{q^{-1}(M\Omega_i)^{-8/3}}{q} \sim q^{-2}(M\Omega_i)^{-8/3}$$



- **Worldtube excision**

- **Excise region** with radius R around m_2
- In excised region, use **tidally perturbed BH** metric
- Determine internal BH-perturbations by matching to NR
- Set NR boundary conditions from internal solution

- Courant limit **increased by factor**
 $R/m_2 \sim 1/q \gg 1$



Dhesi, Rüter, Pound, Barack, HP PRD 104 (2021) 124002
Wittek, Dhesi, Barack, HP, Pound+ PRD 108 (2023) 024041
Wittek, Pound, Barack, HP+ in prep



*Idea from
Bernard Schutz*

Warm up: Scalar charge



- Scalar charge orbiting Schwarzschild BH

$$g^{ab} \nabla_a \nabla_b \Phi(x^c) = -4\pi\rho(x^c)$$

- Goal: handle **point-charge**

$$\rho(x^c) = q \int \frac{\delta^4[x^c - x_p^c(\tau)]}{\sqrt{-g}} d\tau$$

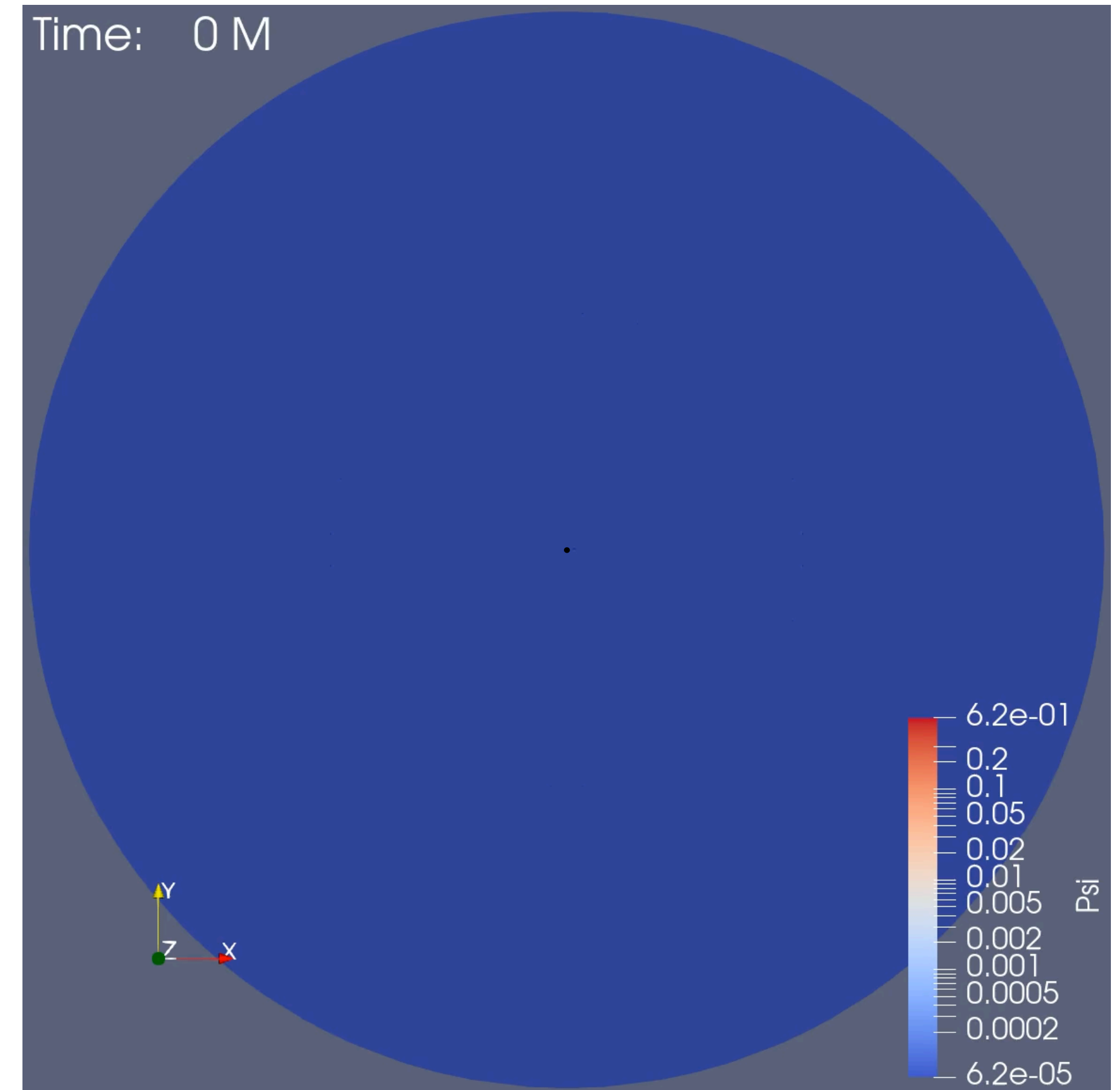
- Perturbative solution inside worldtube

- $\Phi^A = \Phi^S + \Phi^R$

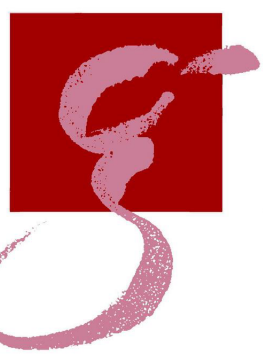
- $\Phi^S(x^i, t)$ (singular part — known analytically)

- $\Phi^R(x^i, t) = \xi_0(t) + \xi_1^i(t)(x^i - c^i) + \xi_2^{ij}(t)(x^i - c^i)(x^j - c^j) + \dots$

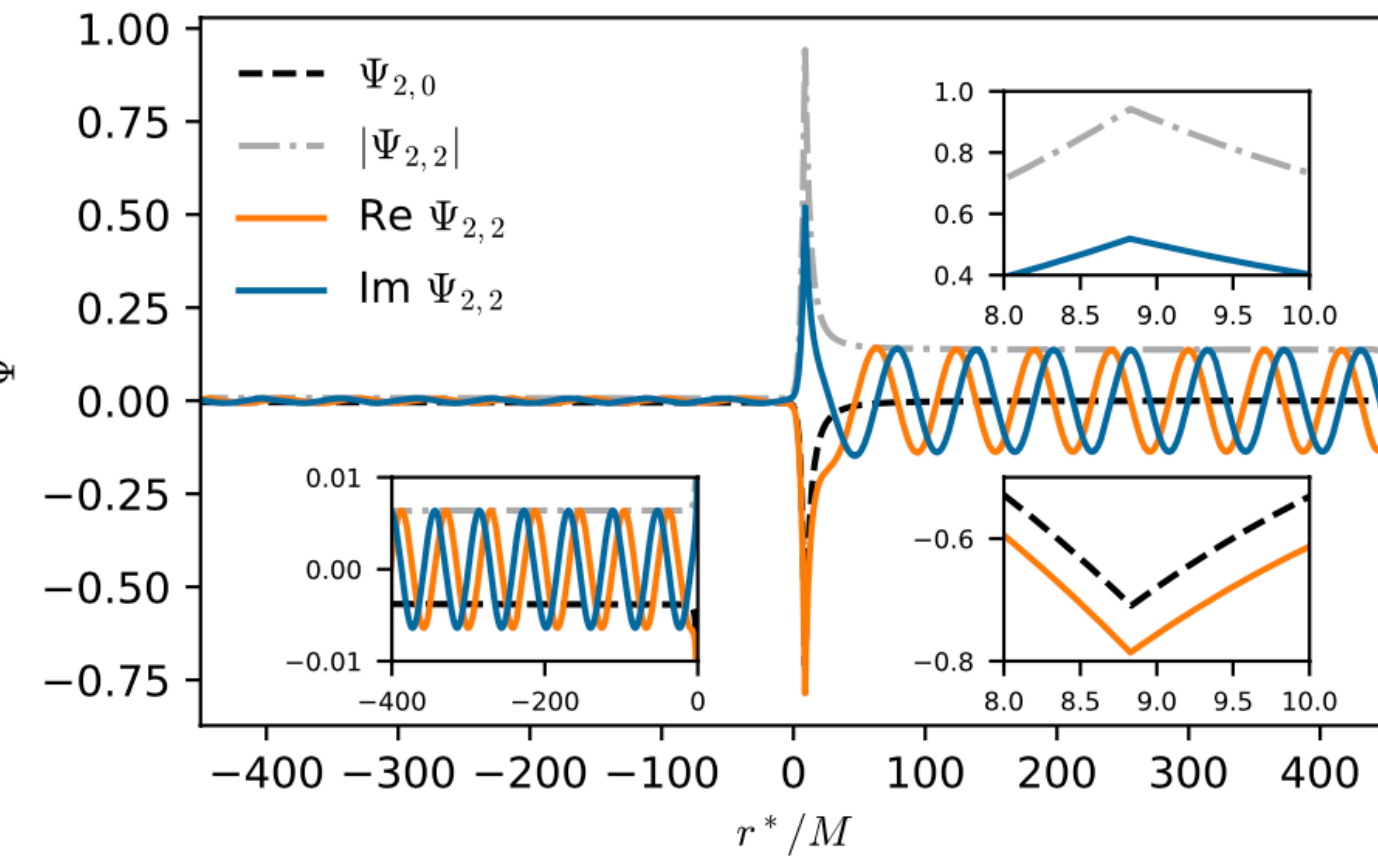
‣ coefficients ξ specify perturbation & need to be determined from NR



So far: Scalar point charge

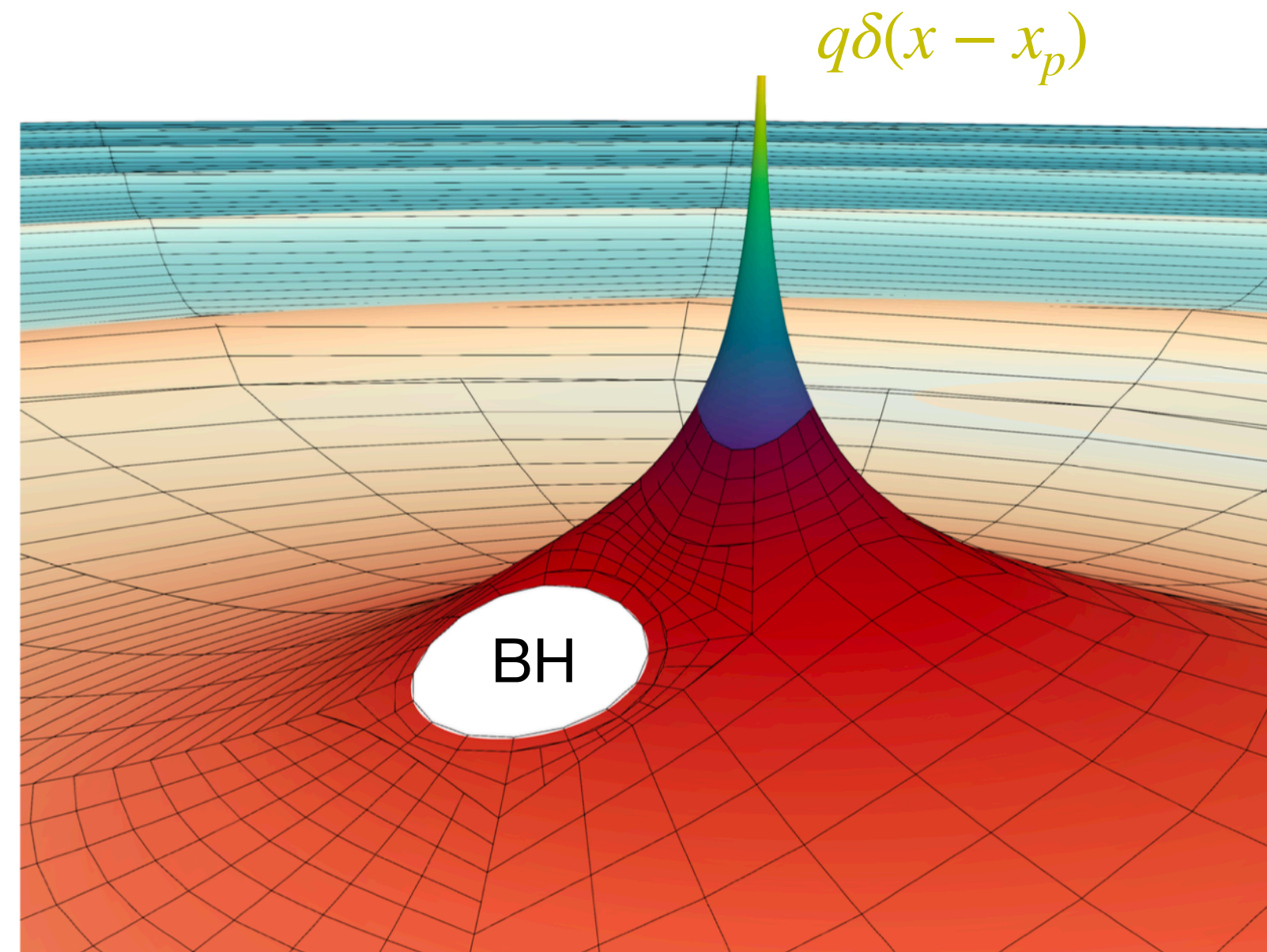


1+1D (circular motion) $\Phi_{lm}(r, t)$

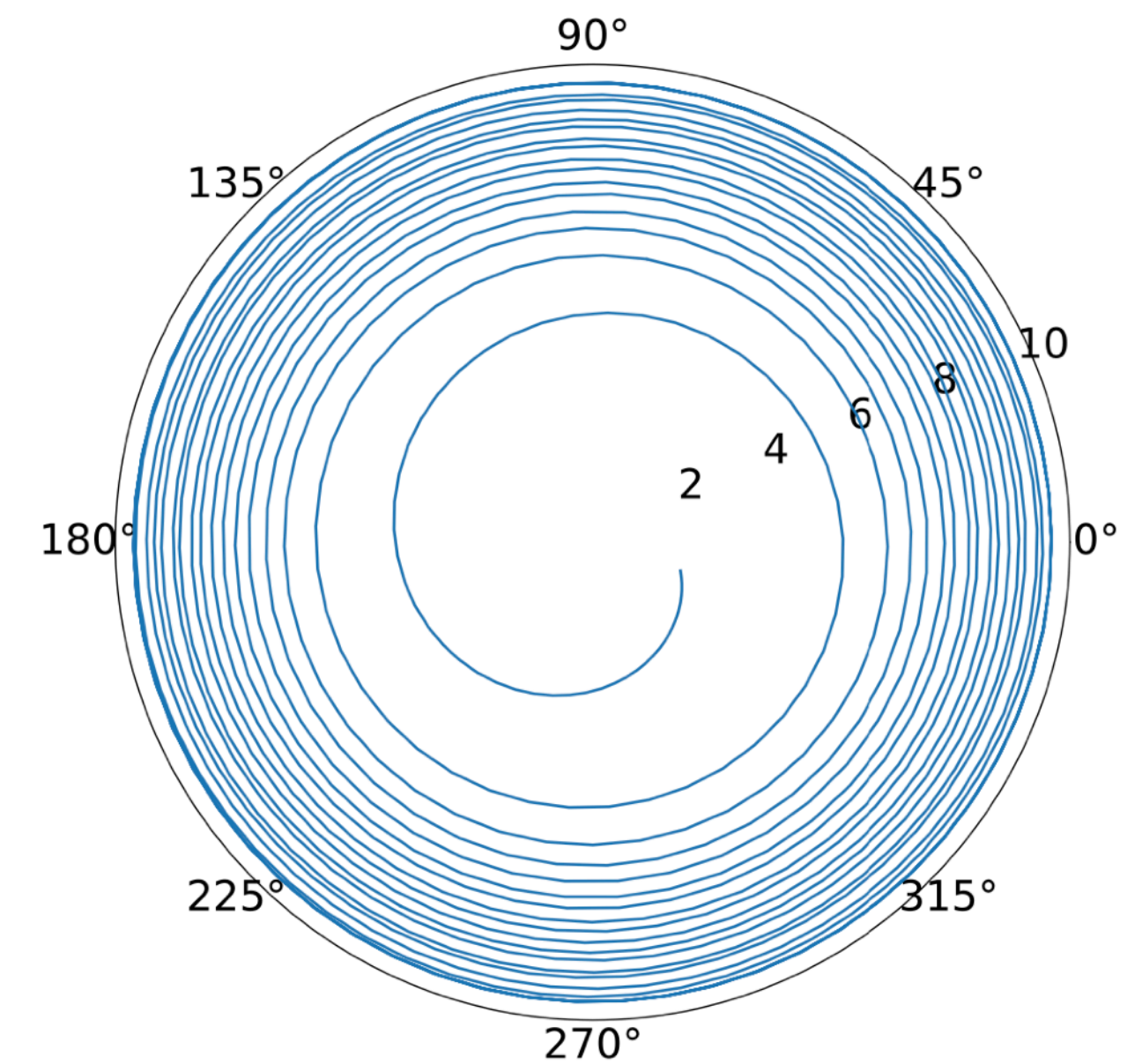


Dhesi, Rüter, Pound, Barack, HP
PRD 104 (2021) 124002

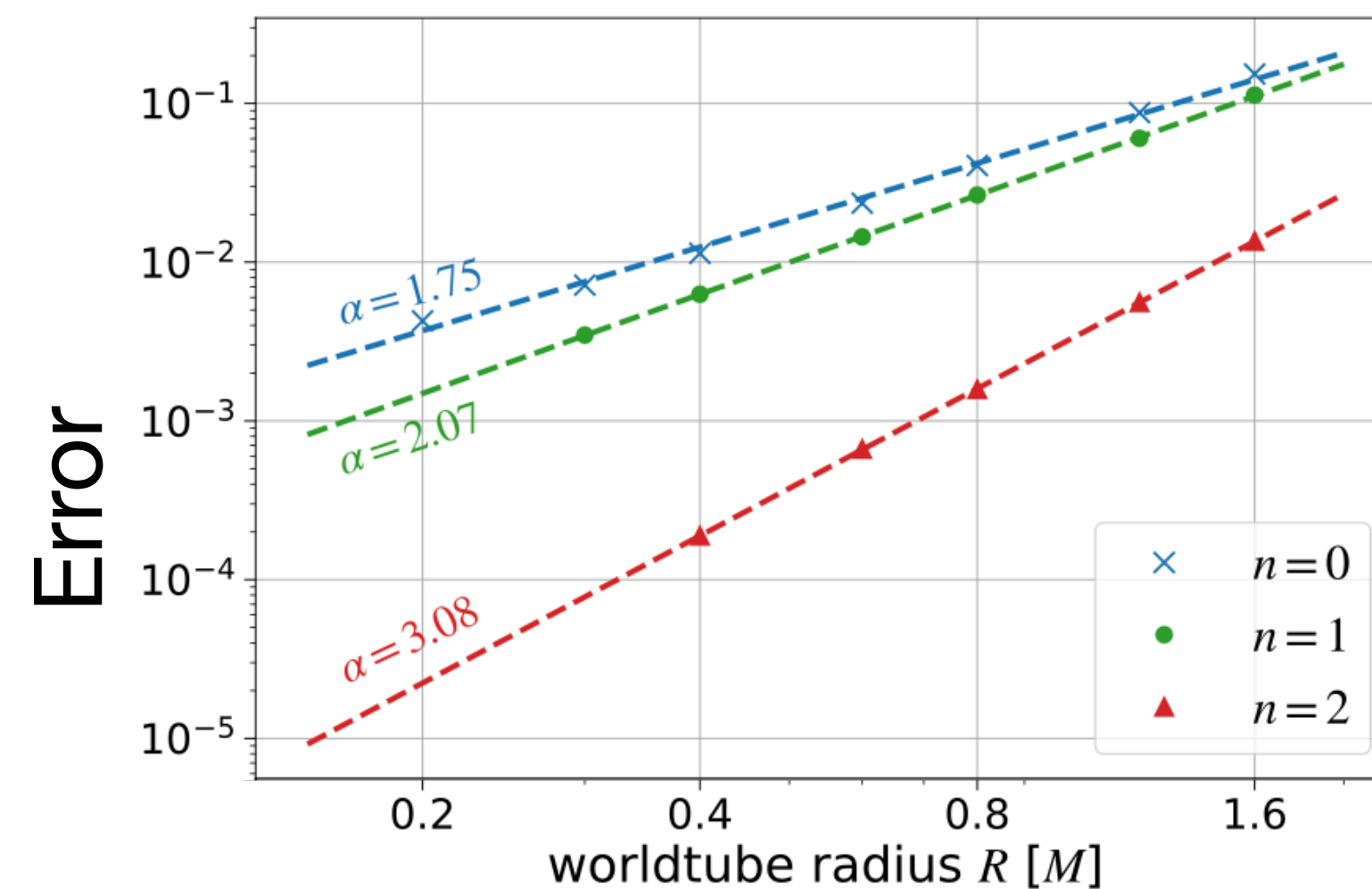
3+1D (circular motion) $\Phi(\vec{x}, t)$



3+1 D (self-force driven inspiral)
 $\vec{x}_p(t)$ and $\Phi(\vec{x}, t)$

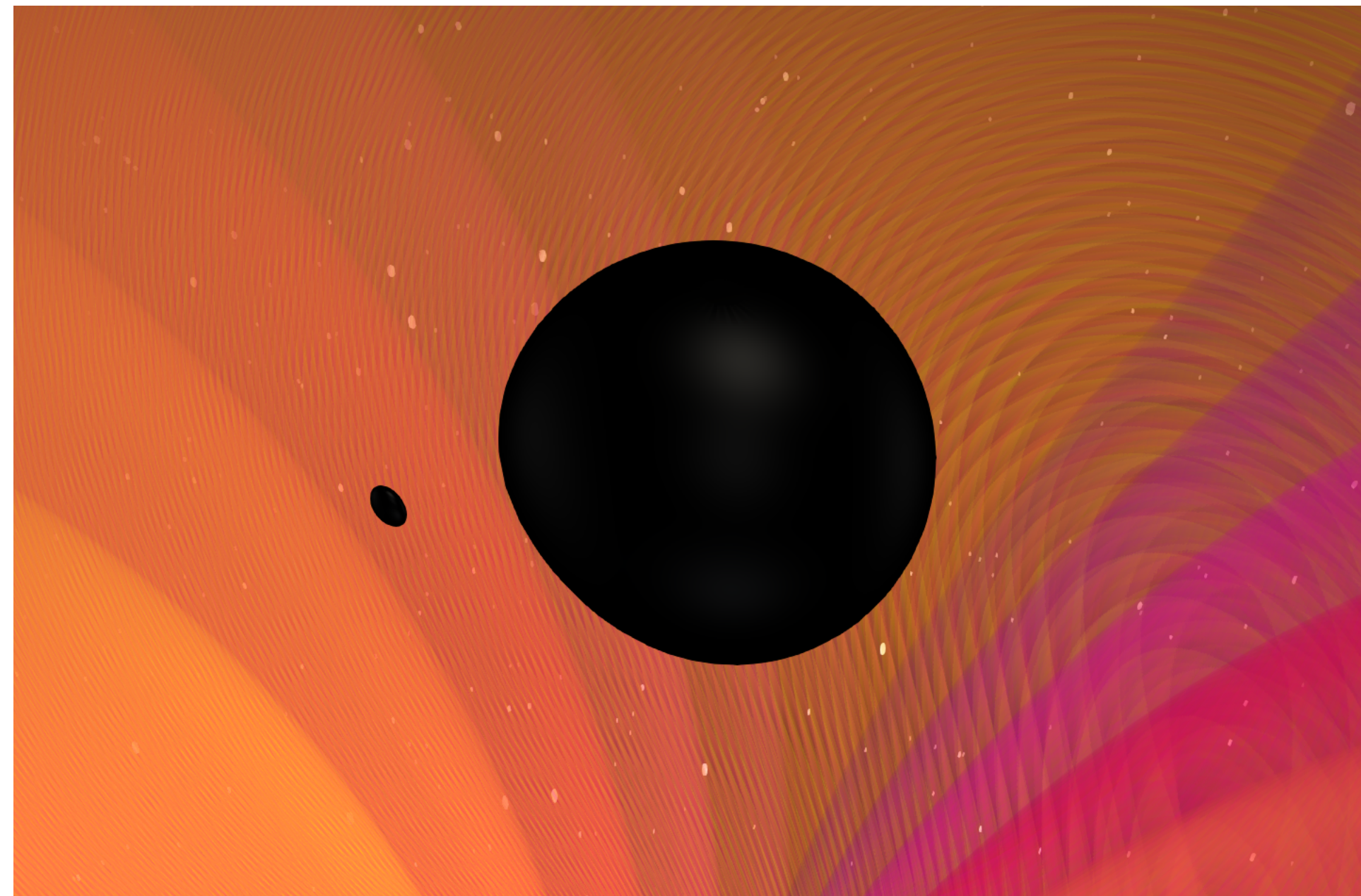


Wittek+ in prep



Wittek, Dhesi, Barack, HP, Pound+
PRD 108 (2023) 024041

Interplay NR & small mass-ratio perturbation theory



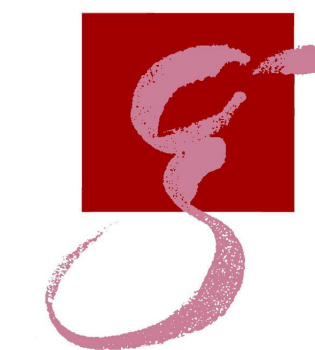
van de Meent & HP, PRL 125 181101 (2020)

— *non-spinning, non-eccentric*

Ramos-Buades, vdMeent, HP+. *PRD 106 124040 (2022)*

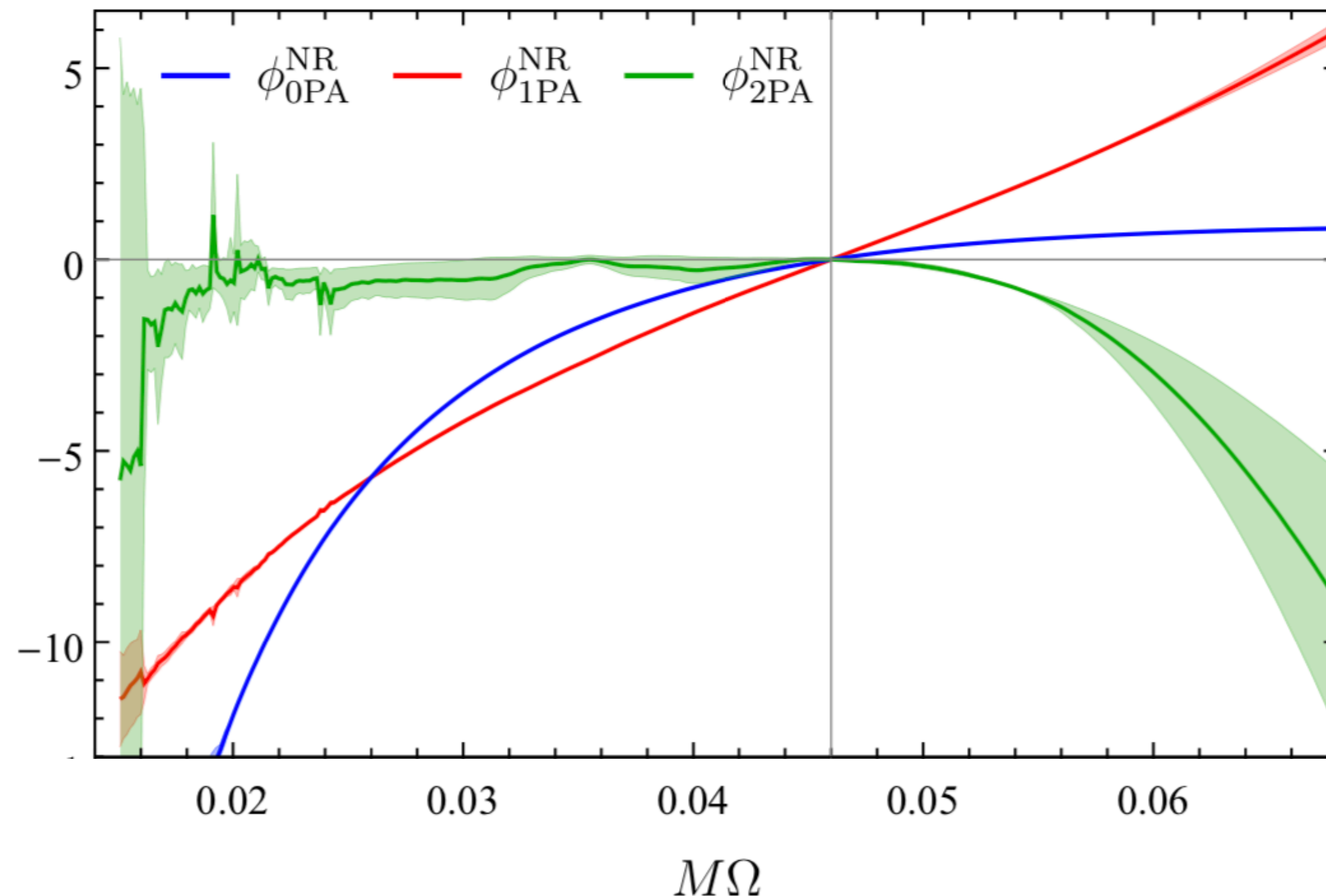
— *non-spinning, eccentric*

Extract GSF information directly from NR



$$\Phi(M\Omega) = \frac{1}{\nu} \Phi_0(M\Omega) + \Phi_1(M\Omega) + \nu \Phi_2(M\Omega) + \dots + \frac{1}{\nu^{1/2}} \Phi_{\text{resonances}} + \frac{1}{\nu^{1/5}} \Phi_{\text{plunge}}$$

- **Fit $\Phi_{\text{NR}}(\Omega_{\text{NR}})$ to SMR expansion**
 - 55 NR sims at different ν
 - non-spinning
- Different orders in ν
 - $\Phi_0(M\Omega)$ - agrees with 0PA
 - $\Phi_1(M\Omega)$ - hereby computed
 - $\Phi_2(M\Omega)$ - remarkably small

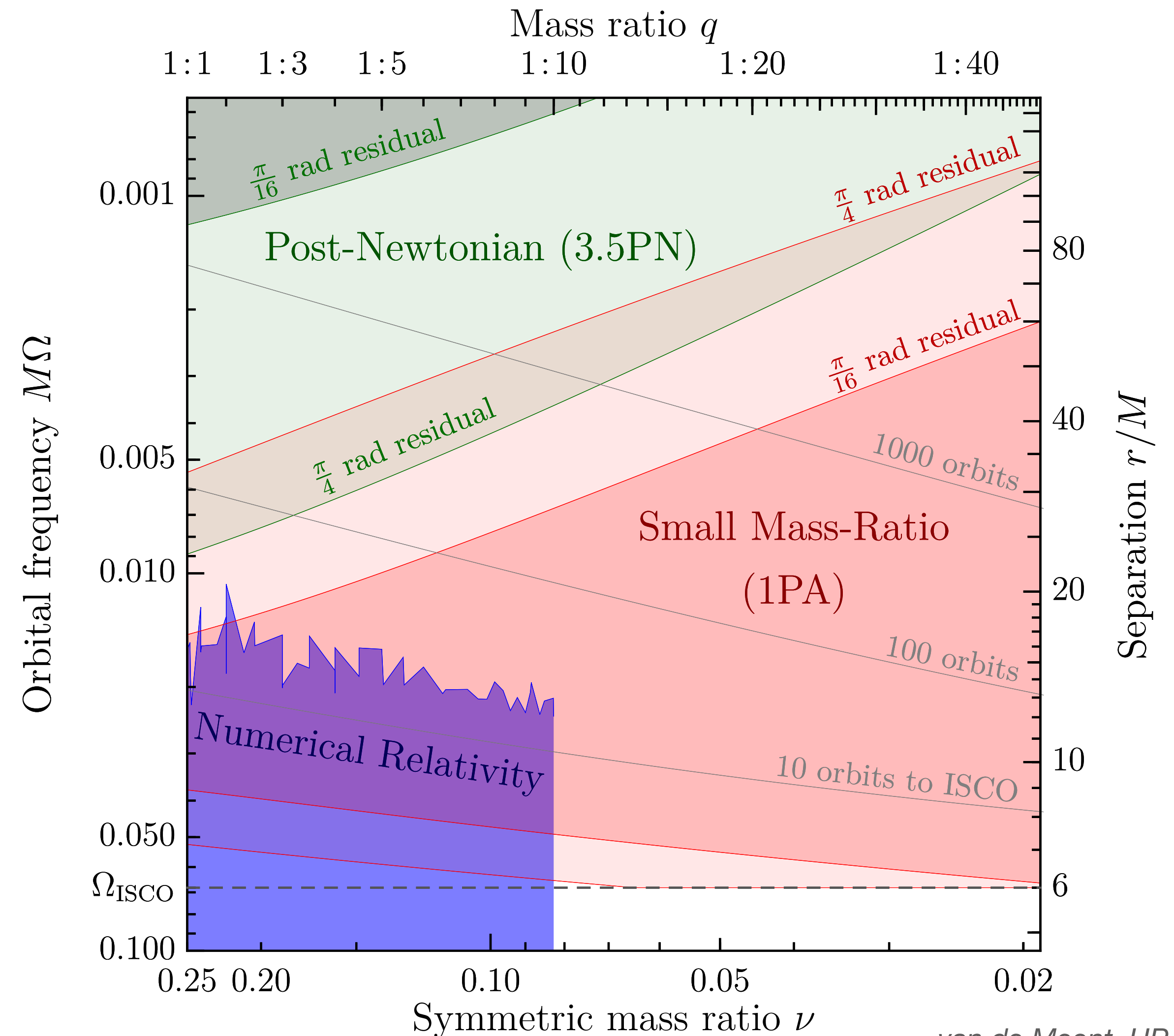


Extract GSF information directly from NR



$$\Phi(M\Omega) = \frac{1}{\nu} \Phi_0(M\Omega) + \Phi_1(M\Omega) + \nu \Phi_2(M\Omega) + \dots + \frac{1}{\nu^{1/2}} \Phi_{\text{resonances}} + \frac{1}{\nu^{1/5}} \Phi_{\text{plunge}}$$

- **Fit $\Phi_{\text{NR}}(\Omega_{\text{NR}})$ to SMR expansion**
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- Different orders in ν
 - $\Phi_0(M\Omega)$ - agrees with 0PA
 - $\Phi_1(M\Omega)$ - hereby computed
 - $\Phi_2(M\Omega)$ - remarkably small
- **Predict region of validity of SMR**
 - quasi-circular & non-spinning

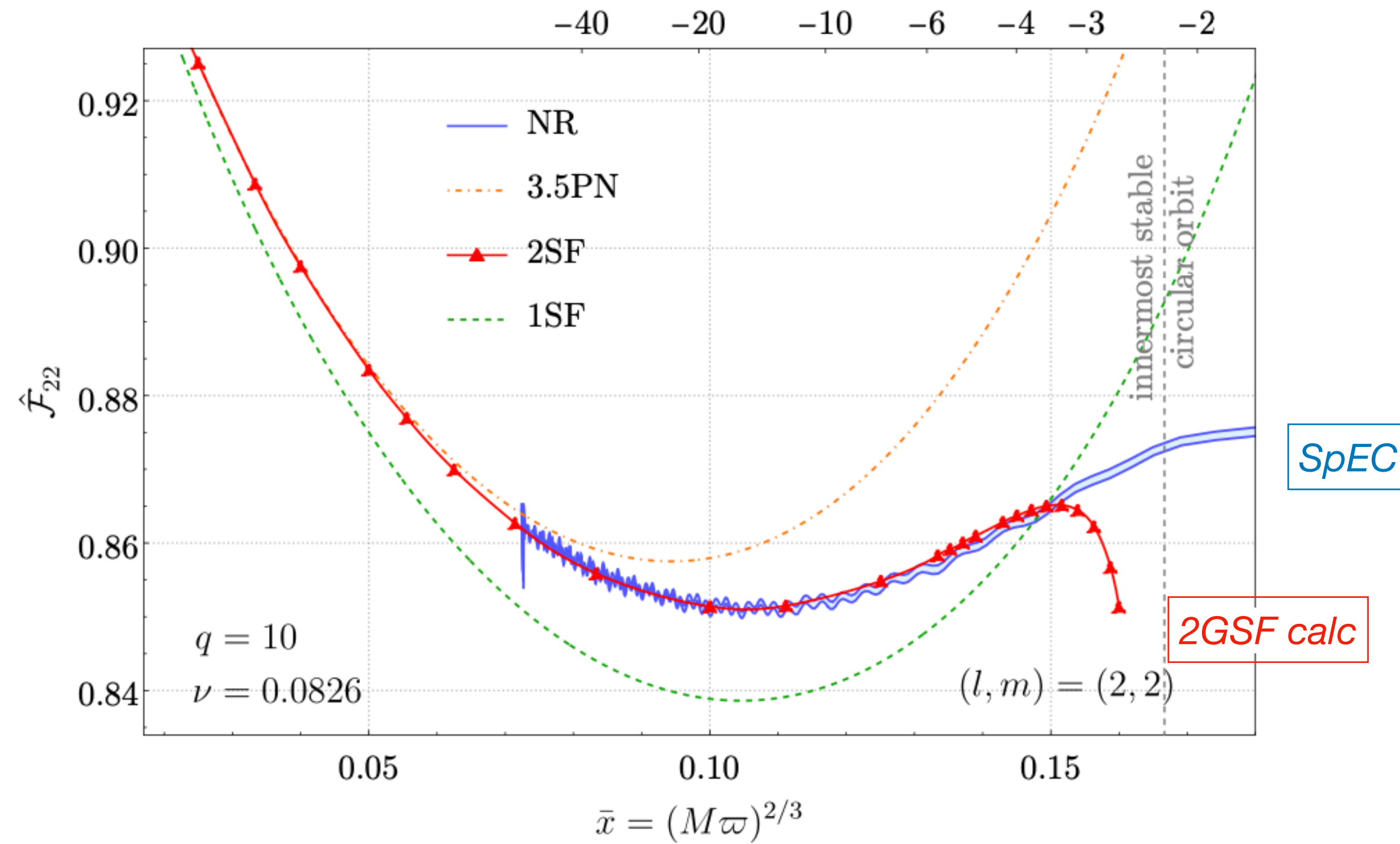


Confirmed by actual 2-SMR calculations

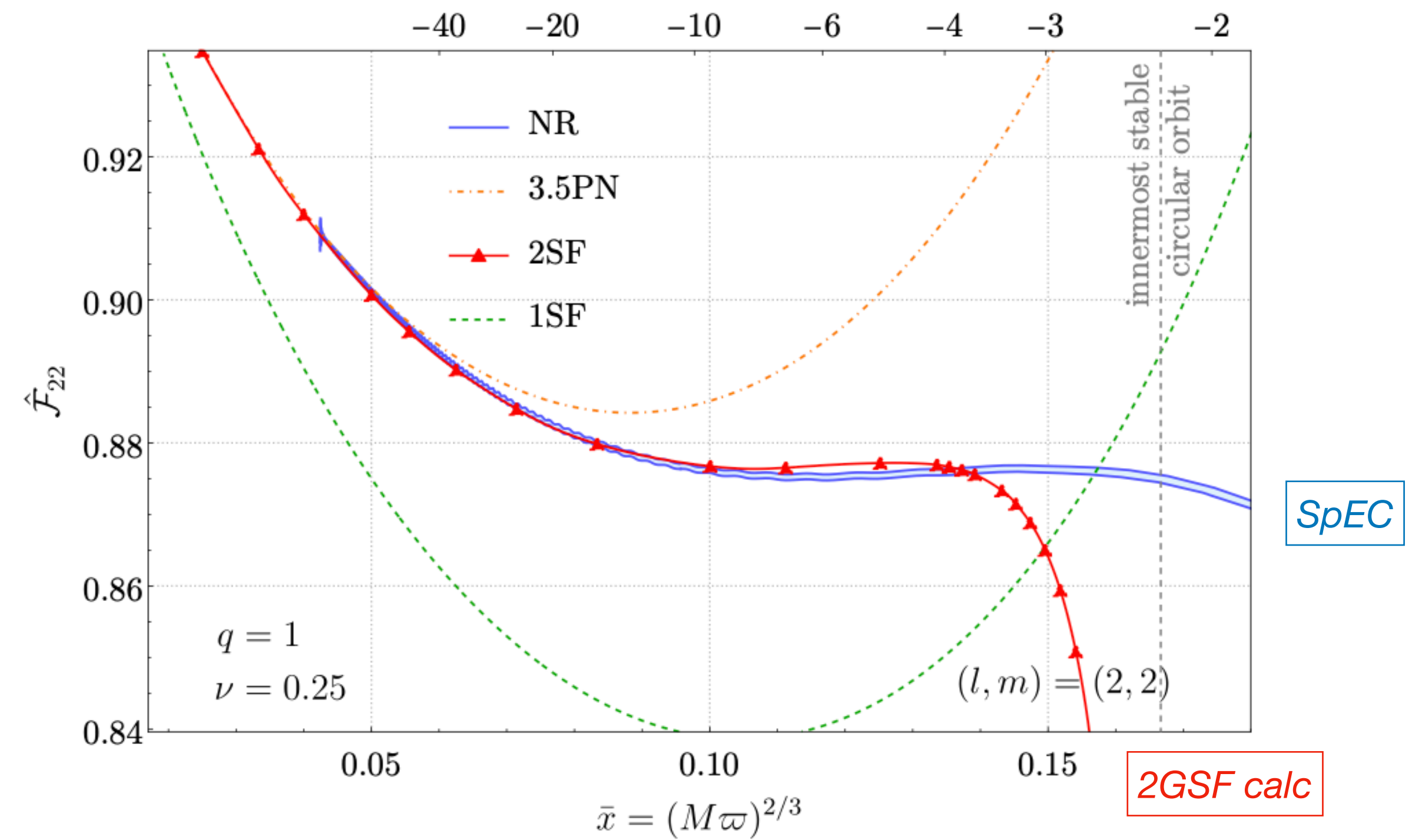


Warburton, Pound, Wardell, Miller, Durkan 21
 Calculation of GW energy flux to $O(q^2)$

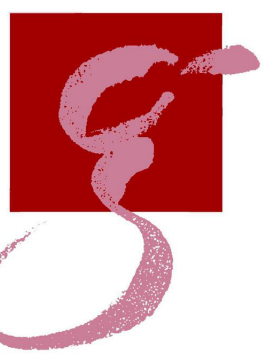
$q=1/6$



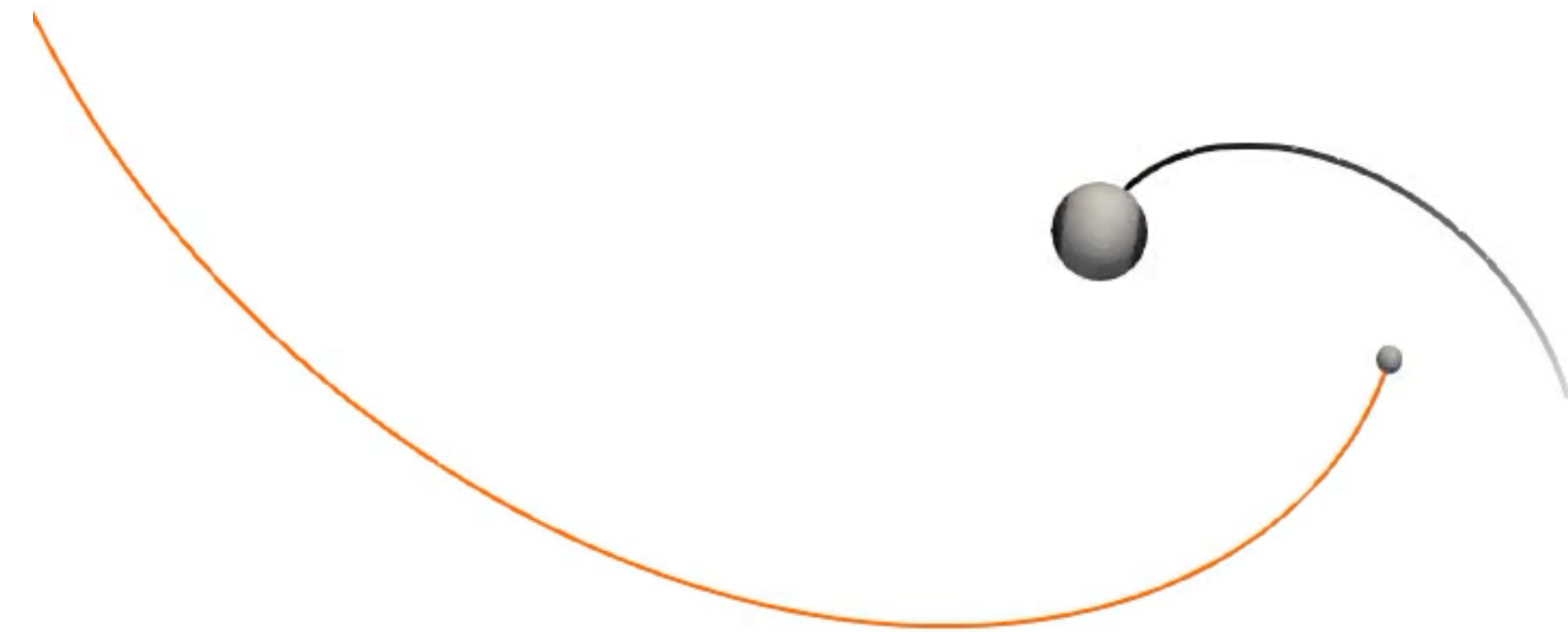
$q=1 !!$



Eccentric BBH simulations



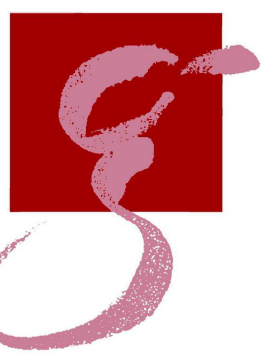
- **Why** eccentric NR?
 - **measure** eccentricity of GW signals
 - **distinguish “deviation from GR”** from “eccentric GR”
 - **EMRIs will be eccentric** in LISA



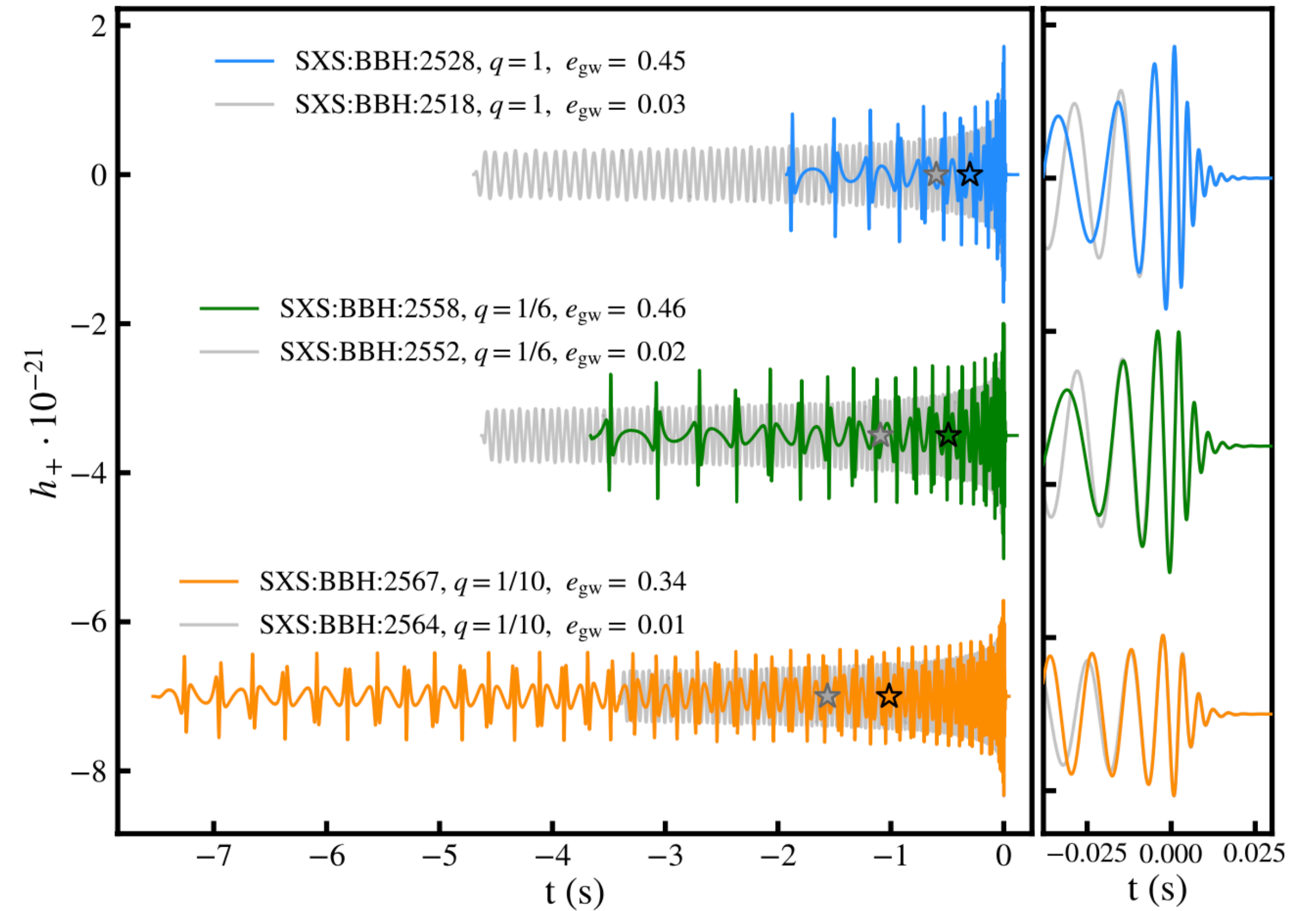
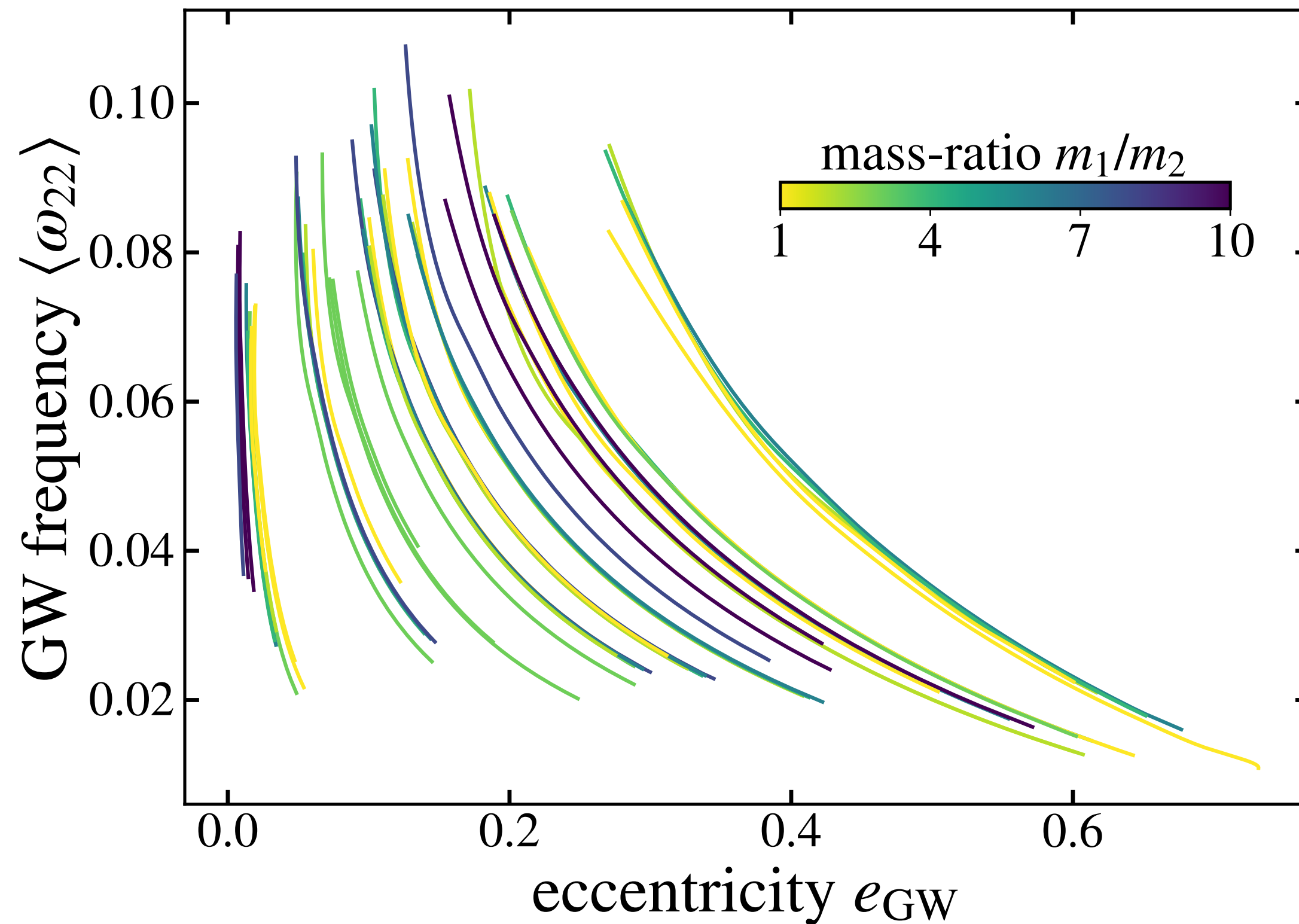
NB:

- *All animations based on full 3-D PDE solutions, not just ODEs for centers of BHs.*
- *BH Horizons shaded by their curvature, showing tidal deformations at periastron passages*

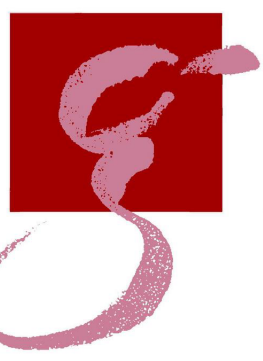
Eccentric NR, $q=1-0.1$



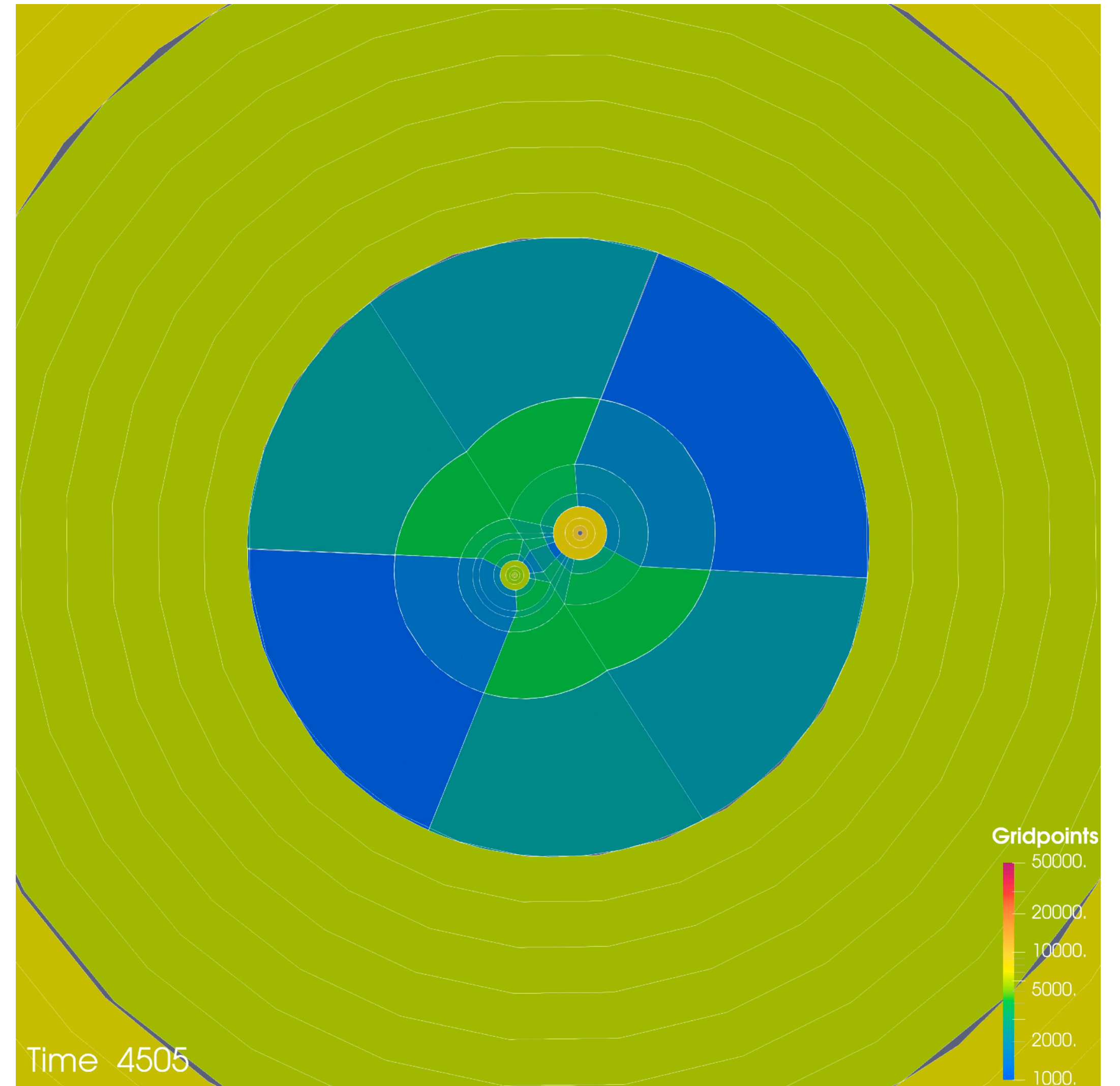
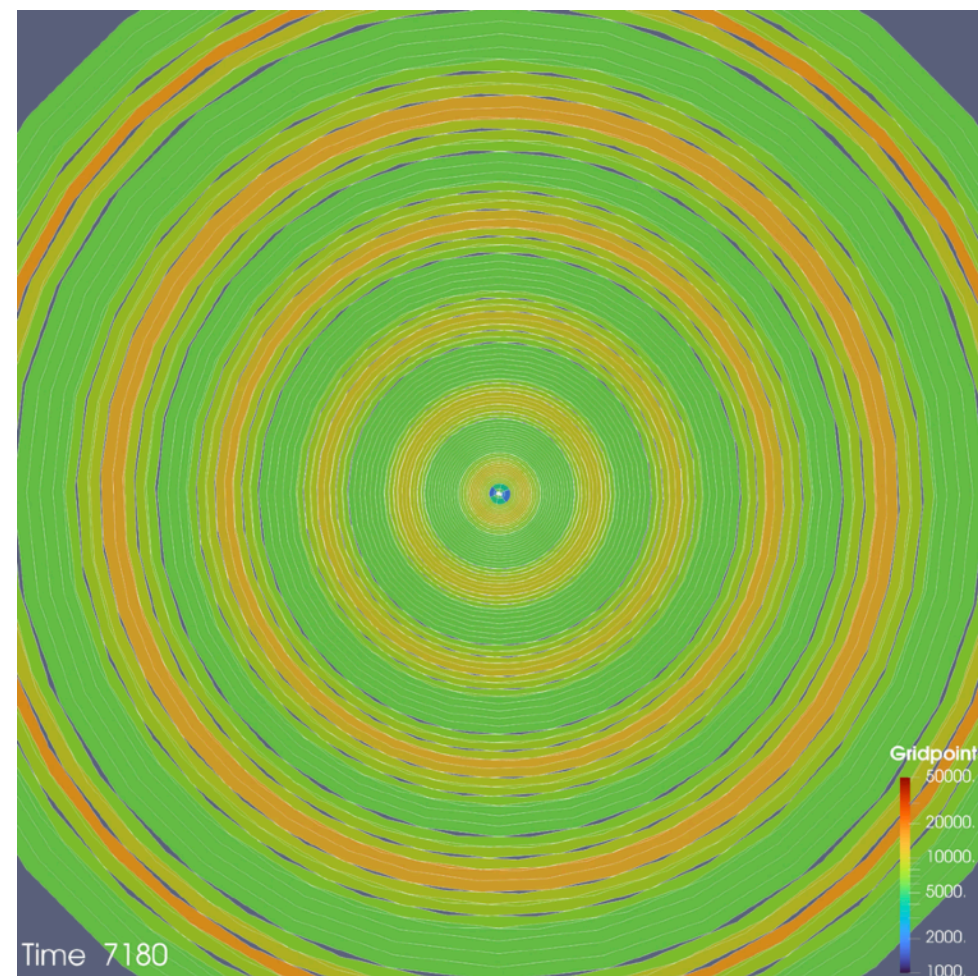
- new NR sims
 - $q=1-0.1$, $e=0-0.7$
 - three resolutions each
 - NR errors analysed



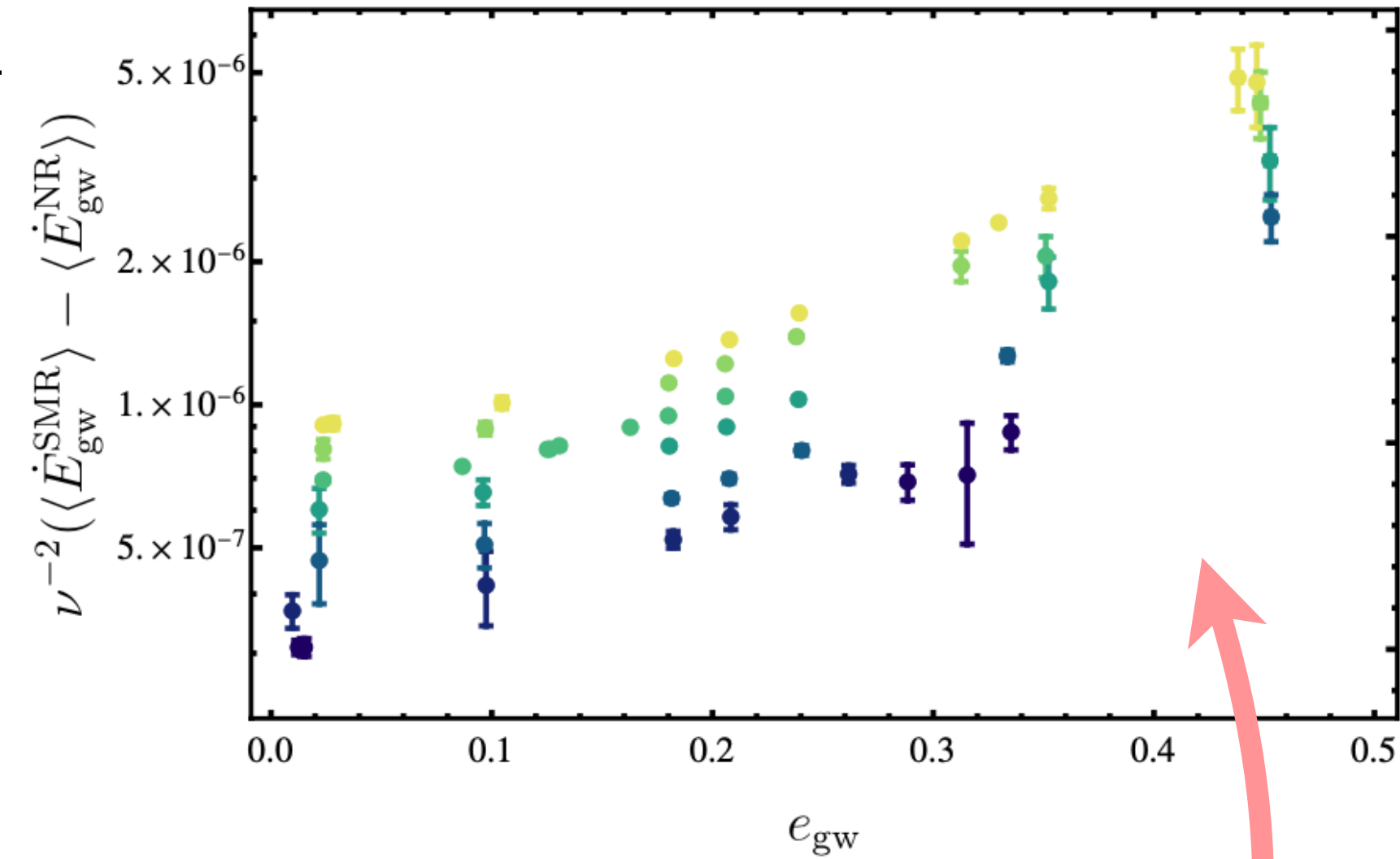
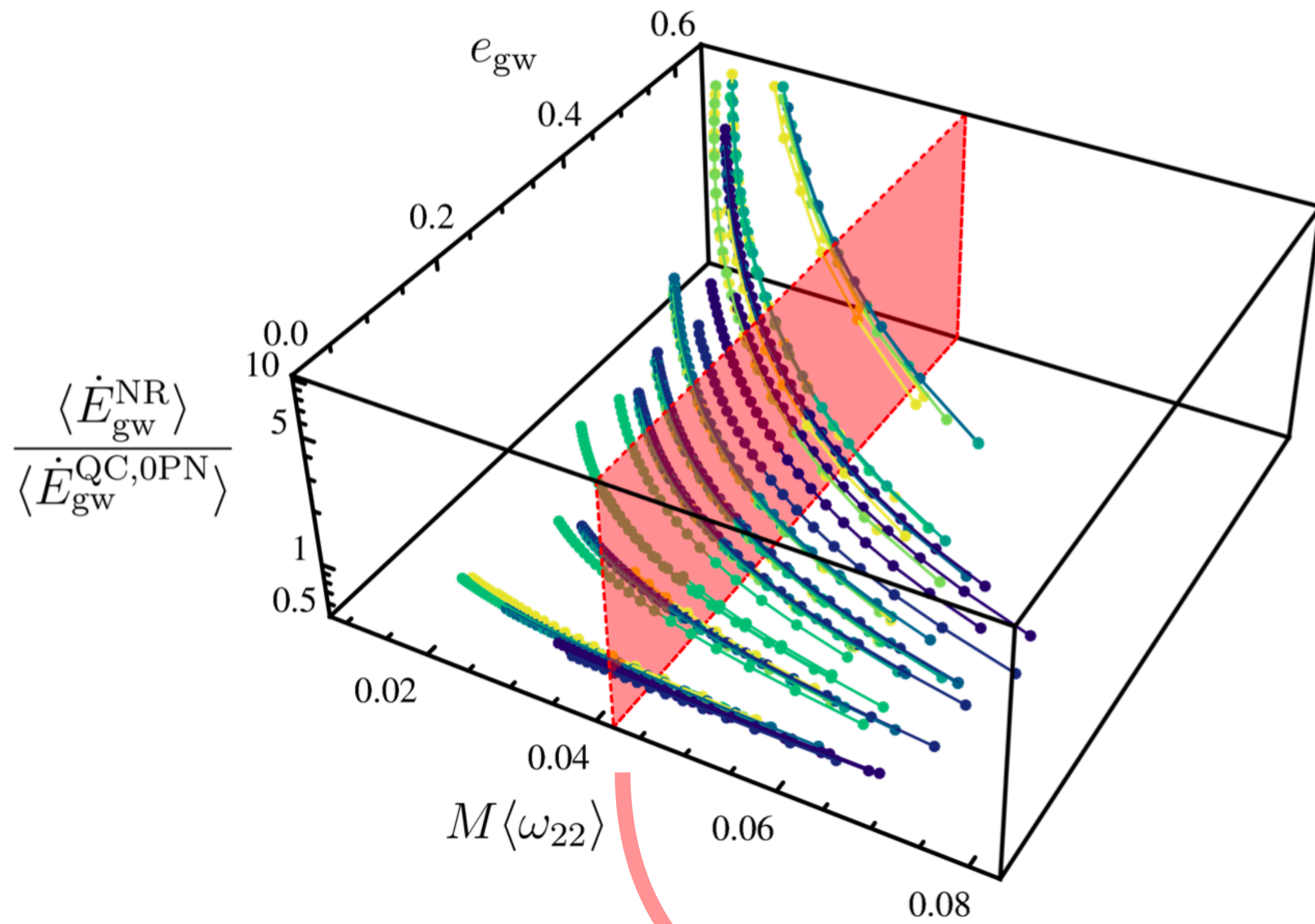
Eccentric NR, $q=1-0.1$



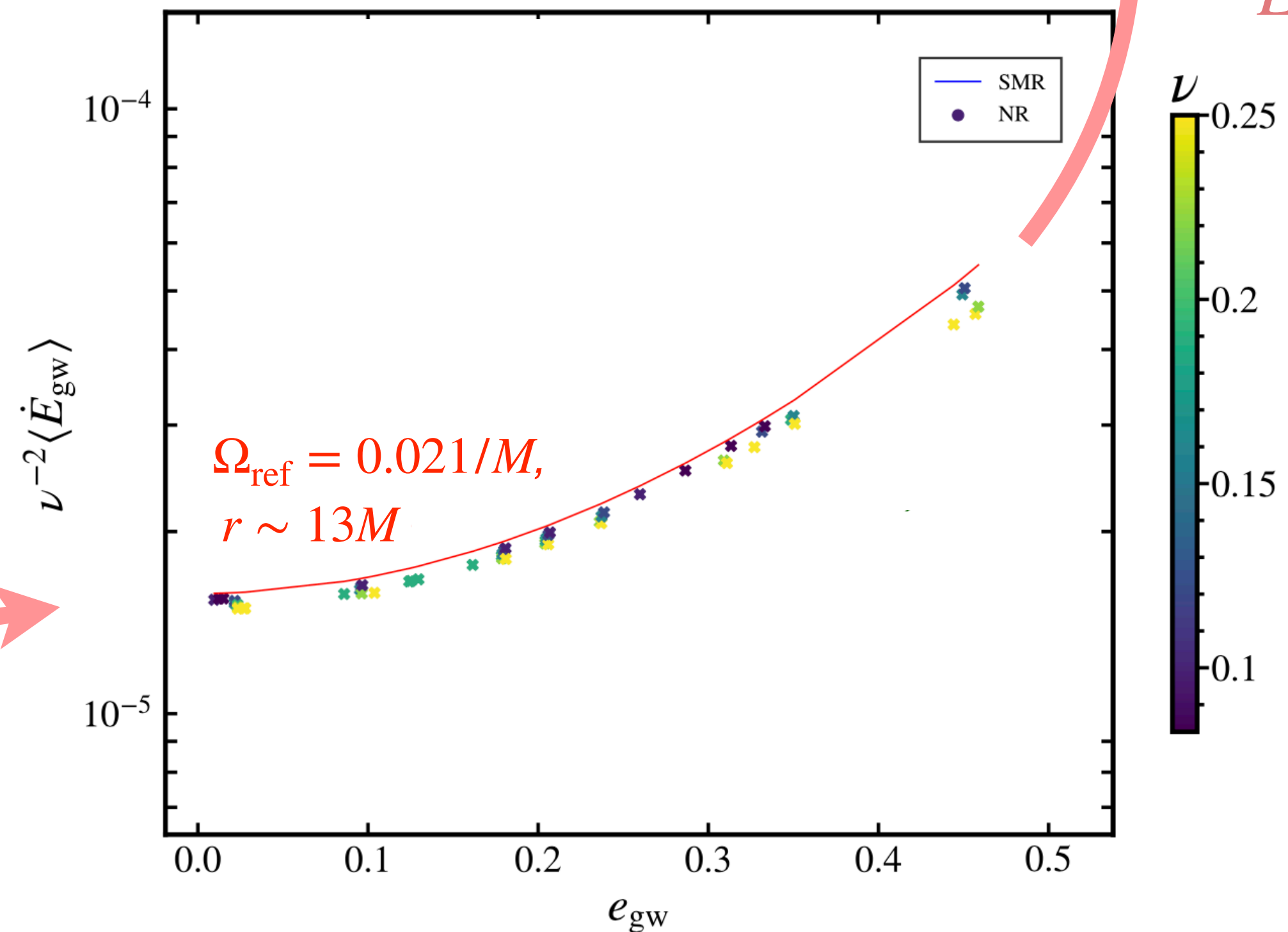
- new NR sims
 - $q=1-0.1$, $e=0-0.7$
 - three resolutions each
 - NR errors analysed
- **GW pulses visible in AMR**



$\langle \dot{E}_{\text{GW, NR}} \rangle$ for 52 simulations

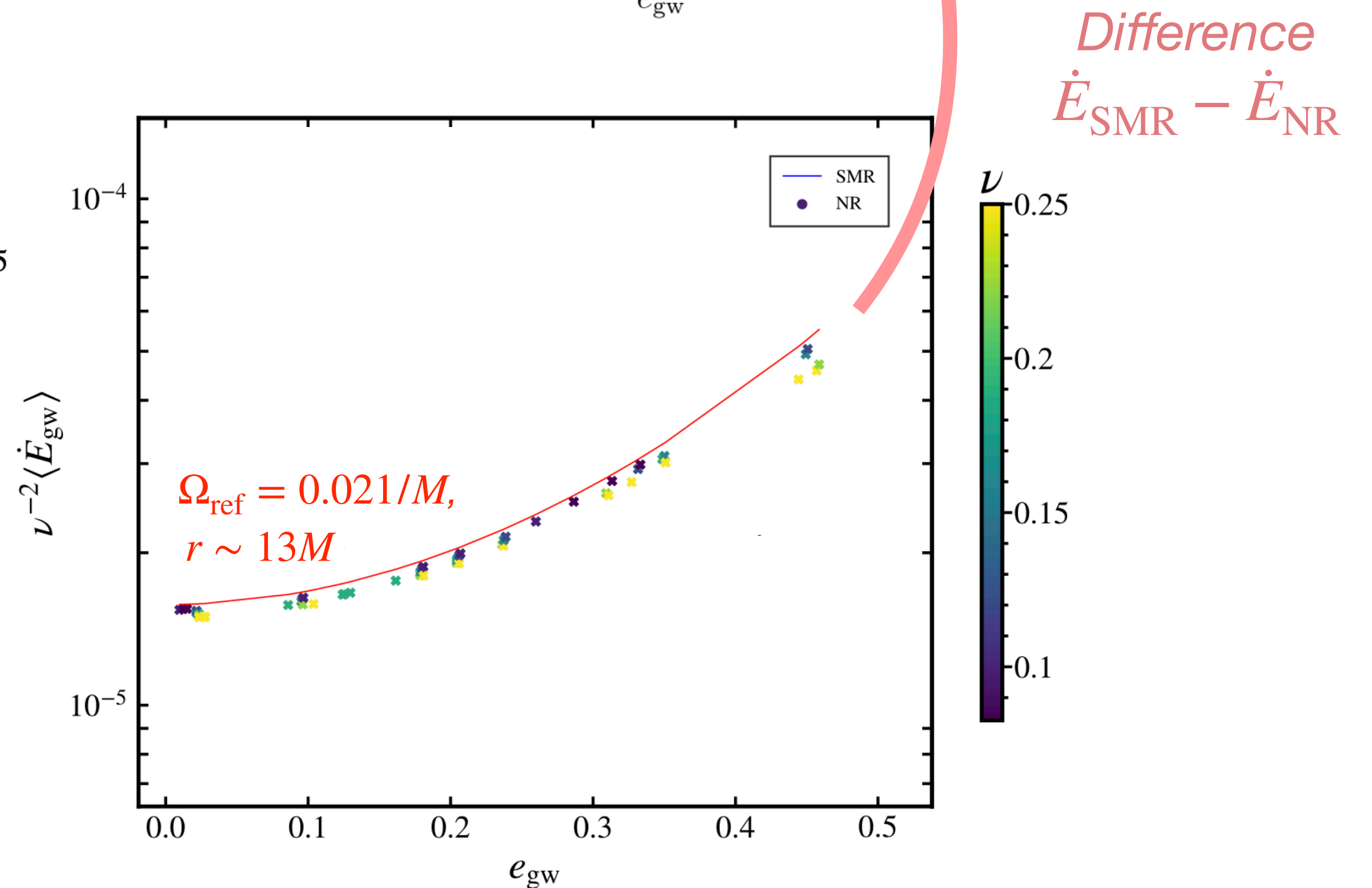
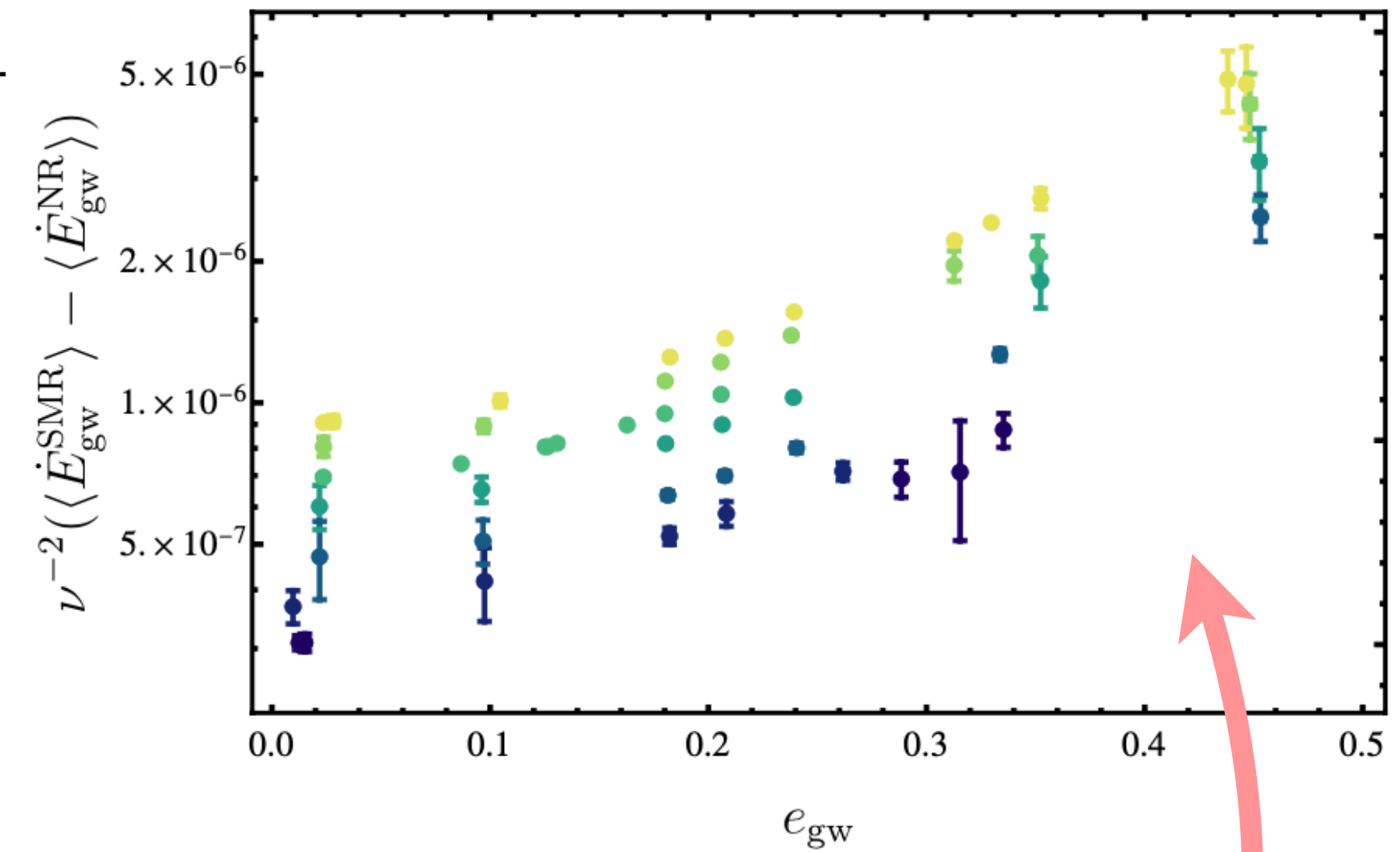
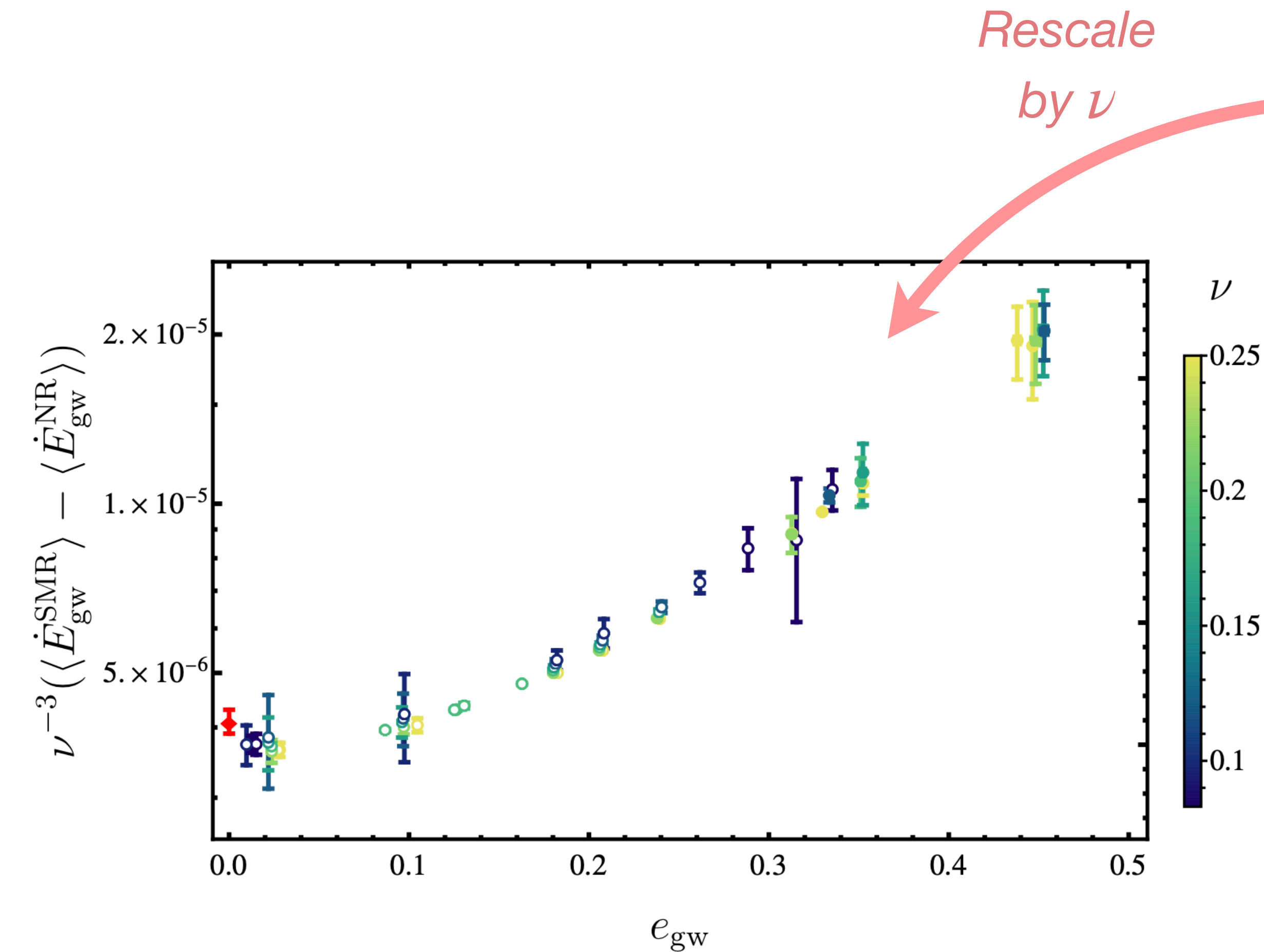


Difference
 $\dot{E}_{\text{SMR}} - \dot{E}_{\text{NR}}$



At fixed Ω_{ref}
accounting for $\dot{E} \propto \nu^2$

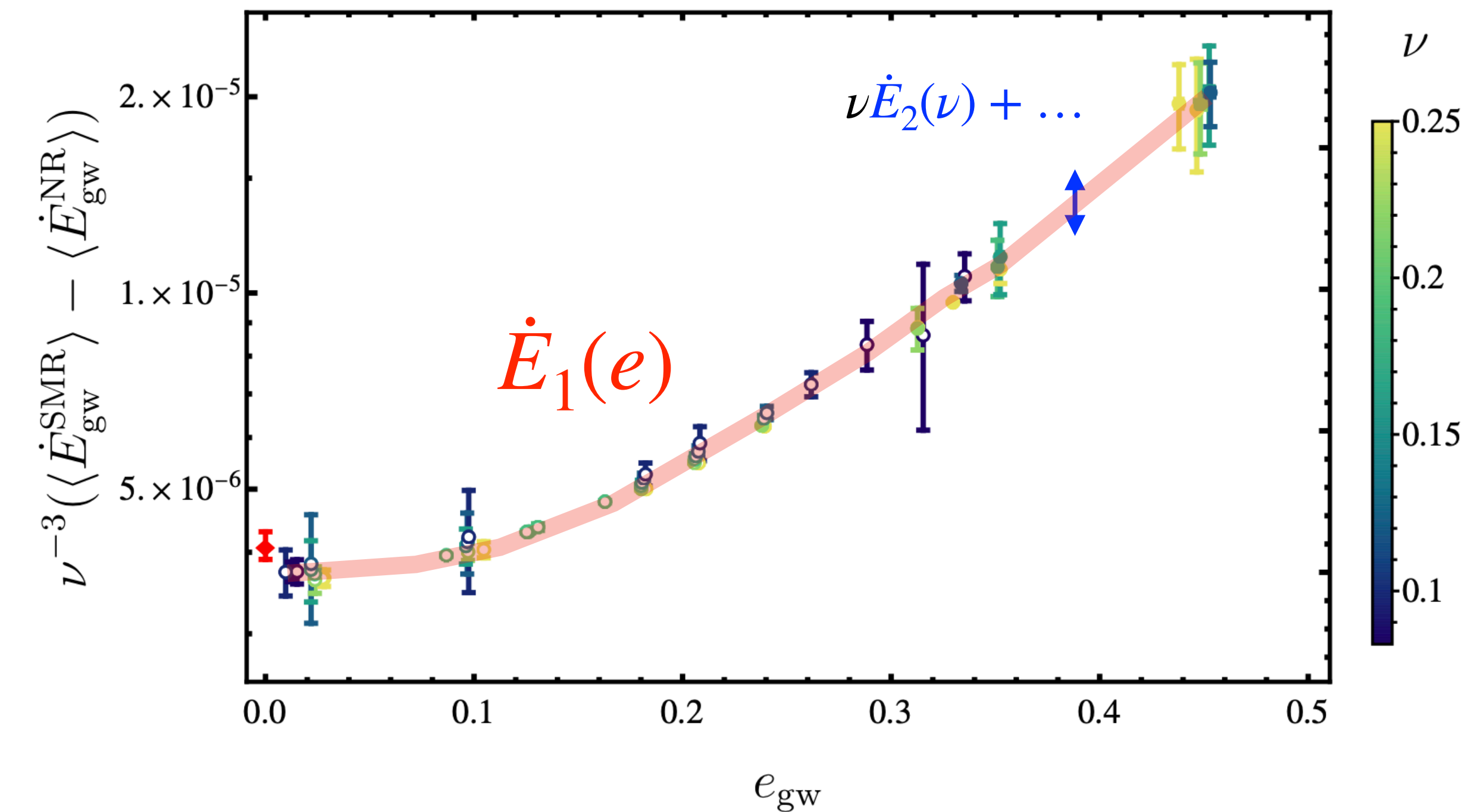
$\langle \dot{E}_{\text{GW, NR}} \rangle$ for 52 simulations





GW fluxes at fixed frequency

$$\dot{E}_{\text{GW}}(\nu, e) = \nu^2 \dot{E}_0(e) + \nu^3 \dot{E}_1(e) + \nu^4 \dot{E}_2(e) + \dots$$



- **Data collapses to one line:**
 $(\text{SMR} - \text{NR}) \propto \nu^3$

⇒ NR recovers \dot{E}_0 from SMR

⇒ collapsed curve gives next-to-leading order contribution

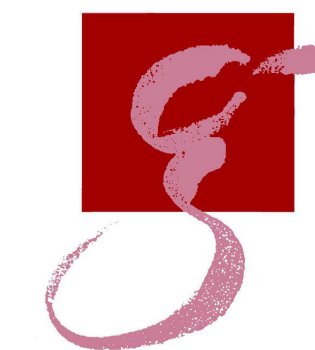
$$\dot{E}_1(e) = \mathcal{O}(\nu^3) - \text{unknown before}$$

⇒ nearly vanishing spread: next-to-next-to-leading-order contribution is small

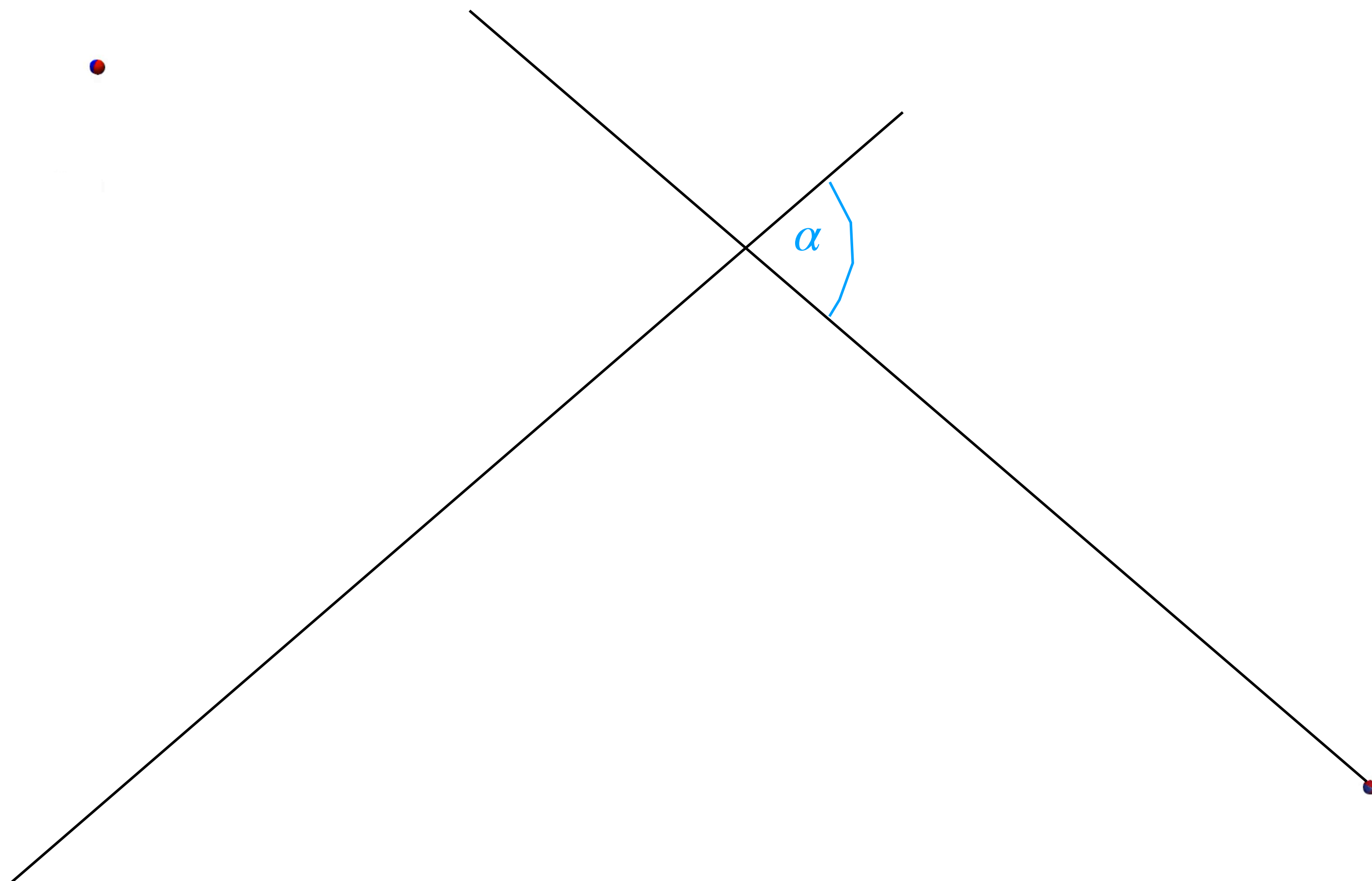
$$\dot{E}_2(e) = \mathcal{O}(\nu^4) - \text{insignificant}$$

BH scattering

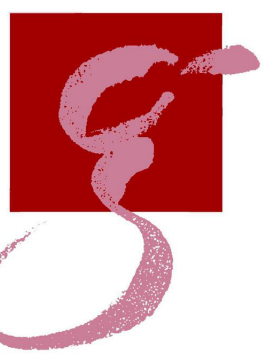
BH Scattering



- **Clean probe** of strong field dynamics
- **Gauge invariant** quantities
 - scattering angle(s) α
 - $E_{\text{GW}}, L_{\text{GW}}, \Delta v_{\infty, i}$
 - $\Delta m_i, \Delta \vec{S}_i$

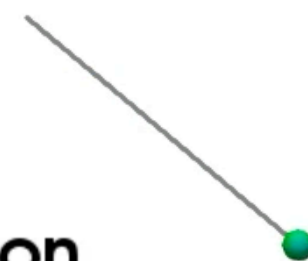
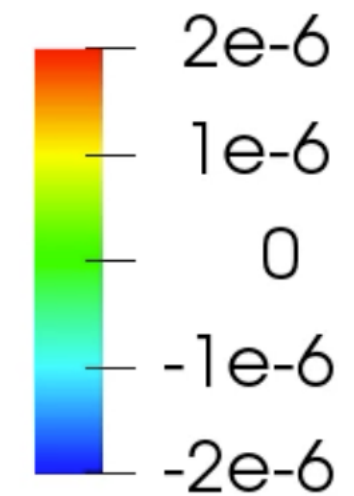


BH Scattering

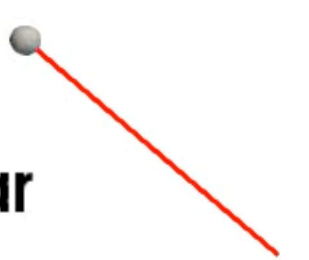
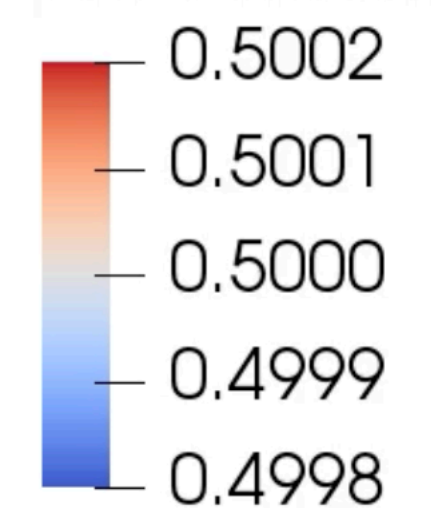


- **Clean probe** of strong field dynamics
- **Gauge invariant** quantities
 - scattering angle(s) α
 - $E_{\text{GW}}, L_{\text{GW}}, \Delta v_{\infty, i}$
 - $\Delta m_i, \Delta \vec{S}_i$
- **Time-dependent data**
 - GWs
 - Horizons

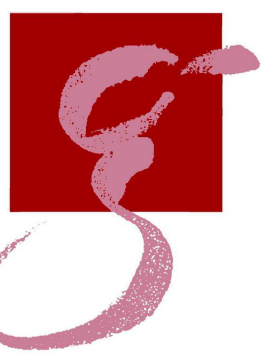
SpinFunction



DimlessRicciScalar

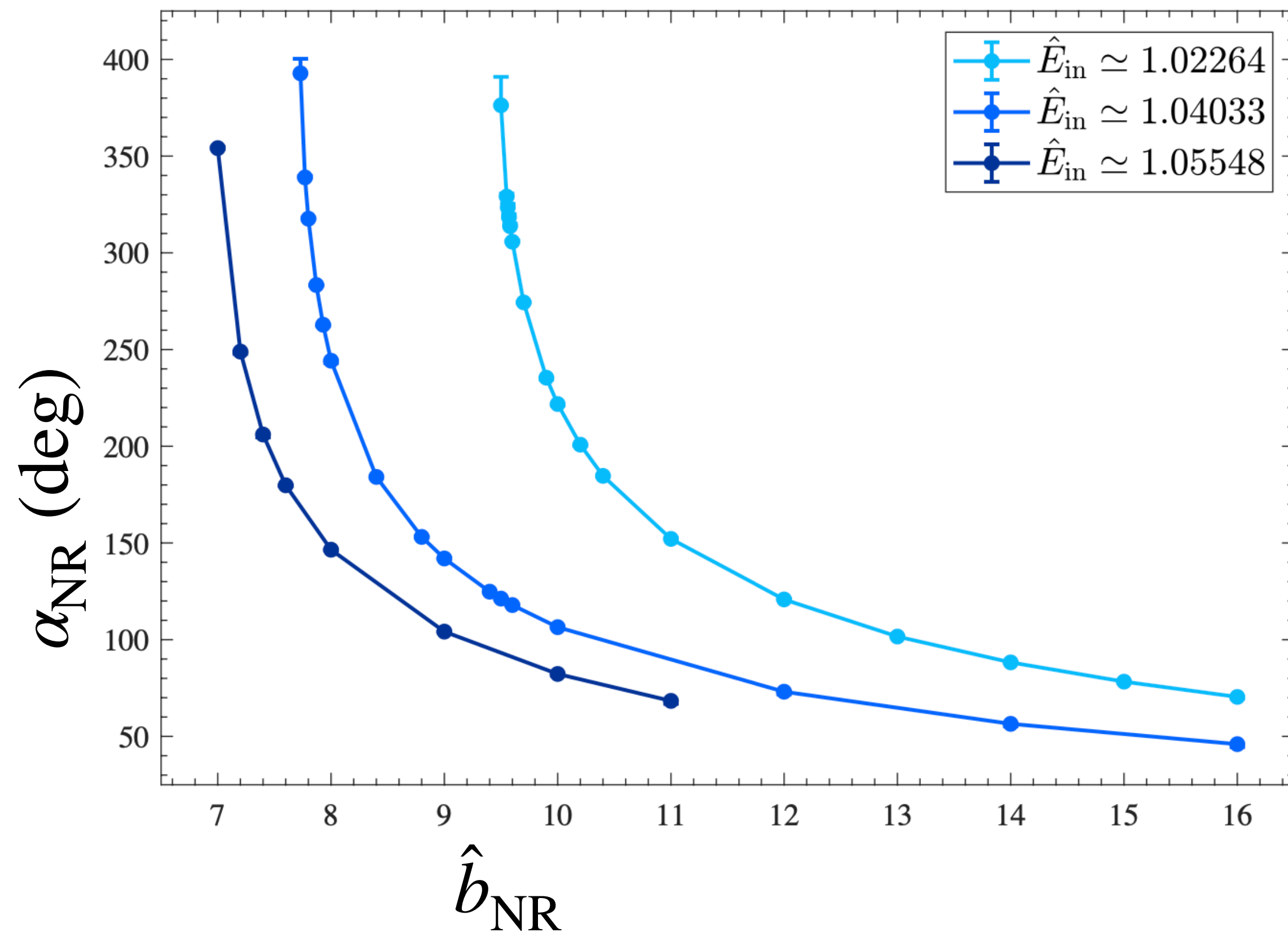


Scattering angles from NR

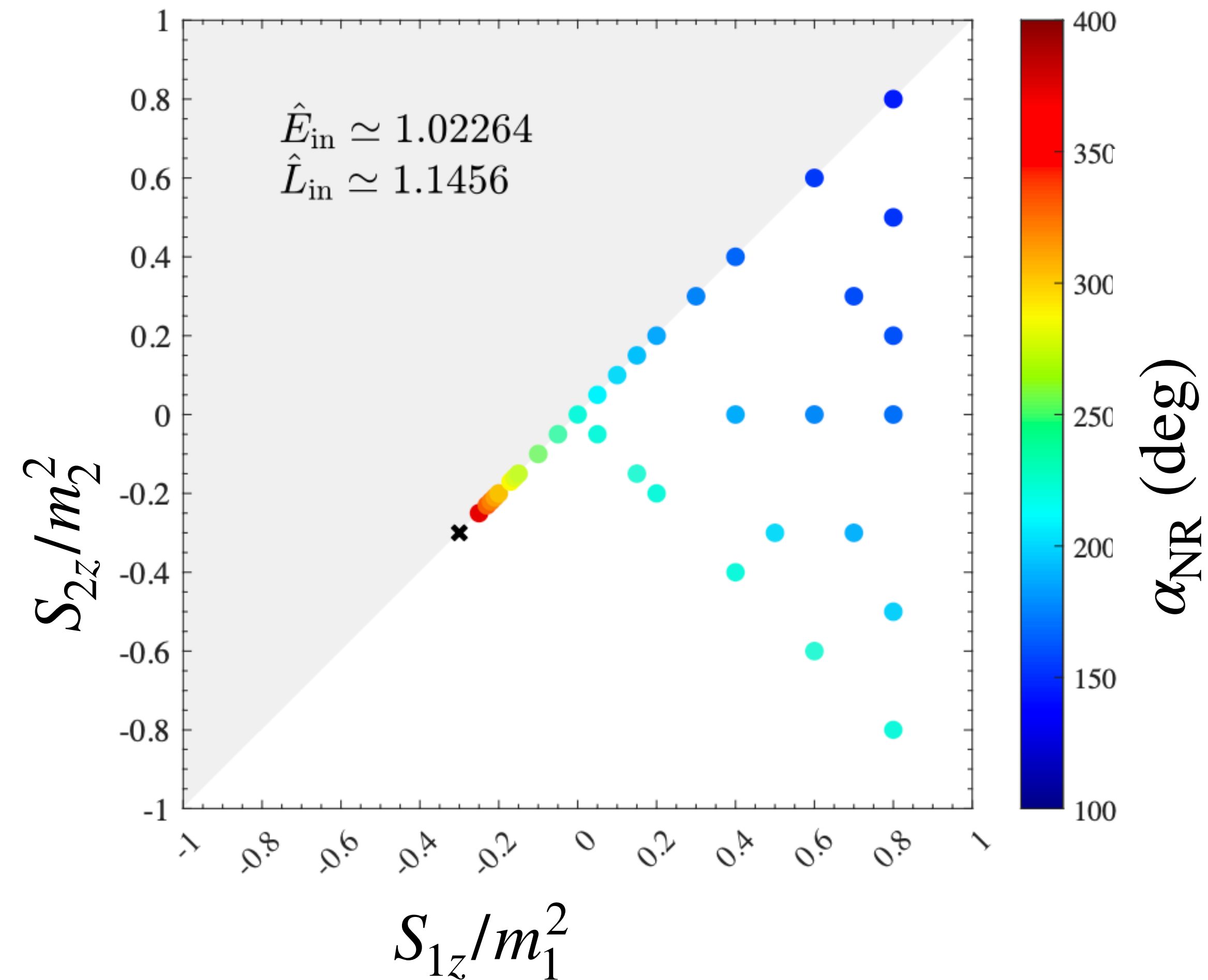


Most impressive results: Rettegno, Pratten, Thomas, Schmidt, Damour 2307.06999

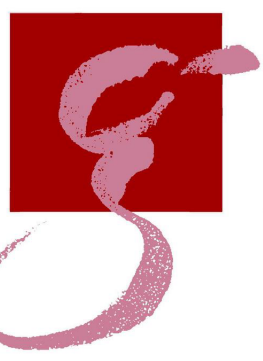
non-spinning



spinning BHs

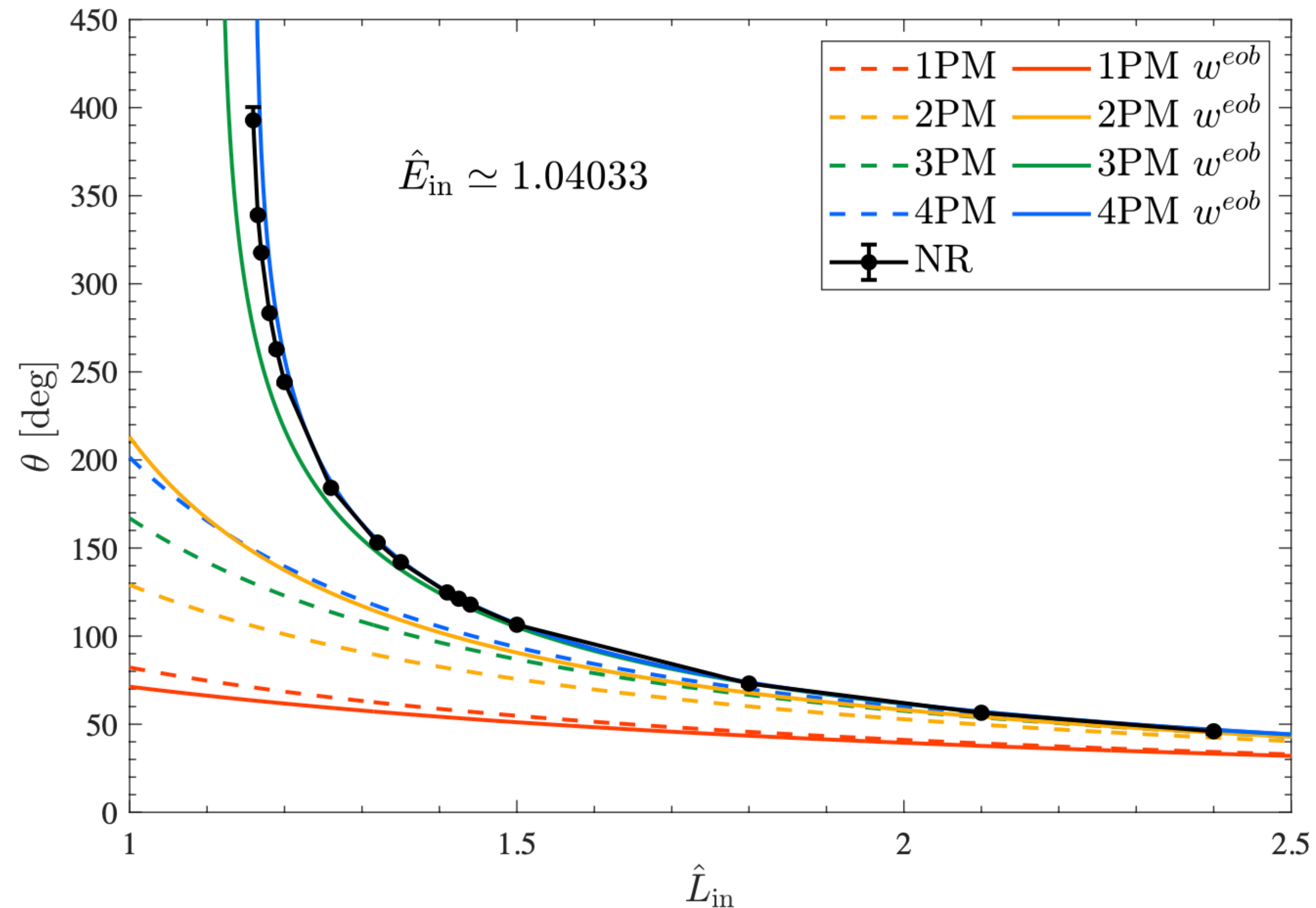


Scattering angles from NR



Most impressive results: Rettegno, Pratten, Thomas, Schmidt, Damour 2307.06999

non-spinning, equal-mass

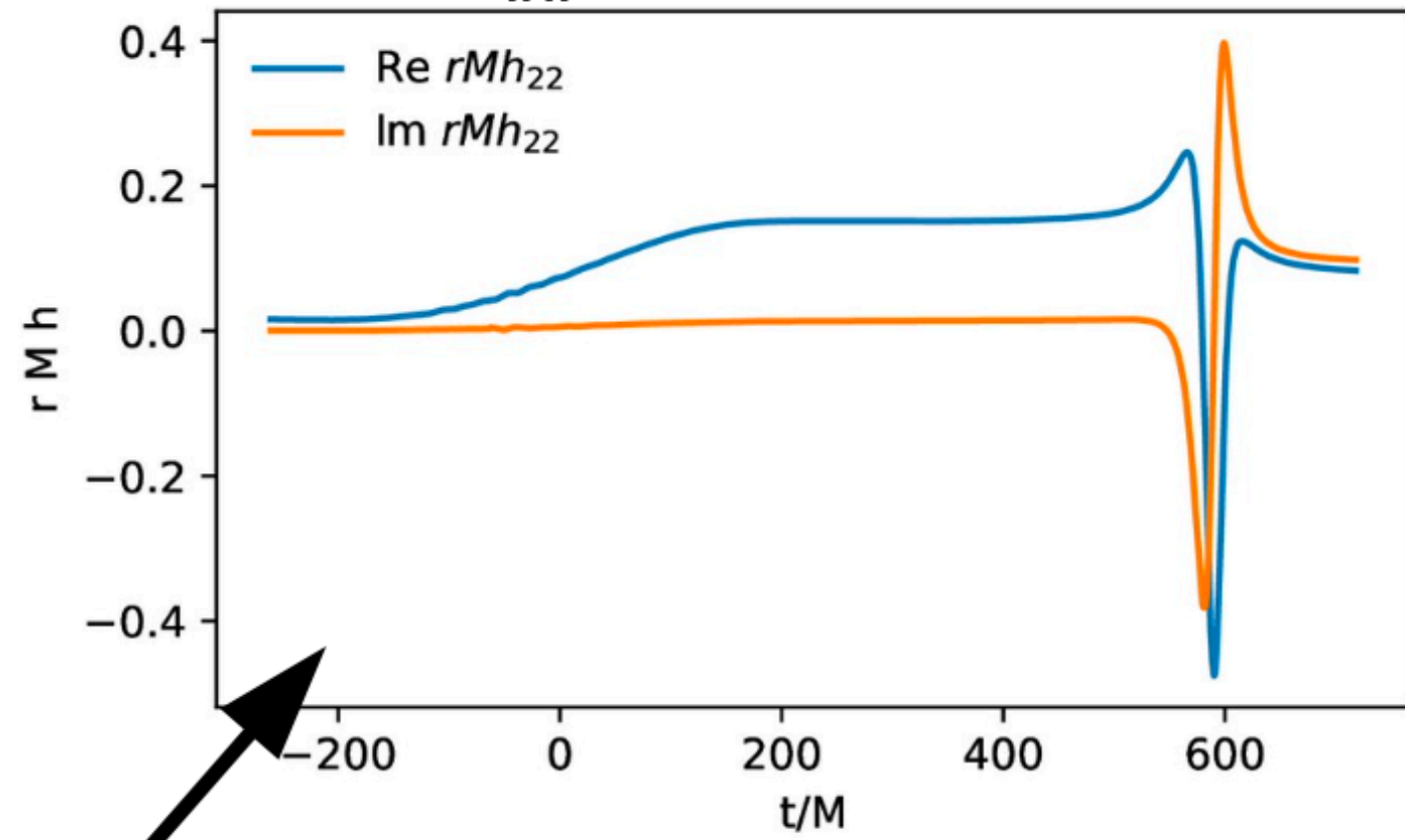


Some first waveforms

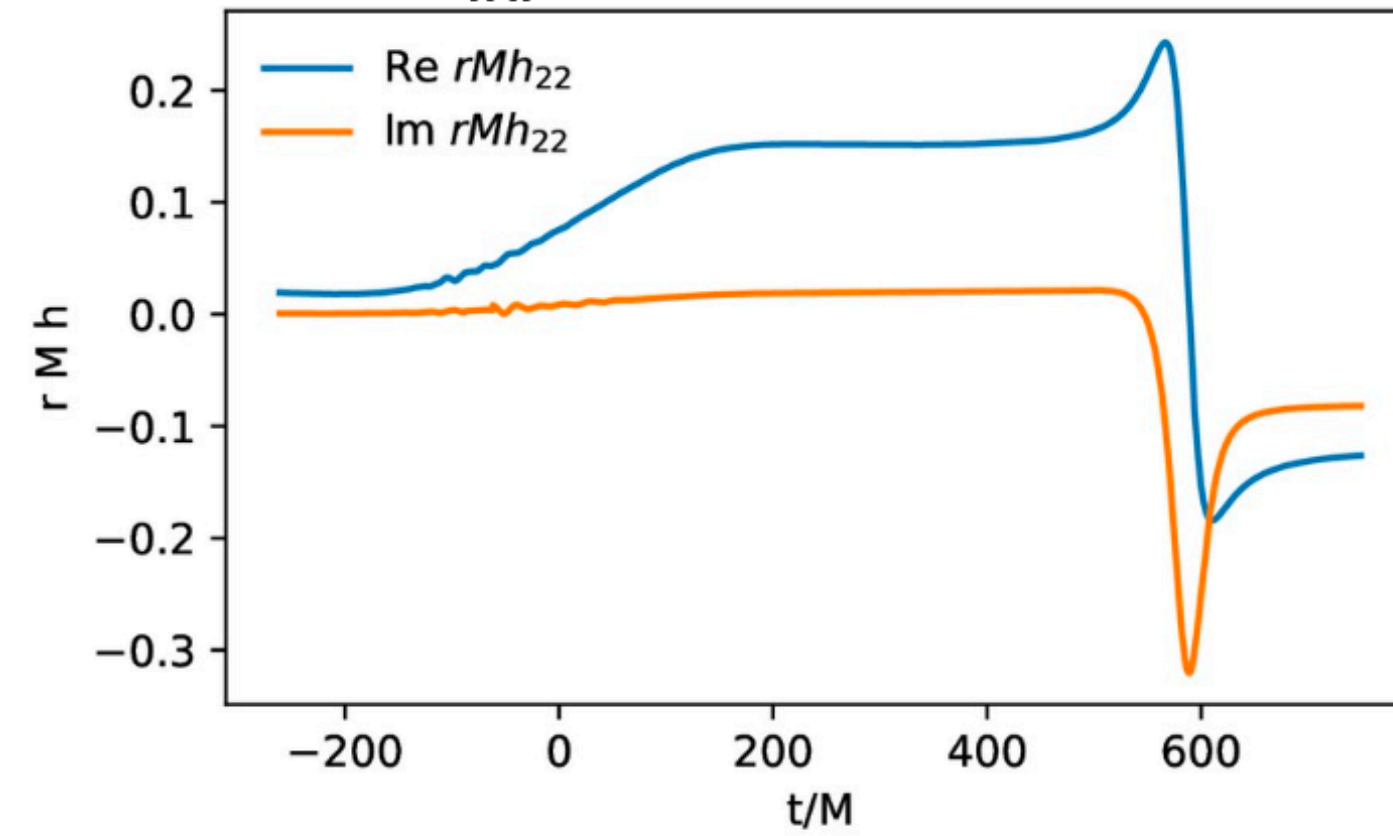


$E=0.0223 M$

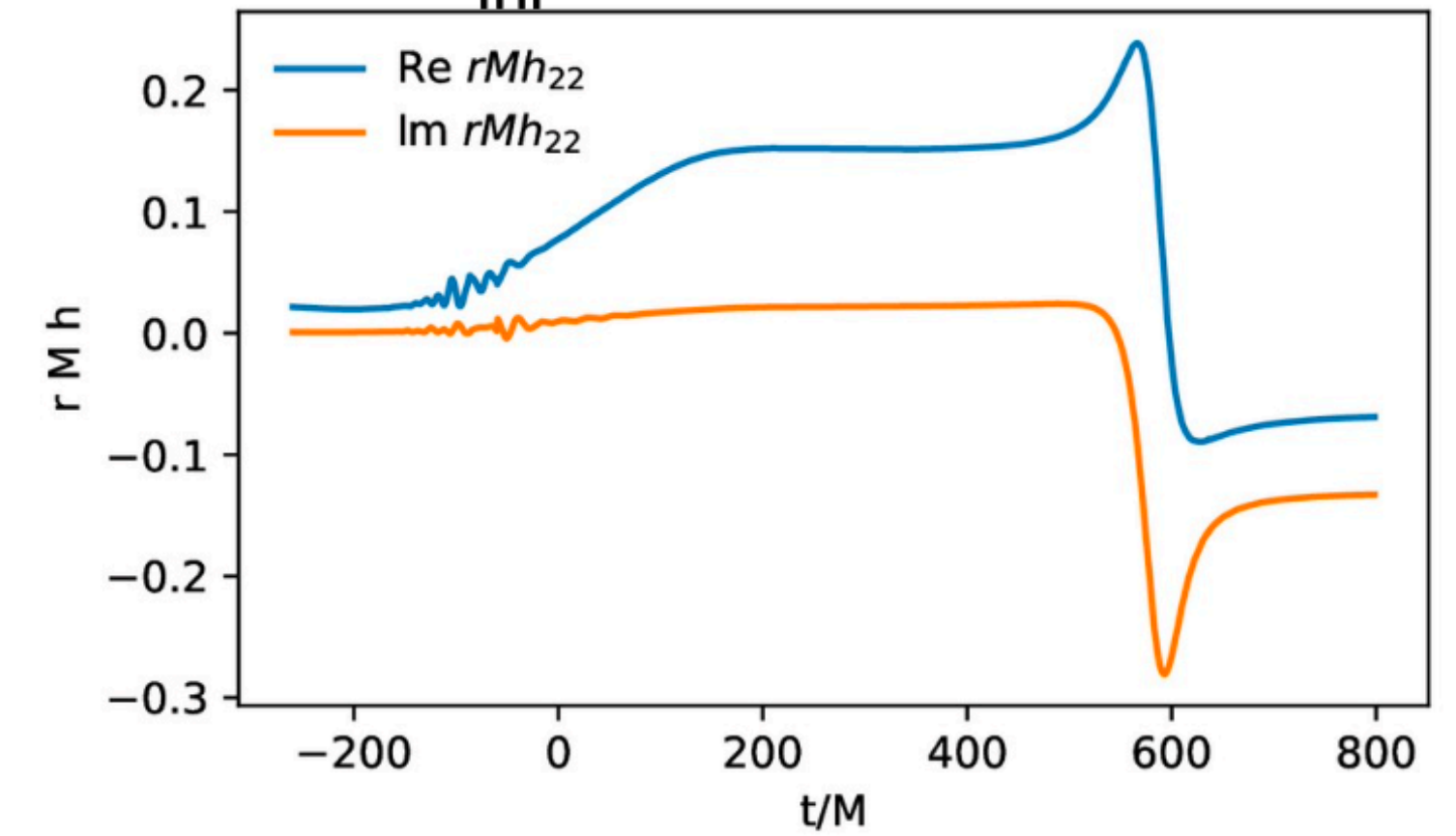
$J_{ini} = 1.103 M^2$



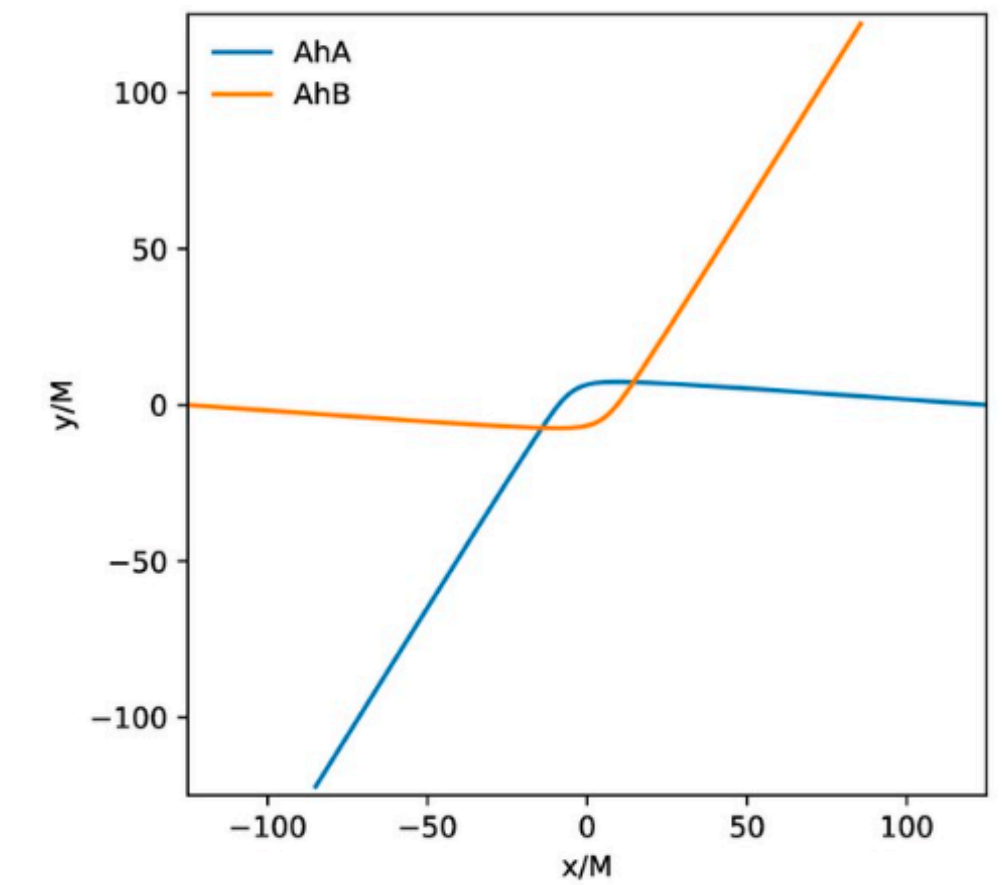
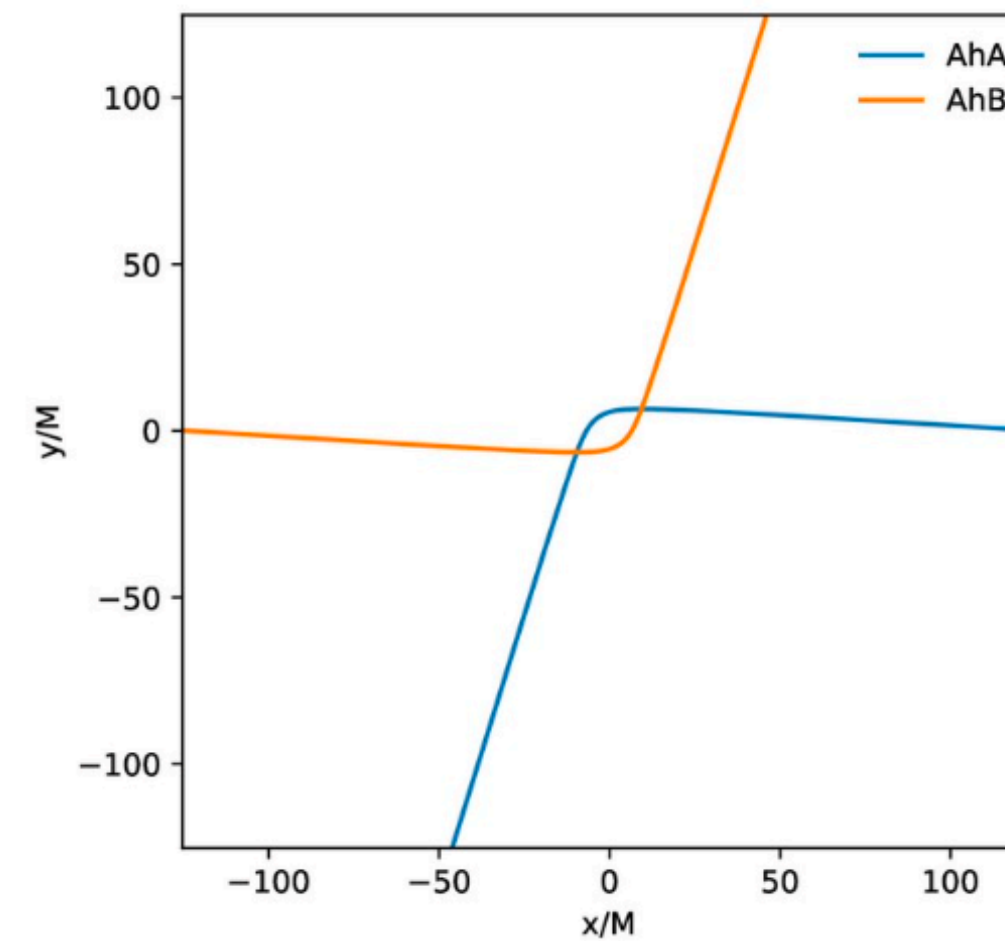
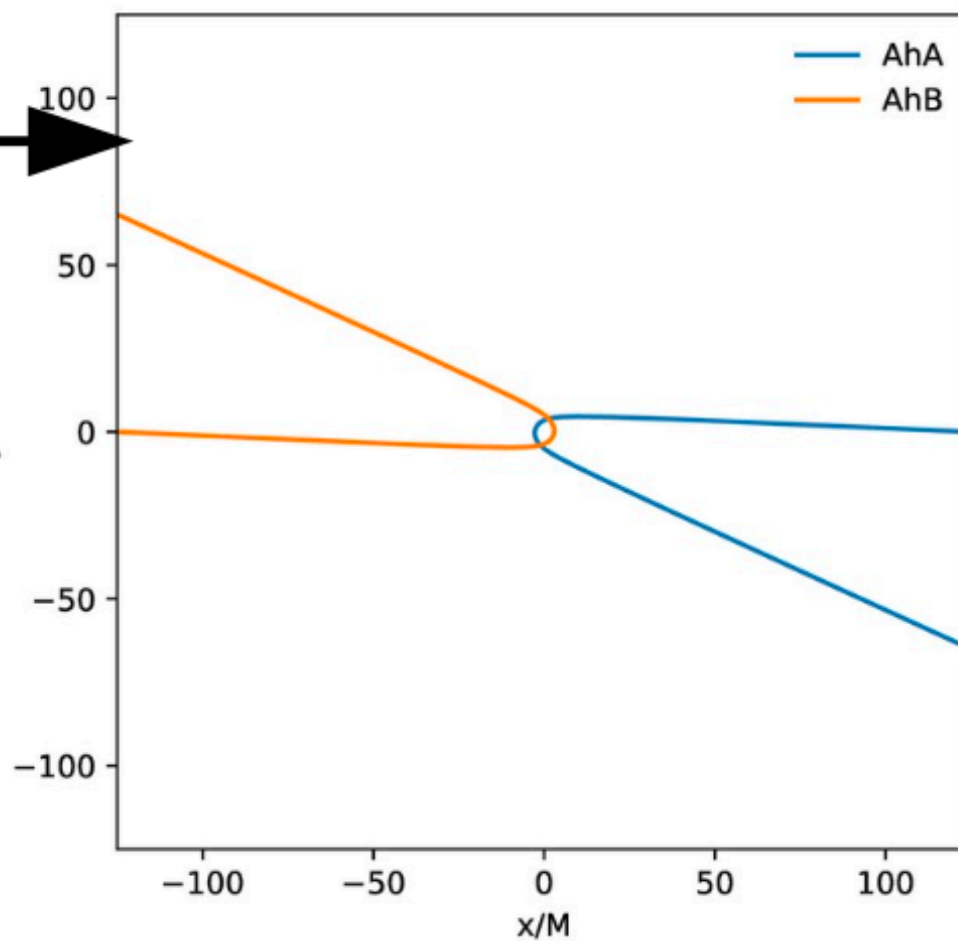
$J_{ini} = 1.558 M^2$



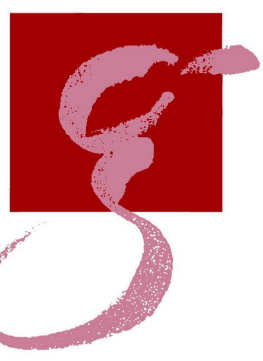
$J_{ini} = 1.780 M^2$



Capture,
but huge
semi-
major
axis



Rüter, HP, SXS in prep



- **NR for quasi-circular binaries mature pillar of GW astronomy**
 - must keep up with improving detectors!
- **High mass-ratio challenging**
 - new methods promising
 - growing evidence that **2nd order SMR may be surprisingly accurate**
- **Eccentric and hyperbolic systems** gain in attention
 - super-large parameter spaces

