

Reading the footprints of the flavor anomalies

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The Standard Model (SM)

Extremely successful description of Nature...

... yet many unanswered questions

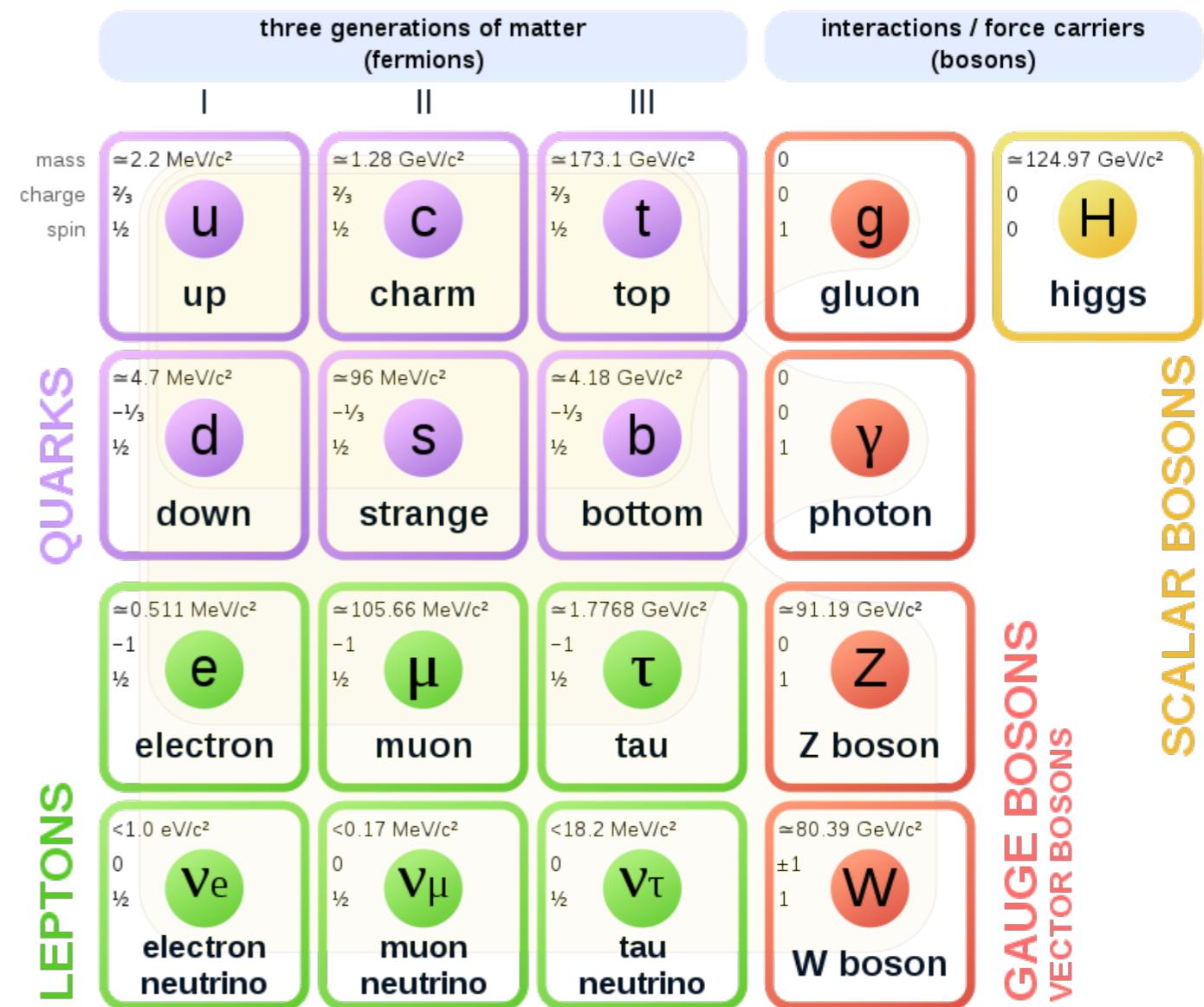
- Neutrino masses
- Dark matter
- Naturalness problems
- ...

What we know after LHC runs I & II:

- ★ The Higgs boson is SM-like and relatively light [**SM is complete!**]
- ★ No direct evidence for new particles [**presence of a mass gap?**]

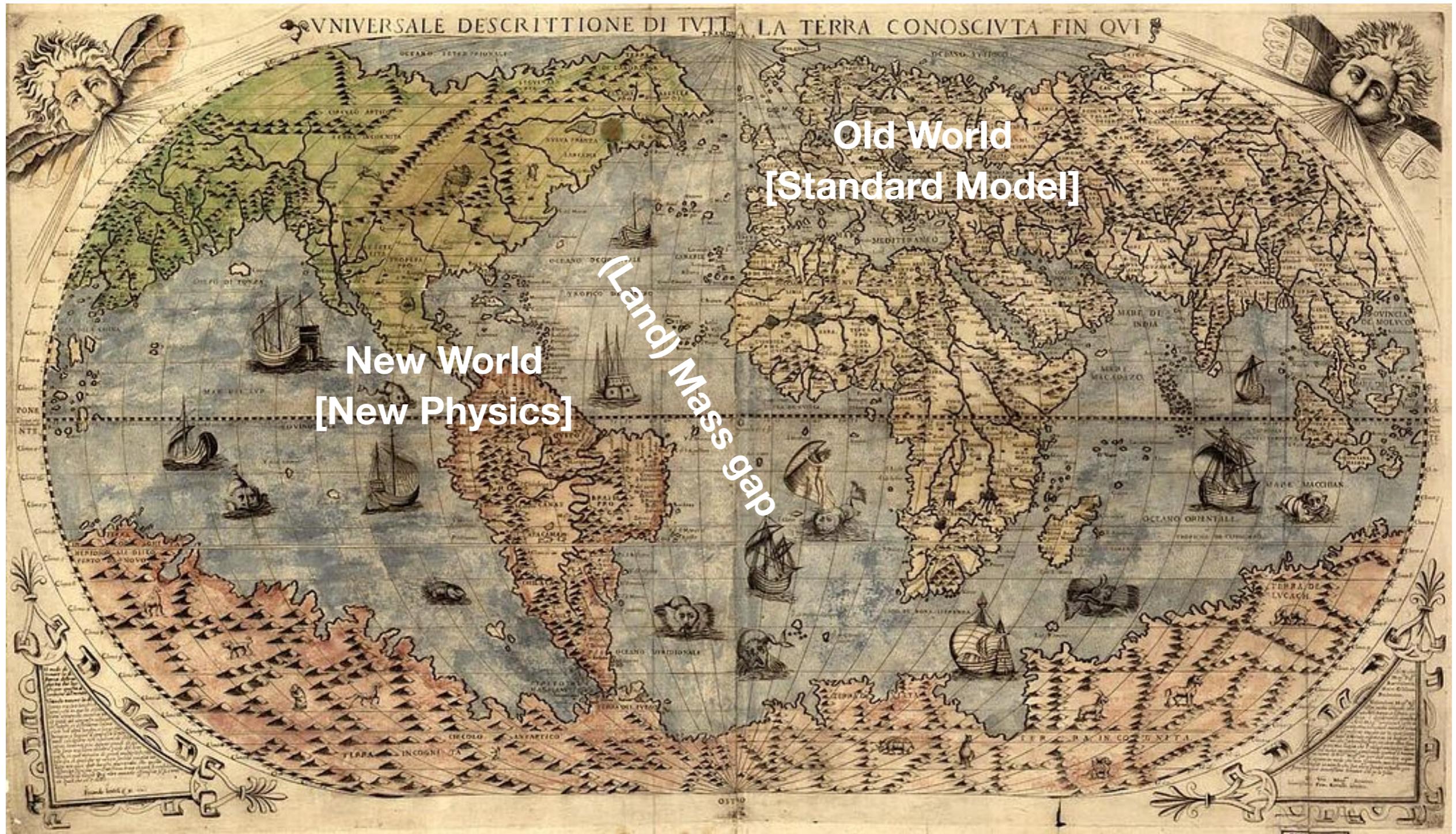
$$\mathcal{G}_{\text{SM}} \equiv SU(3)_c \times SU(2)_L \times U(1)_Y$$

Standard Model of Elementary Particles



The search for Terra Incognita

Much like late Middle Age Europe, particle physics has entered an age of exploration

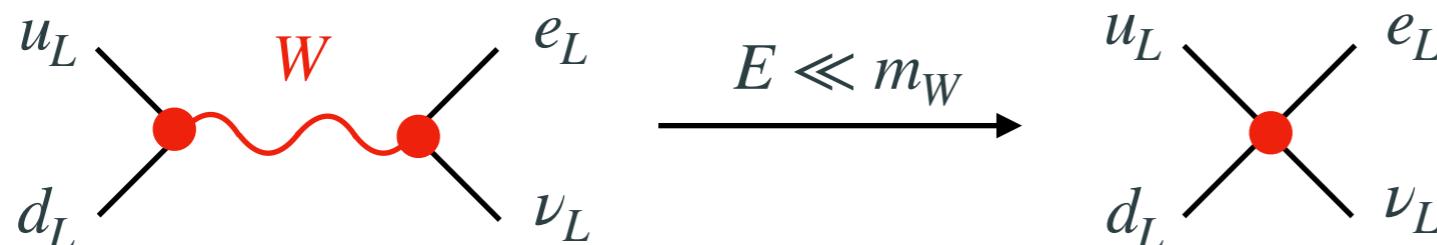


The Fermi theory

Not the first time we have faced a mass gap in particle physics

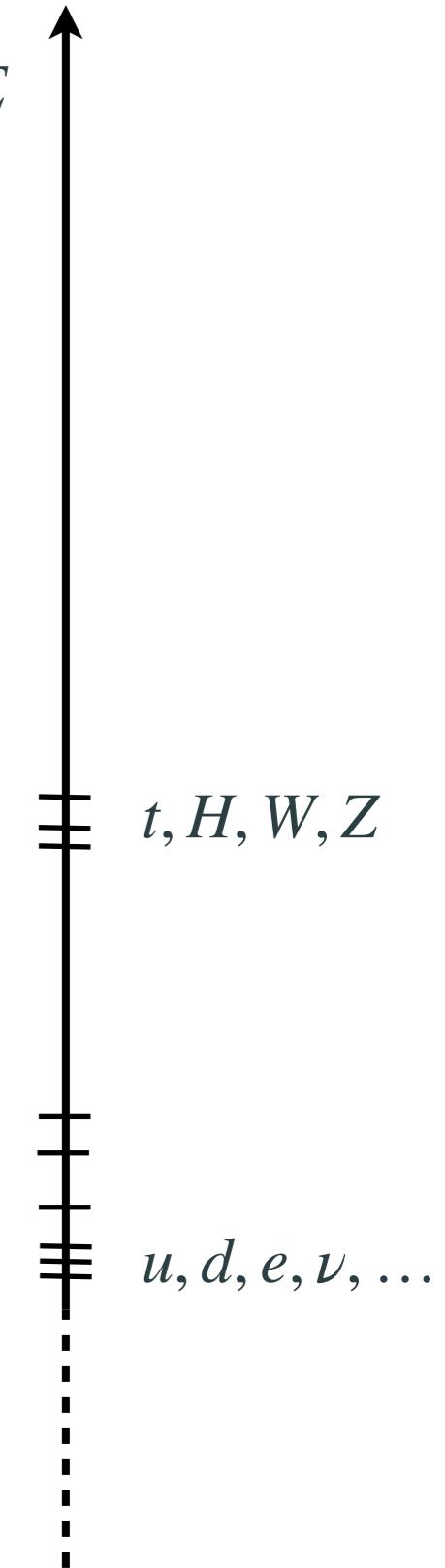
Fermi Theory [$E \ll m_W$]

$$\mathcal{L}_{WET} = \mathcal{L}_{QED} + \mathcal{L}_{QCD} - \frac{4G_F}{\sqrt{2}} (\bar{u}_L \gamma_\mu d_L)(\bar{e}_L \gamma^\mu \nu_L) + \dots \quad [G_F \sim g_W^2/M_W^2]$$



Reconstructing a UV theory from its low-energy imprints is a **very difficult task** (no unique solution due to limited information)

[It took **more than 30 years** to arrive to the SM from the Fermi theory]



The SM effective theory (SMEFT)

If New Physics is heavy (mass gap), we can write the analog of the Fermi theory for the SM

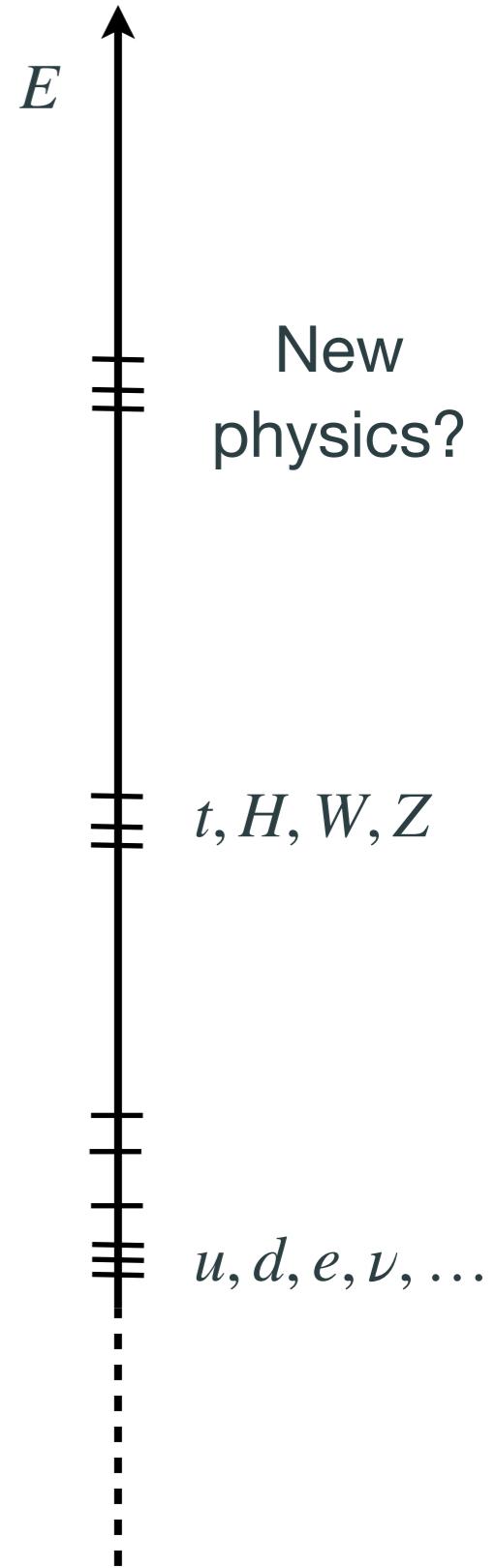
SMEFT [$E \ll M_{\text{NP}}$]

$$\mathcal{L}_{\text{SMEFT}} = \mathcal{L}_{\text{SM}} + C_{\ell q}^{\alpha\beta ij} (\ell_L^\alpha \gamma_\mu \ell_L^\beta)(q_L^i \gamma^\mu q_L^j) + \dots \quad [C_{\ell q} \sim \Lambda^{-2} \sim g_{\text{NP}}^2/m_{\text{NP}}^2]$$



- ★ 59 new possible interactions (2499 new couplings) at $\mathcal{O}(\Lambda^{-2})$
- ★ New physics is unlikely to produce them all with the same strength

Can we infer anything about them from the SM couplings?



The SM Lagrangian: Naturalness problems

$$\begin{aligned} \mathcal{L} = & -\frac{1}{4} F_{\mu\nu} F^{\mu\nu} \\ & + i \bar{\psi} \not{D} \psi \\ & + \frac{1}{2} |\partial_\mu \phi|^2 - V(\phi) \\ & + \bar{\psi}_i y_{ij} \psi_j \phi + h.c. \end{aligned}$$

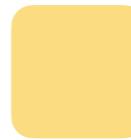
The SM Lagrangian contain two **unnatural features** pointing towards new physics



Higgs hierarchy problem

[Instability of the Higgs mass under quantum corrections]

TeV-scale new physics?



SM flavor puzzle

[Very hierarchical structure in the Yukawa couplings]

Similar structure also for new physics?

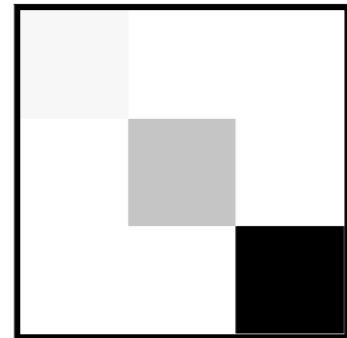
Are these two features correlated?

The SM flavor puzzle

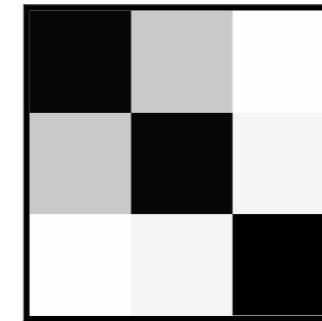
The SM Yukawa sector is characterized by **13** parameters (for massless neutrinos)
[**3** lepton masses + **6** quark masses + **3+1** CKM parameters]

... whose values span **5 orders of magnitude** and do not look **at all accidental**

$$M_{u,d,e} \sim$$



$$V_{\text{CKM}} \sim$$



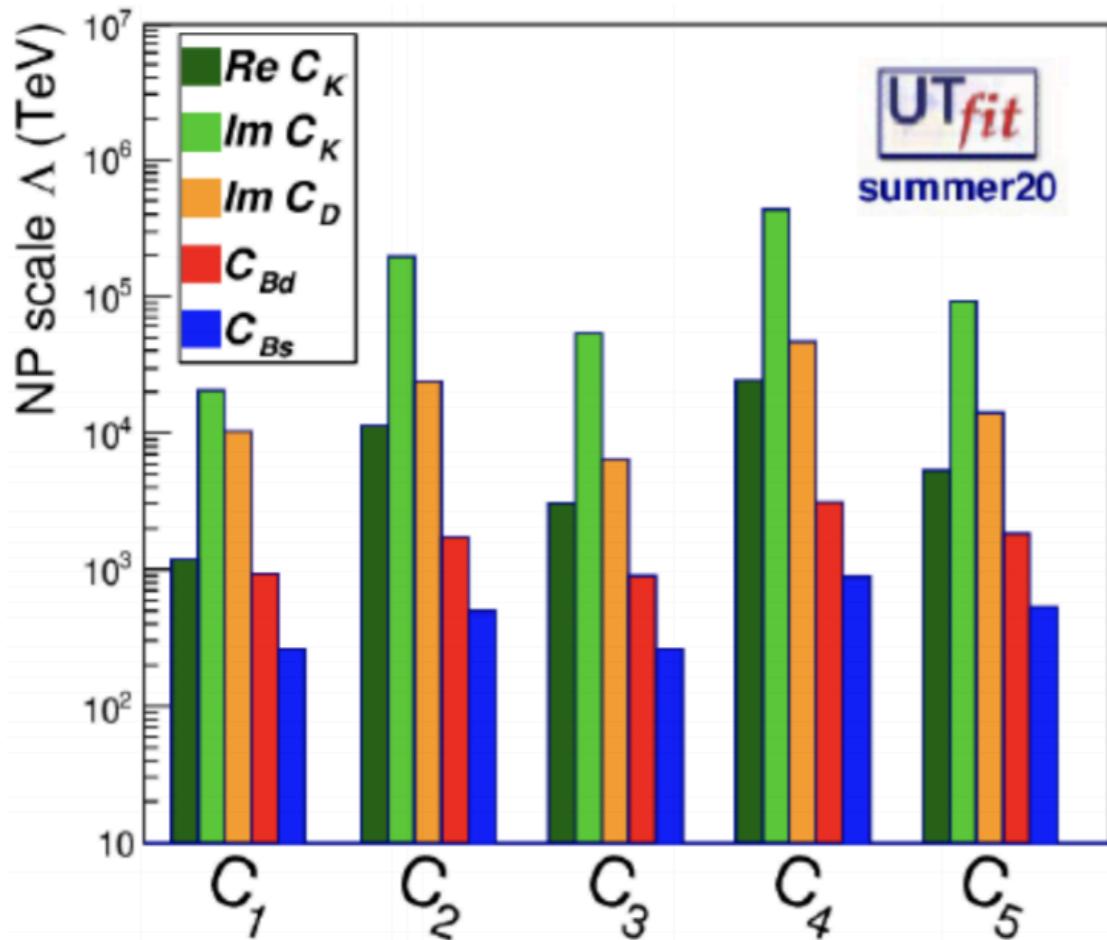
Possible explanations/approaches to the observed hierarchies:

1. They are accidental and/or arise from very high-scale new physics out of our current reach
2. They are low-energy imprints of new symmetries/dynamics deeply connected to the next layer of new physics (possibly related to the Higgs hierarchy problem)

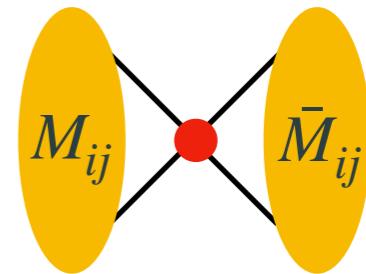
The new physics flavor problem

$$\mathcal{L}_{\text{SMEFT}} = \mathcal{L}_{\text{Gauge}} + \mathcal{L}_{\text{Higgs}} + \mathcal{L}_{\text{Yukawa}} + \sum_{i,d} \frac{1}{\Lambda_i^{d-4}} C_i \mathcal{O}_i^d$$

Very stringent bounds on the new physics scale if it has a **generic flavor structure**
(far too heavy to be directly probed or to stabilize the Higgs)



$$\frac{1}{\Lambda^2} (\bar{\psi}_i \psi_j)^2$$



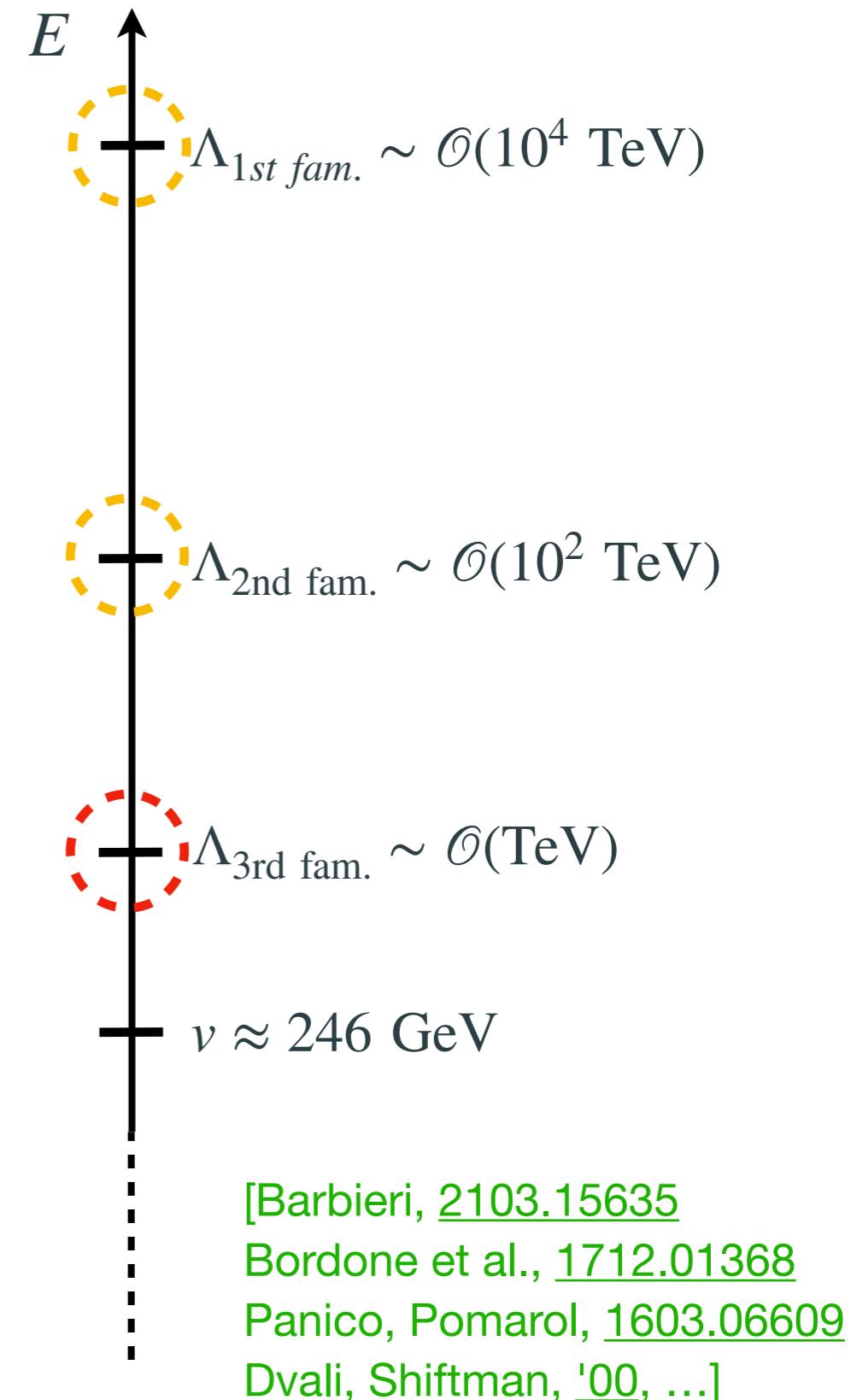
1. **Minimal flavor violation:** SM Yukawas are the only source of flavor violation [new physics is flavor blind/universal]
[D'Ambrosio, Giudice, Isidori, Strumia, ['02](#)]
2. **New physics is flavor non-universal** and possibly connected to the origin of the Yukawa hierarchies

Multi-scale solution of the flavor problem/puzzle

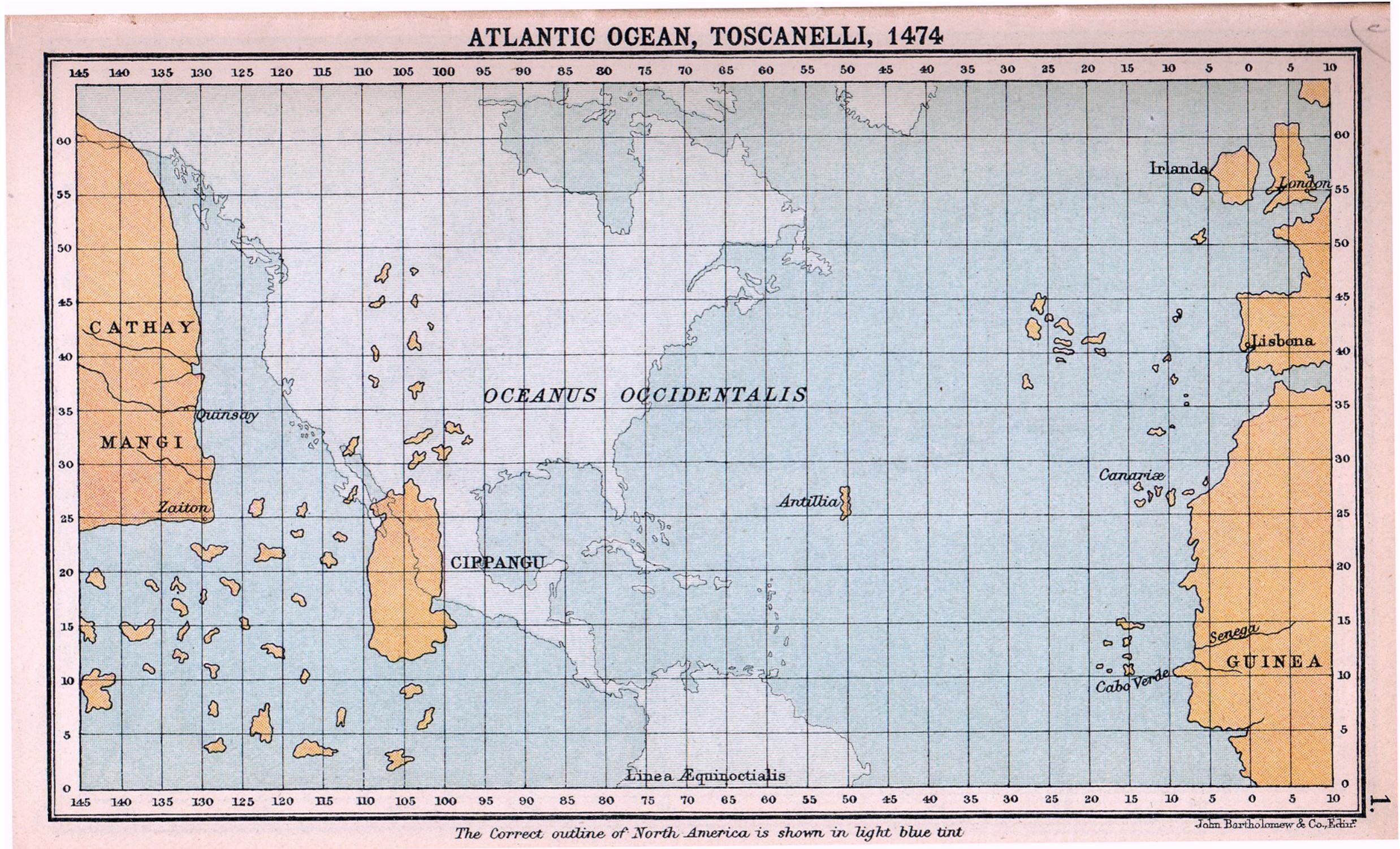
$$\mathcal{L}_{\text{SMEFT}} = \mathcal{L}_{\text{Gauge}} + \mathcal{L}_{\text{Higgs}} + \mathcal{L}_{\text{Yukawa}} + \sum_{i,d} \frac{1}{\Lambda_i^{d-4}} C_i \mathcal{O}_i^d$$

Non-trivial UV imprints

- ★ The SM Yukawas are very different because they originate at separate scales!
- ★ TeV-scale new physics dominantly coupled to the third and (to a lesser extent) second families
- ★ Direct production of new physics states at the LHC is naturally more suppressed [new physics scale can be lower]



A closer look to the data and EFT analysis

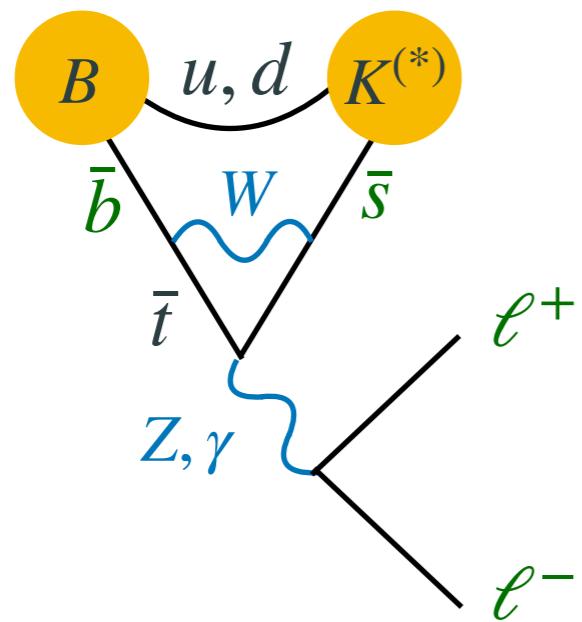


The B anomalies

Hints of Lepton Flavour Universality Violation (LFUV) in semileptonic B decays

$$b \rightarrow s \ell^+ \ell^-$$

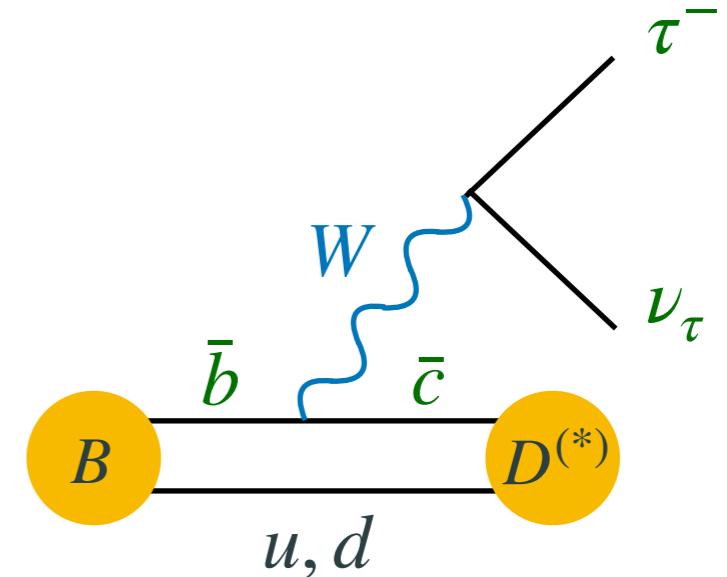
μ/e universality



$> 4.5 \sigma$

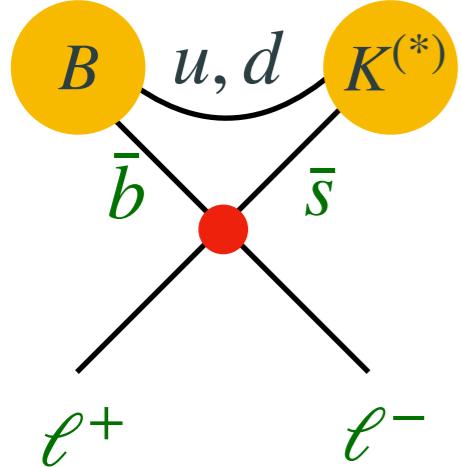
$$b \rightarrow c \tau \nu$$

$\tau/\mu, e$ universality



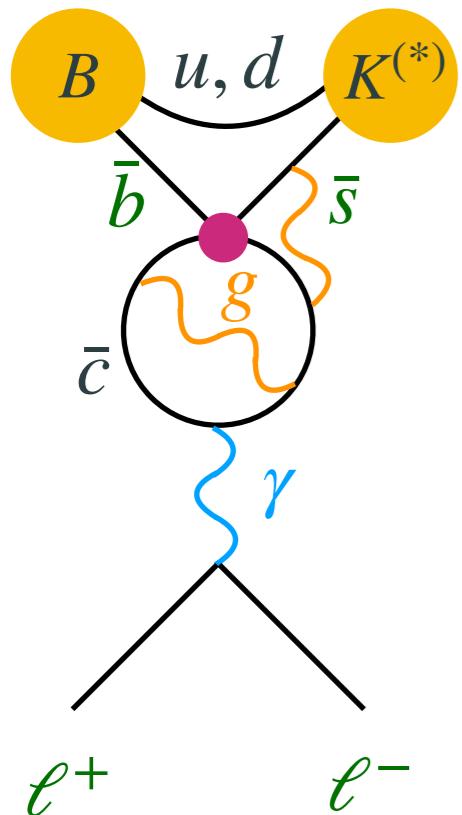
$\sim 3 \sigma$

Anatomy of $b \rightarrow s\ell^+\ell^-$ decays



Short-distance
(semileptonic int.)

“Easy” to compute



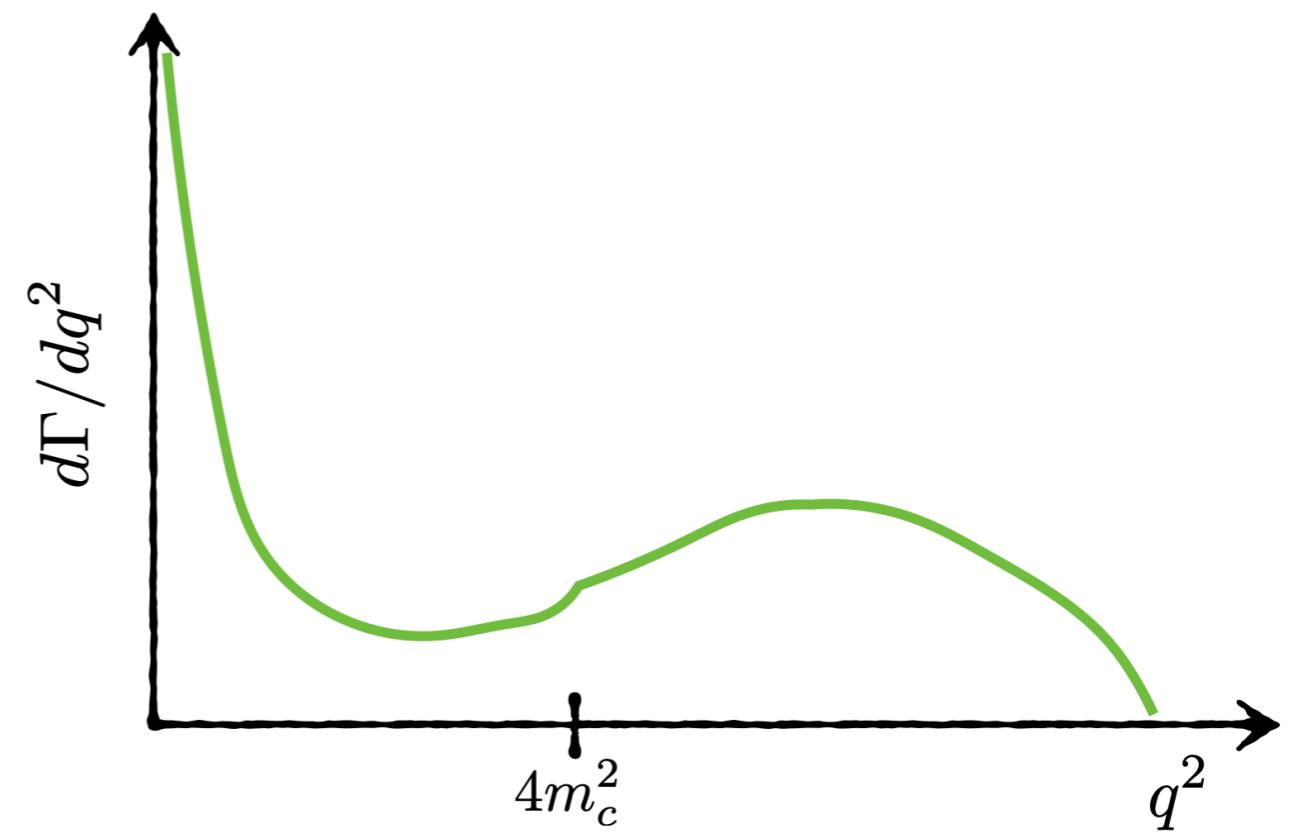
Long-distance
(four-quark int.)

Difficult to estimate

Induces a vectorial
and lepton-universal
contribution

$$\mathcal{O}_9^\ell = (\bar{s}_L \gamma_\mu b_L)(\bar{\ell} \gamma^\mu \ell)$$

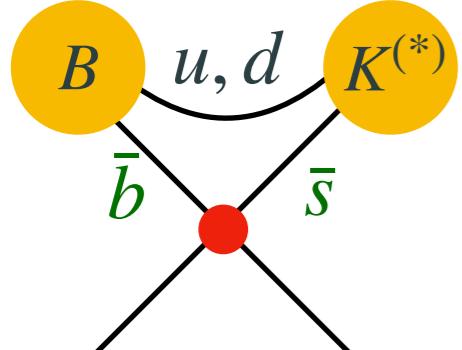
In an ideal world



[Figure from Uli Haisch]

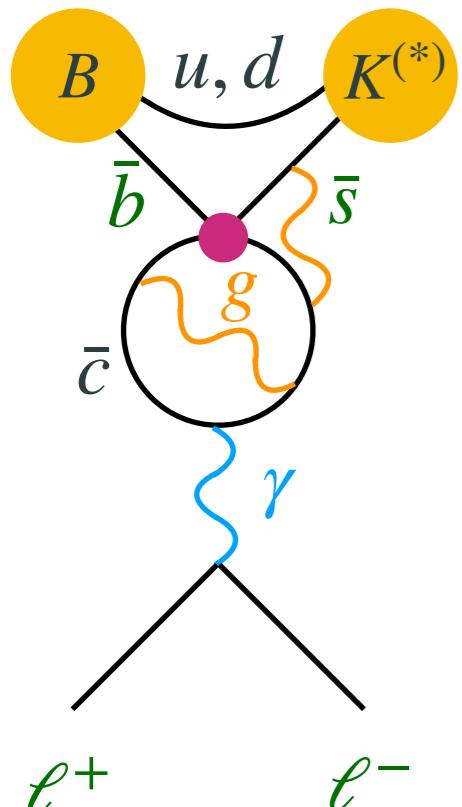
- ★ Long-distance effects cannot induce:
 - Breakings of lepton universality
 - Axial-current contributions
(no effect in $B_s \rightarrow \mu^+ \mu^-$)

Anatomy of $b \rightarrow s\ell^+\ell^-$ decays



Short-distance
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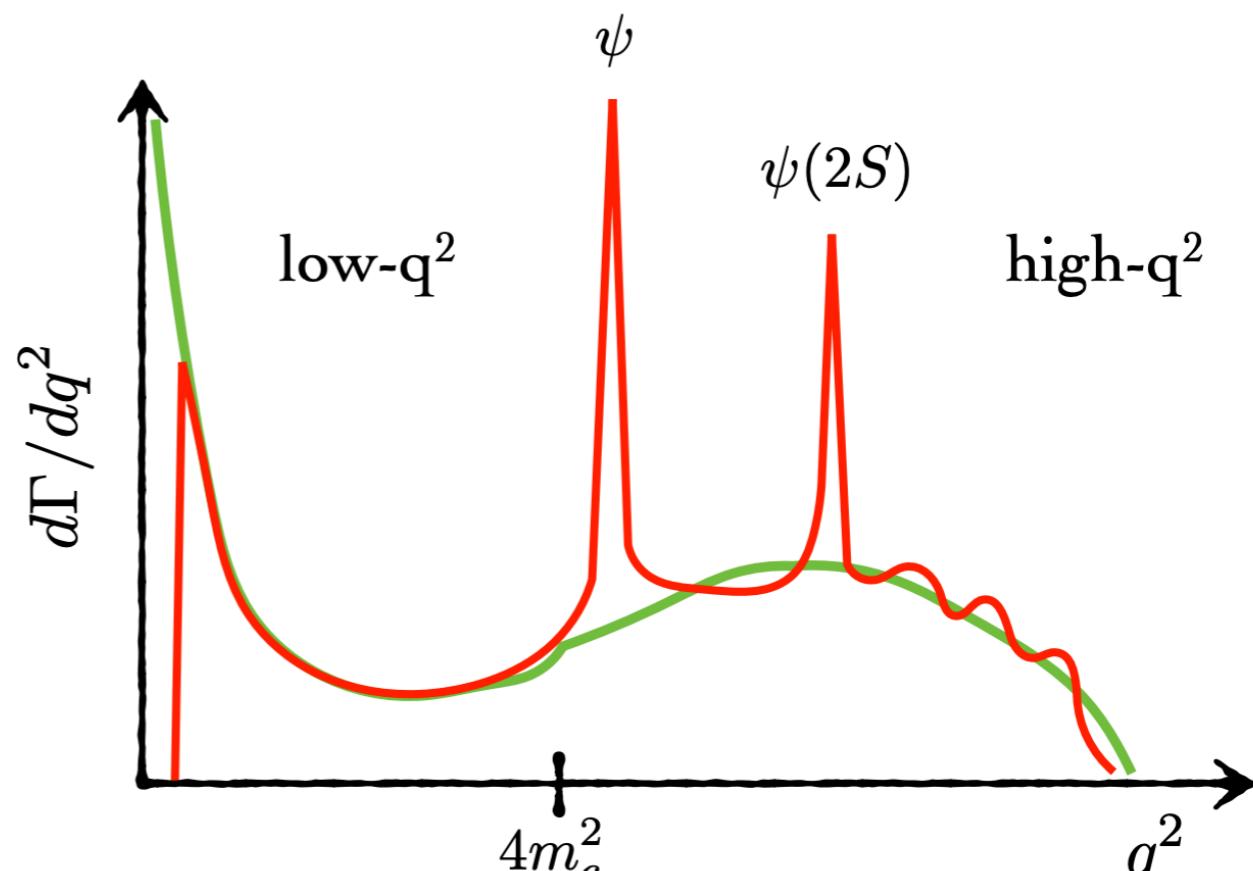
Long-distance
(four-quark int.)

Difficult to estimate

Induces a vectorial
and lepton-universal
contribution

$$\mathcal{O}_9^\ell = (\bar{s}_L \gamma_\mu b_L)(\bar{\ell} \gamma^\mu \ell)$$

... but in reality



[Figure from Uli Haisch]

- ★ Long-distance effects cannot induce:
 - Breakings of lepton universality
 - Axial-current contributions
(no effect in $B_s \rightarrow \mu^+ \mu^-$)

The $b \rightarrow s\ell^+\ell^-$ anomalies

Since 2013, several “anomalies” started to appear in B -meson decays involving the $b \rightarrow s\ell^+\ell^-$ [$\ell = e, \mu$] semileptonic transition [**plenty of updates this year!**]

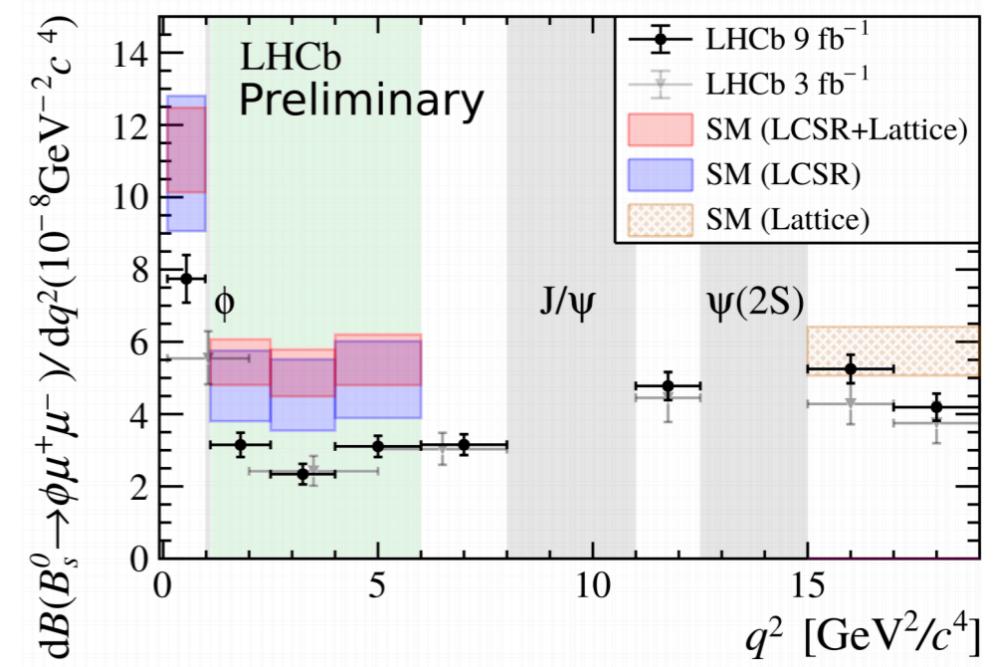
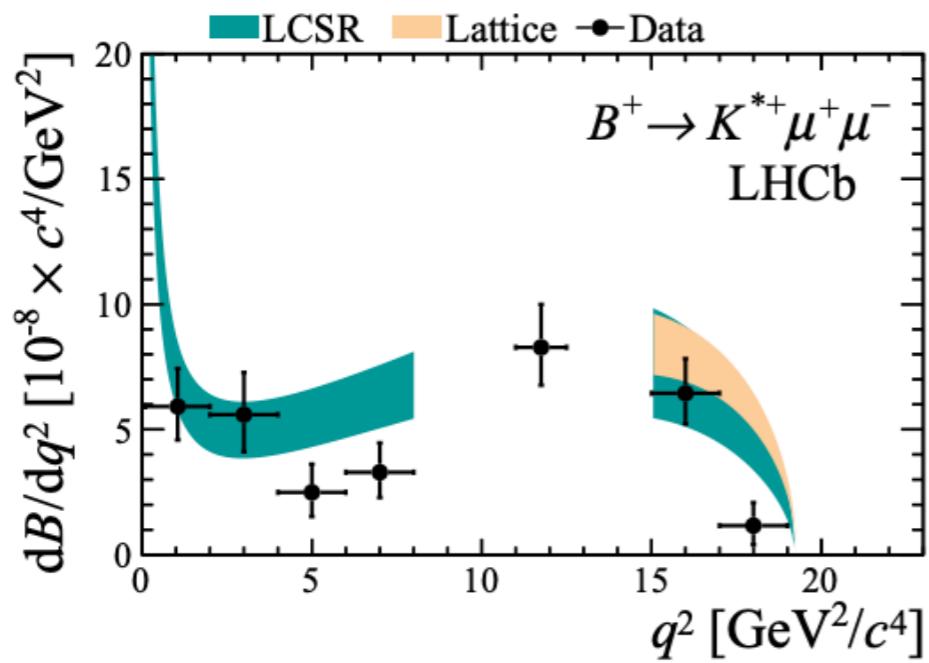
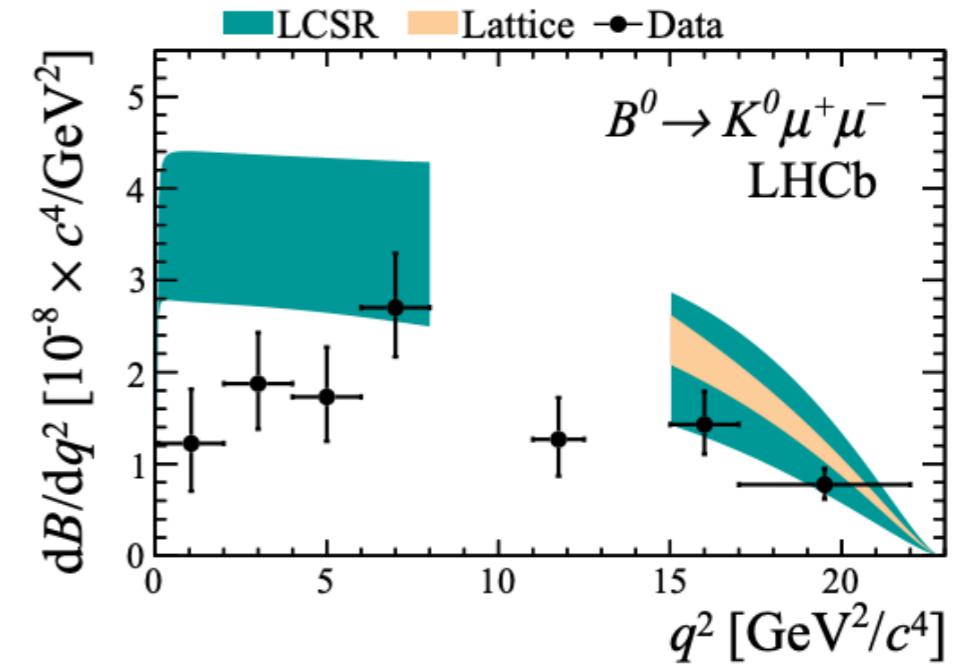
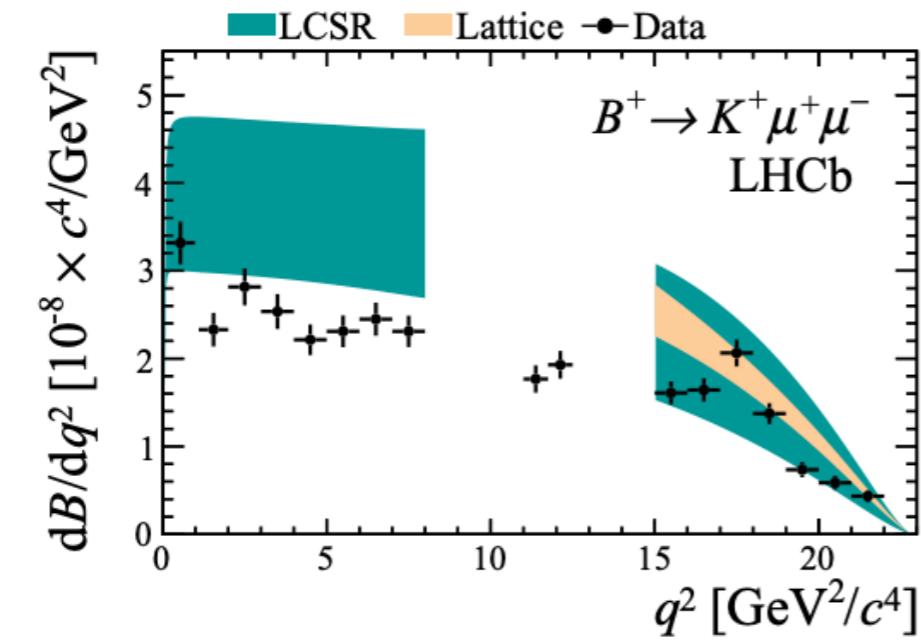
- ★ Deficit in the exclusive $B \rightarrow K^{(*)}\mu^+\mu^-$ and $B_s \rightarrow \phi\mu^+\mu^-$ decay rates
[sizable QCD uncertainties (form factors and long-distance contributions)]
- ★ Deviations in optimized angular observables in $B \rightarrow K^*\mu^+\mu^-$ (P'_5 anomaly)
[reduced theory errors but still sizable QCD uncertainties]
- ★ Deficit in the leptonic $B_s \rightarrow \mu^+\mu^-$ decay rate
[theoretically “clean”]: theory error below 5 % (no charm loop)
- ★ Departures from unity in lepton flavor universality ratios with $B \rightarrow K^{(*)}\ell^+\ell^-$
[theoretically very “clean”]: 1 % theory error (QED and lepton mass effects)]

Theoretical robustness

Remarkably, all this data point to **coherent NP effects**

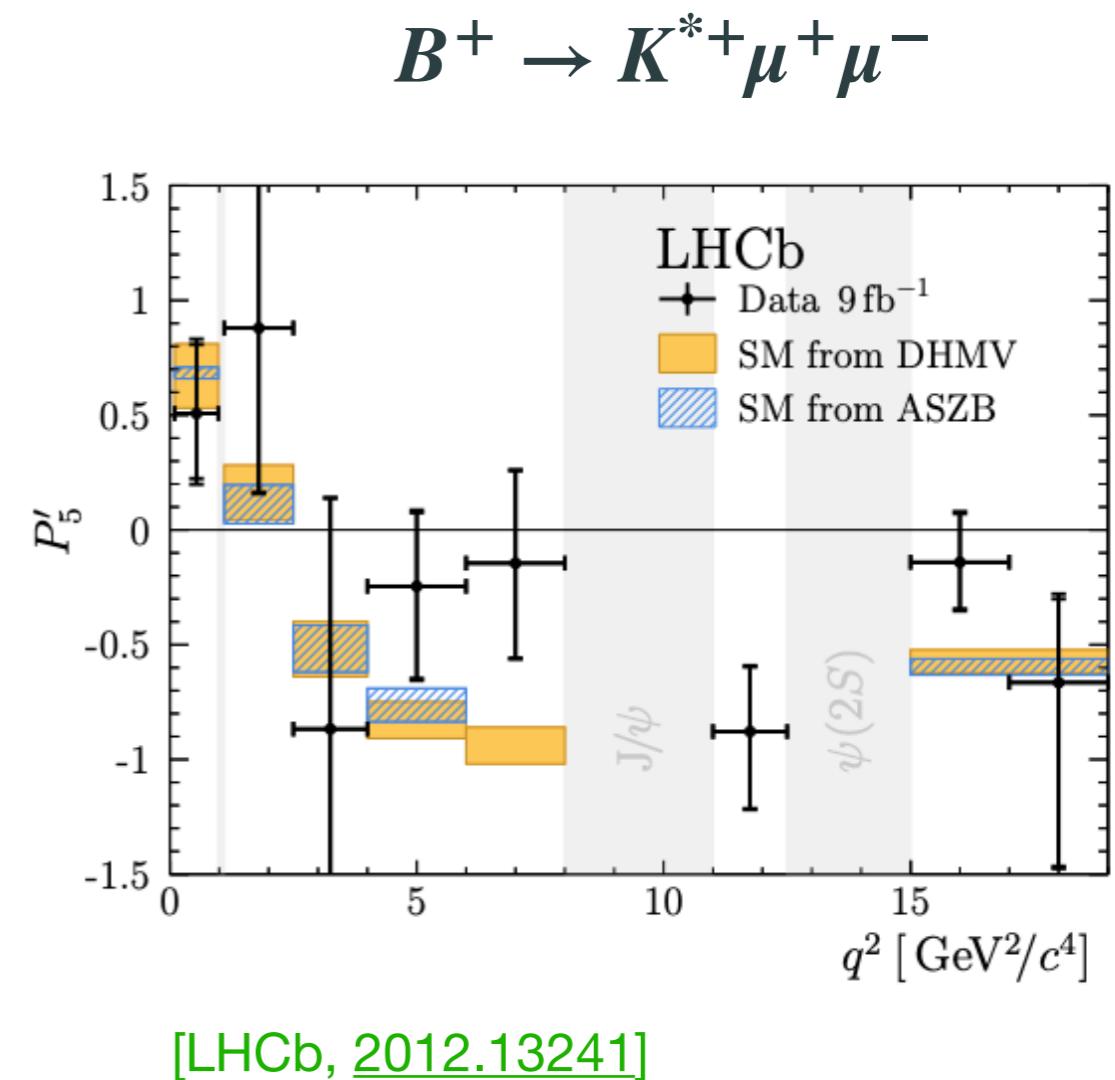
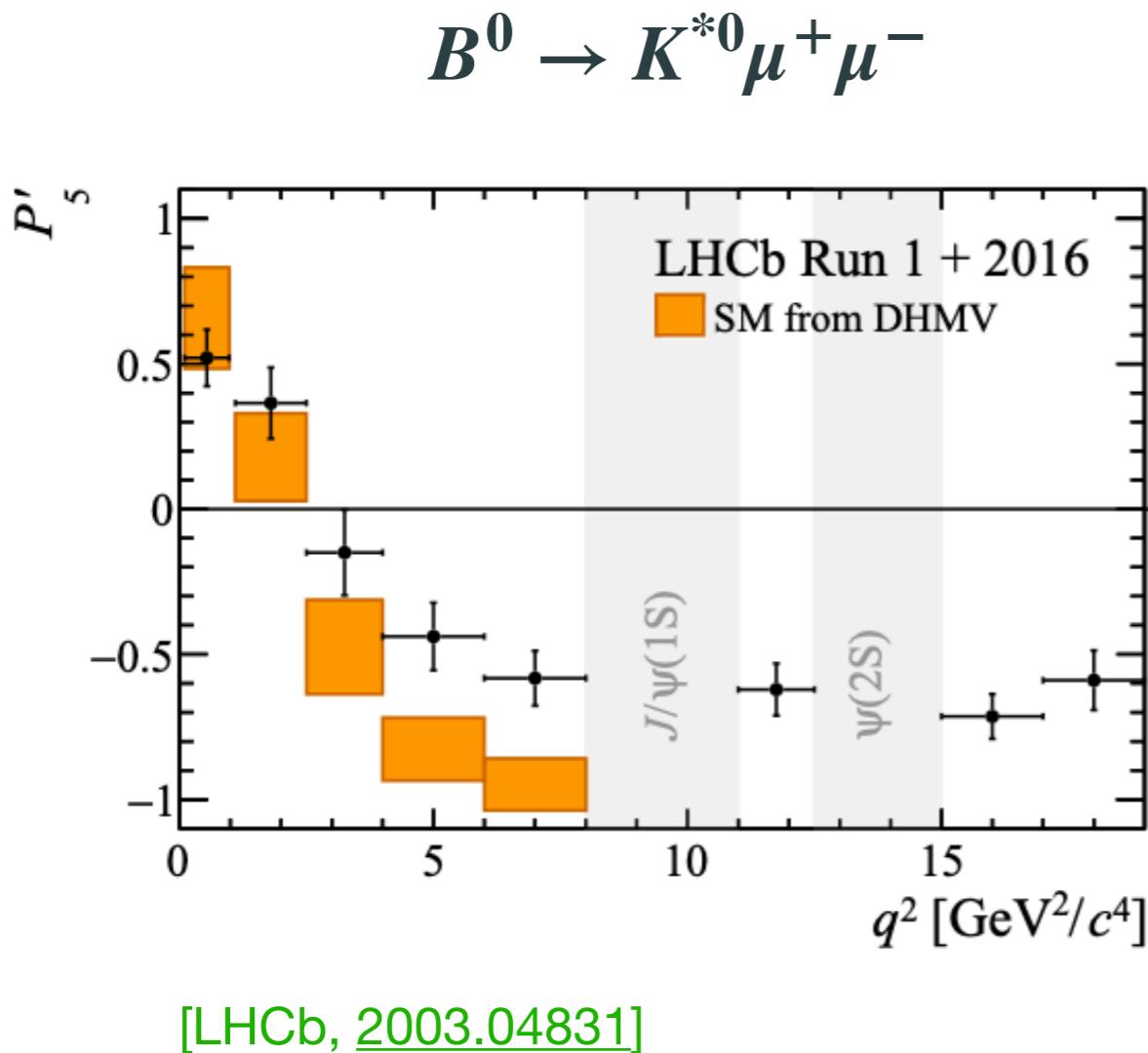
The $b \rightarrow s\ell^+\ell^-$ anomalies

- ★ Deficit in the exclusive $B \rightarrow K^{(*)}\mu^+\mu^-$ and $B_s \rightarrow \phi\mu^+\mu^-$ decay rates
[sizable QCD uncertainties from factors and long-distance contributions]



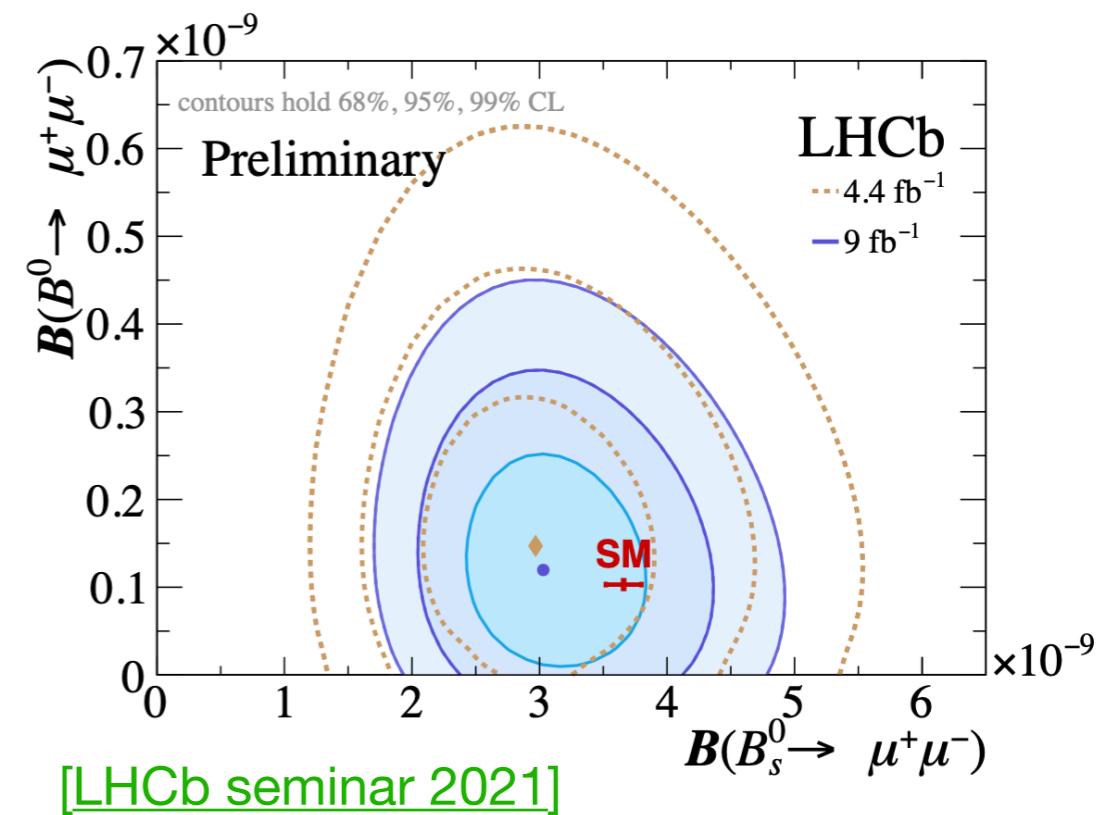
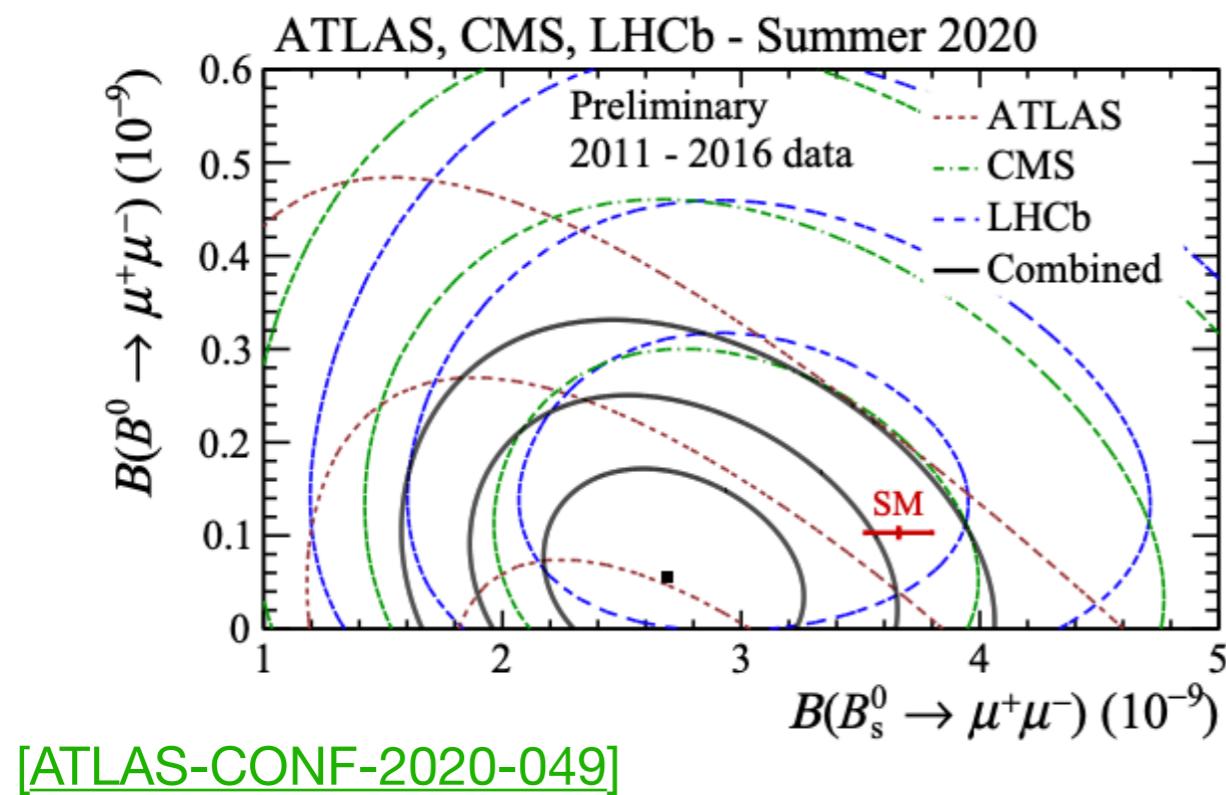
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The $b \rightarrow s\ell^+\ell^-$ anomalies

- ★ Deficit in the leptonic $B_s \rightarrow \mu^+\mu^-$ decay rate
 - [theoretically “clean”: theory error below 5 % (no charm loop)]



$$\mathcal{B}(B_s \rightarrow \mu^+\mu^-)_{\text{exp}} = (2.85^{+0.32}_{-0.31}) \times 10^{-9}$$

$$\mathcal{B}(B_s \rightarrow \mu^+\mu^-)_{\text{SM}} = (3.66 \pm 0.14) \times 10^{-9}$$

2.3σ deviation from the SM prediction

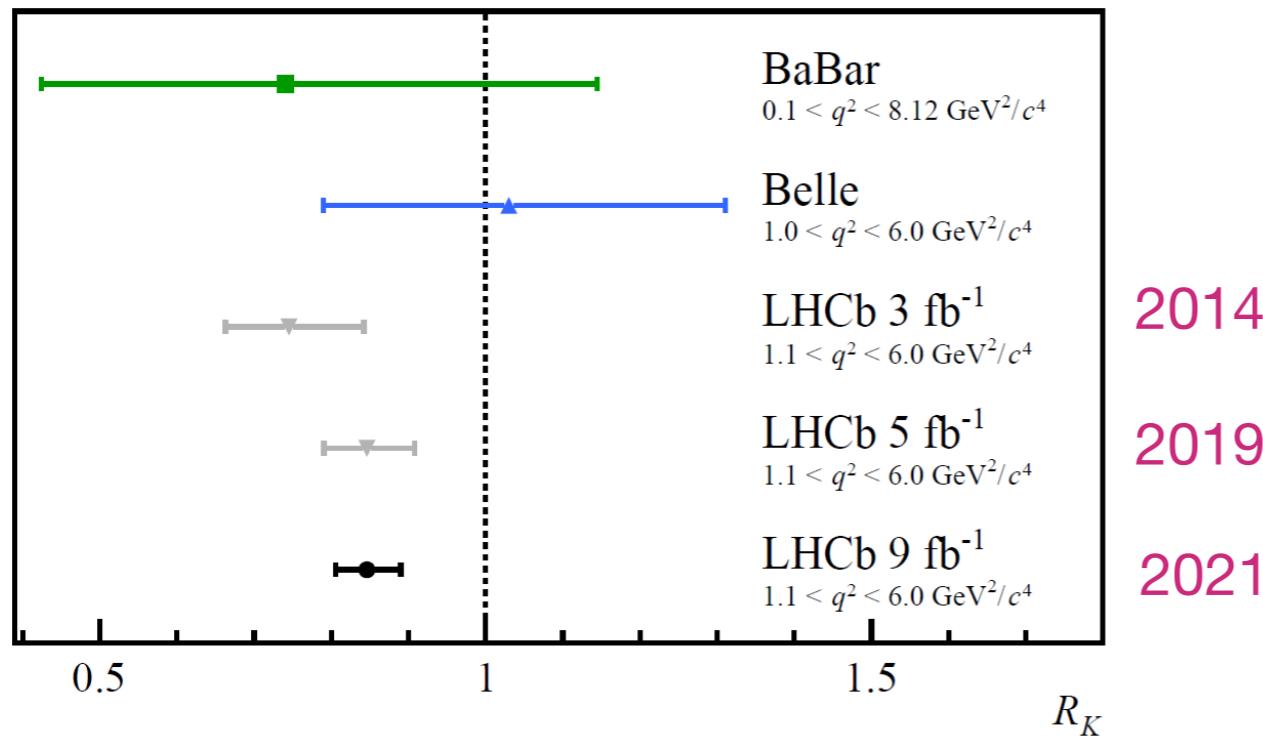
The $b \rightarrow s\ell^+\ell^-$ anomalies

- ★ Departures from unity in lepton flavor universality ratios with $B \rightarrow K^{(*)}\ell^+\ell^-$
[theoretically very “clean”: 1 % theory error (QED and lepton mass effects)]

$$R_{K^{(*)}}^{[q_{\min}^2, q_{\max}^2]} \equiv \frac{\int_{q_{\min}^2}^{q_{\max}^2} d\Gamma(B \rightarrow K^{(*)}\mu^+\mu^-)}{\int_{q_{\min}^2}^{q_{\max}^2} d\Gamma(B \rightarrow K^{(*)}e^+e^-)}$$

$$R_{K^{(*)}}^{[1.1, 6] \text{ GeV}^2} = 1.00 \pm 0.01$$

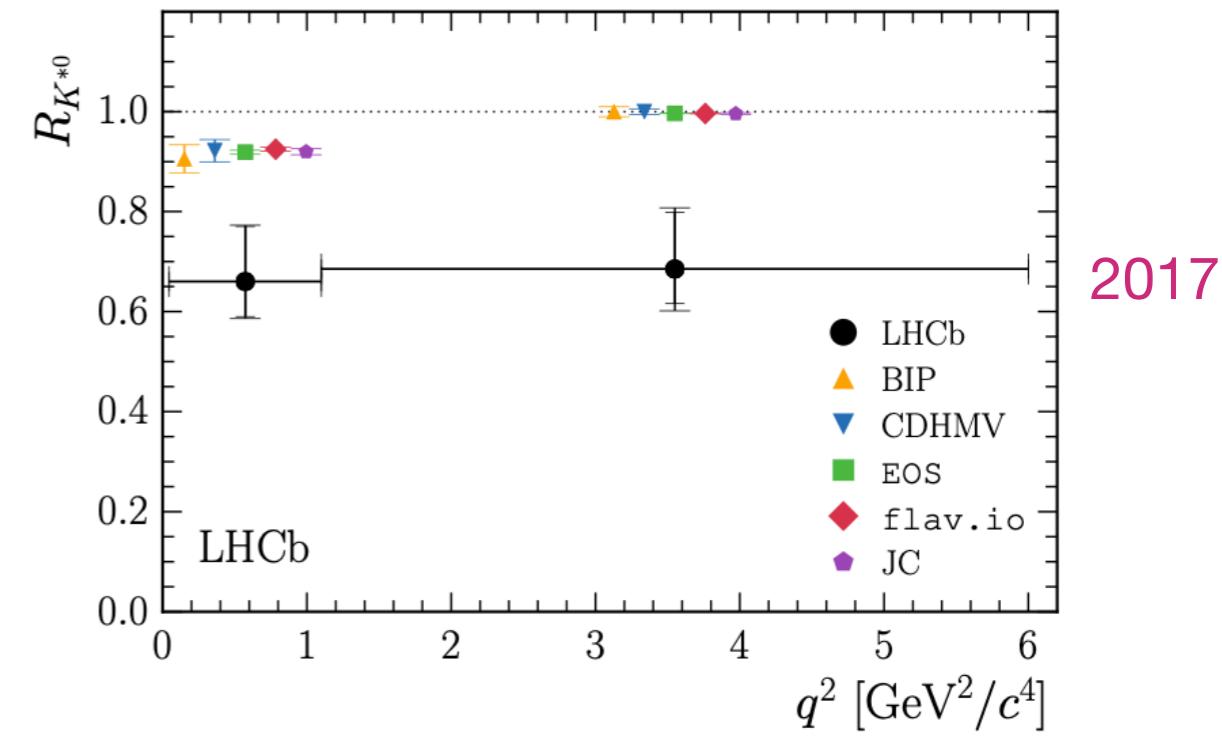
[Isidori, Bordone, Pattori, [1605.07633](#)]



Evidence for LFUV in R_K [3.1σ]

[LHCb, [2103.11769](#)]

2014
2019
2021



$2\sigma - 2.5\sigma$ deviations in R_{K^*}

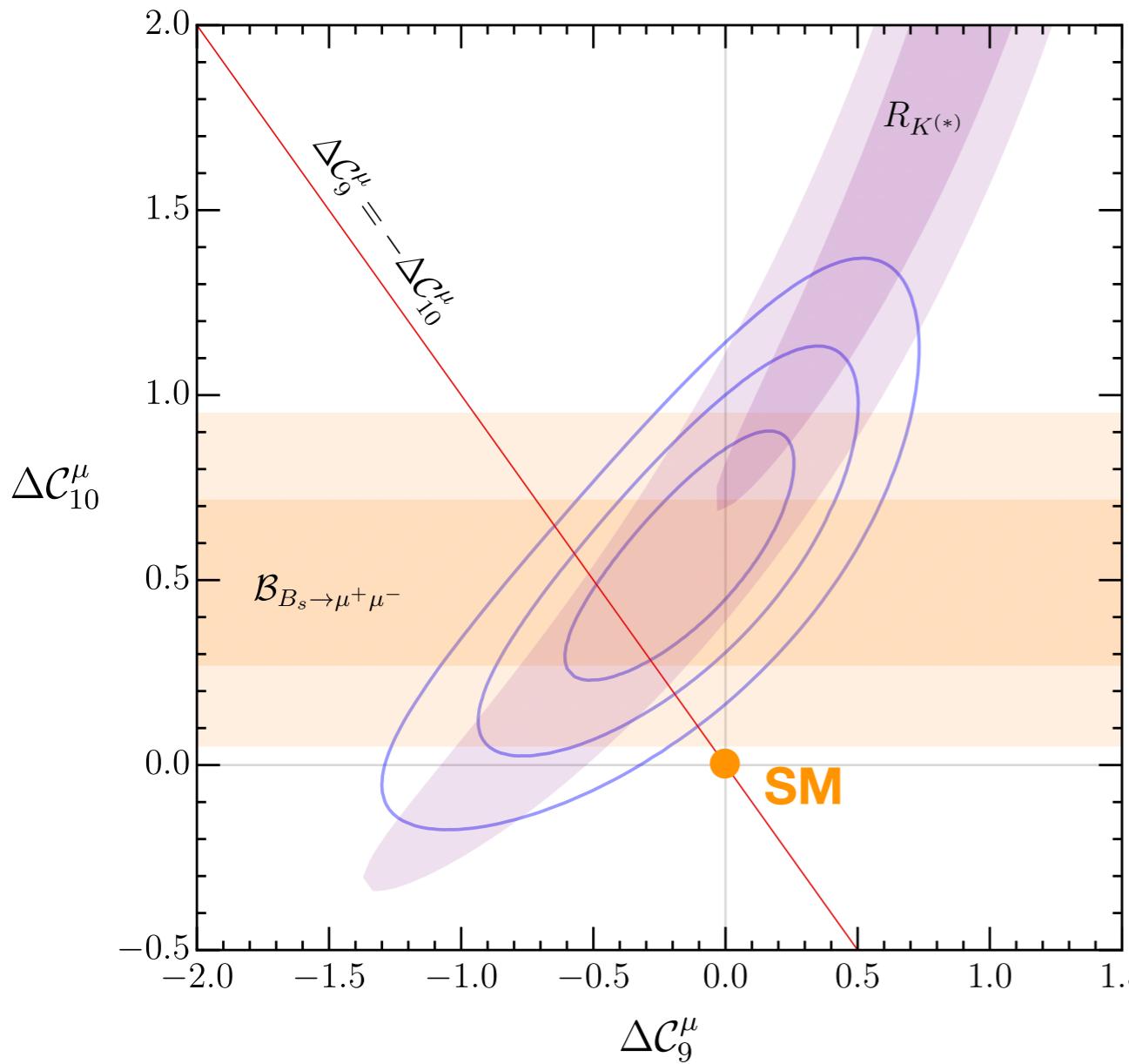
[LHCb, [1705.05802](#)]

2017

The $b \rightarrow s\ell^+\ell^-$ anomalies

Conservative fit using “clean observables” only

$$[C_i^\mu = C_i^{\text{SM}} + \Delta C_i^\mu + \Delta C_9^\mu, C_i^e = C_i^{\text{SM}} + \Delta C_9^e]$$



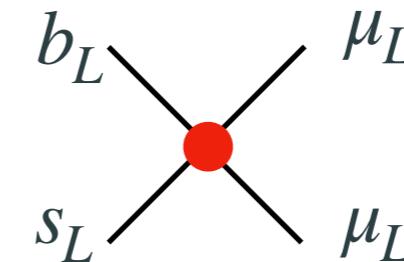
[Cornella et al., [2103.16558](#)]

$$\mathcal{L}_{\text{eff}} = -\frac{4G_F}{\sqrt{2}} V_{ts}^* V_{tb} \frac{\alpha}{4\pi} \sum_i C_i \mathcal{O}_i$$

$$\mathcal{O}_9^\mu = (\bar{s}_L \gamma_\mu b_L)(\bar{\mu} \gamma^\mu \mu) \quad C_9^{\text{SM}} \approx 4.1$$

$$\mathcal{O}_{10}^\mu = (\bar{s}_L \gamma_\mu b_L)(\bar{\mu} \gamma^\mu \gamma_5 \mu) \quad C_{10}^{\text{SM}} \approx -4.2$$

Left-handed new physics hypothesis
($\Delta C_9^\mu = -\Delta C_{10}^\mu$) preferred over the SM
by **4.6 σ**



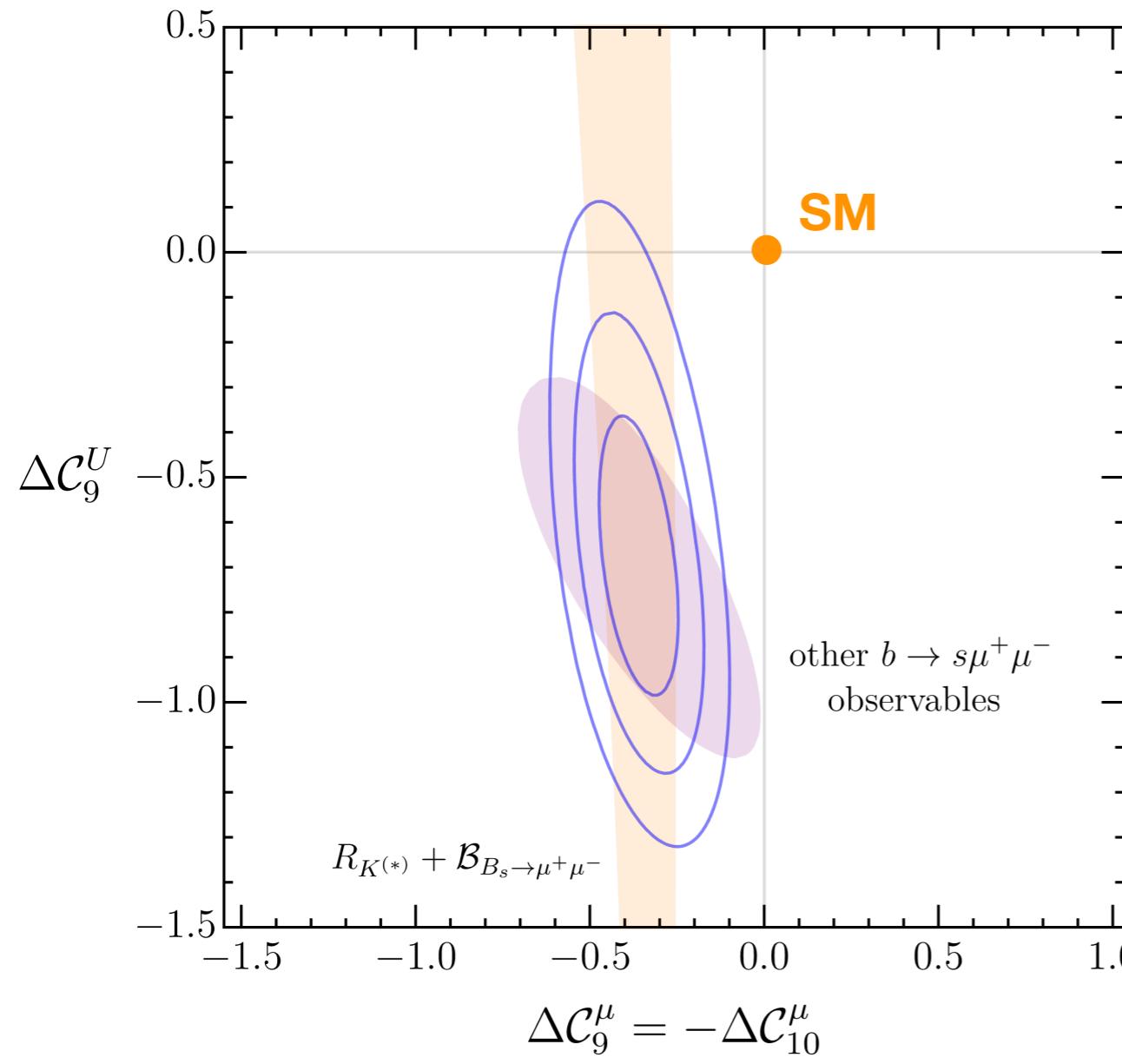
$$\sim 3 \times 10^{-5} G_F$$

$$\implies \frac{g_{\text{NP}}^2}{M_{\text{NP}}^2} \sim \frac{1}{(40 \text{ TeV})^2}$$

The $b \rightarrow s\ell^+\ell^-$ anomalies

Fit using all observables

$$[C_i^\mu = C_i^{\text{SM}} + \Delta C_i^\mu + \Delta C_9^U, C_i^e = C_i^{\text{SM}} + \Delta C_9^U]$$



[Cornella et al., [2103.16558](#)]

$$\mathcal{L}_{\text{eff}} = \frac{4G_F}{\sqrt{2}} V_{tb}^* V_{ts} \frac{\alpha}{4\pi} \sum_i C_i \mathcal{O}_i$$

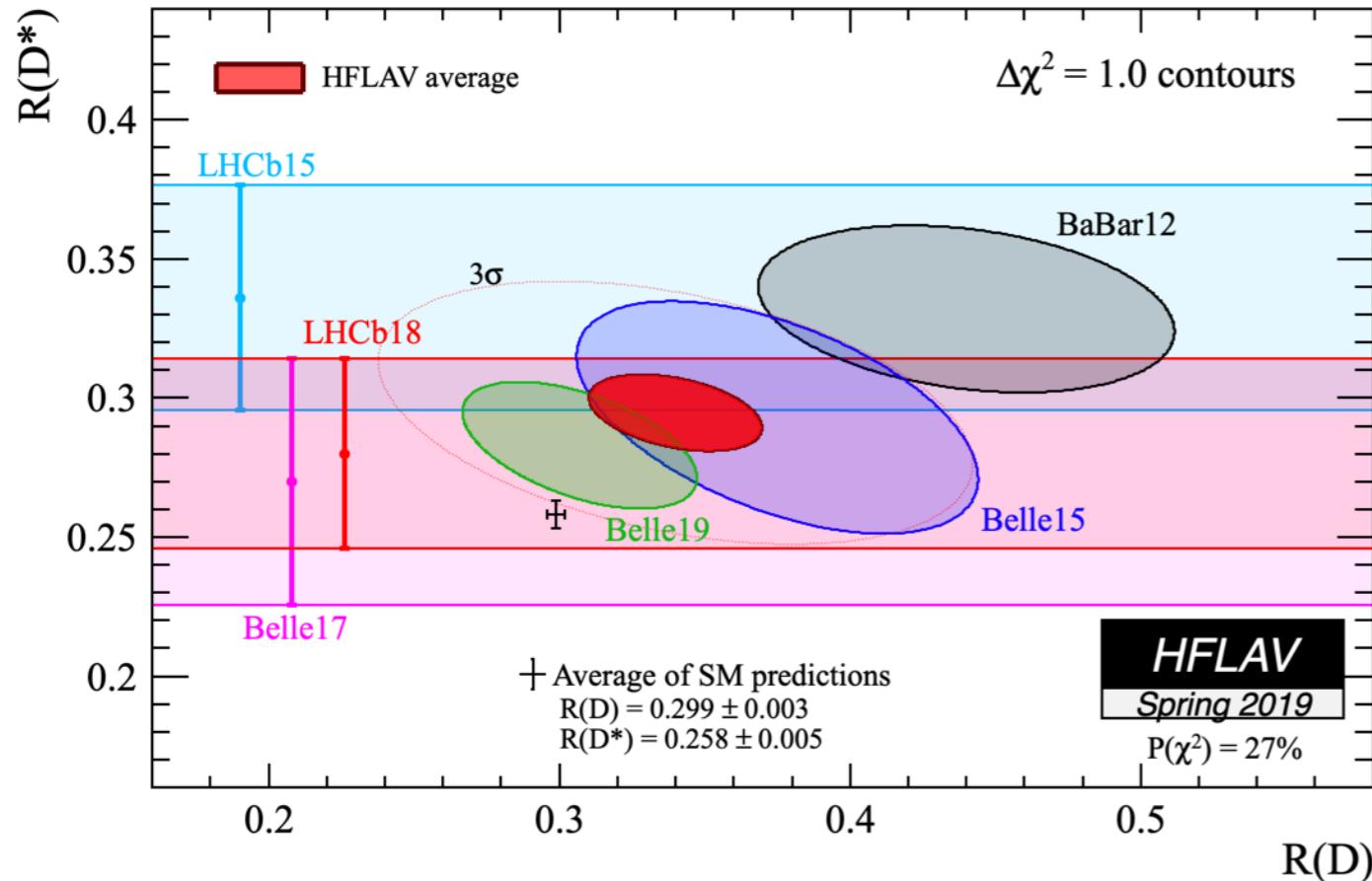
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$$\mathcal{O}_{10}^\mu = (\bar{s}_L \gamma_\mu b_L)(\bar{\mu} \gamma^\mu \gamma_5 \mu) \quad C_{10}^{\text{SM}} \approx -4.2$$

Adding all $b \rightarrow s\ell\ell$ measurements [$\mathcal{O}(100)$ observables] with the current best estimate of the QCD contributions, the significance of the new physics hypothesis grows to well over 5σ

[Altmannshofer, Stangl, [2103.13370](#); Algueró et al, [2104.08921](#); Hurth et al. [2104.10058...](#)]

The $b \rightarrow c\tau\bar{\nu}$ anomalies

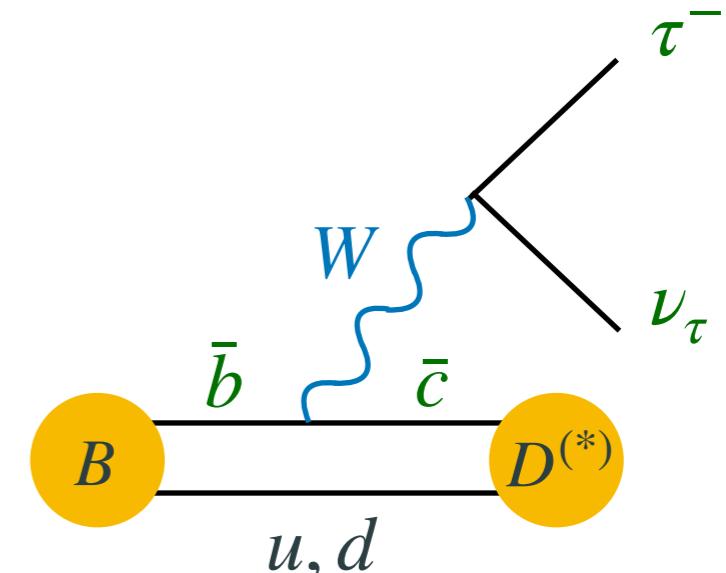


$$R_{D^{(*)}} \equiv \frac{\mathcal{B}(B \rightarrow D^{(*)}\tau\bar{\nu})}{\mathcal{B}(B \rightarrow D^{(*)}\ell\bar{\nu})} \quad [\ell = e, \mu]$$

$$R_{D^*}/R_{D^*}^{\text{SM}} = 1.14 \pm 0.05$$

$$R_D/R_D^{\text{SM}} = 1.14 \pm 0.10$$

$$\rho = -0.38$$

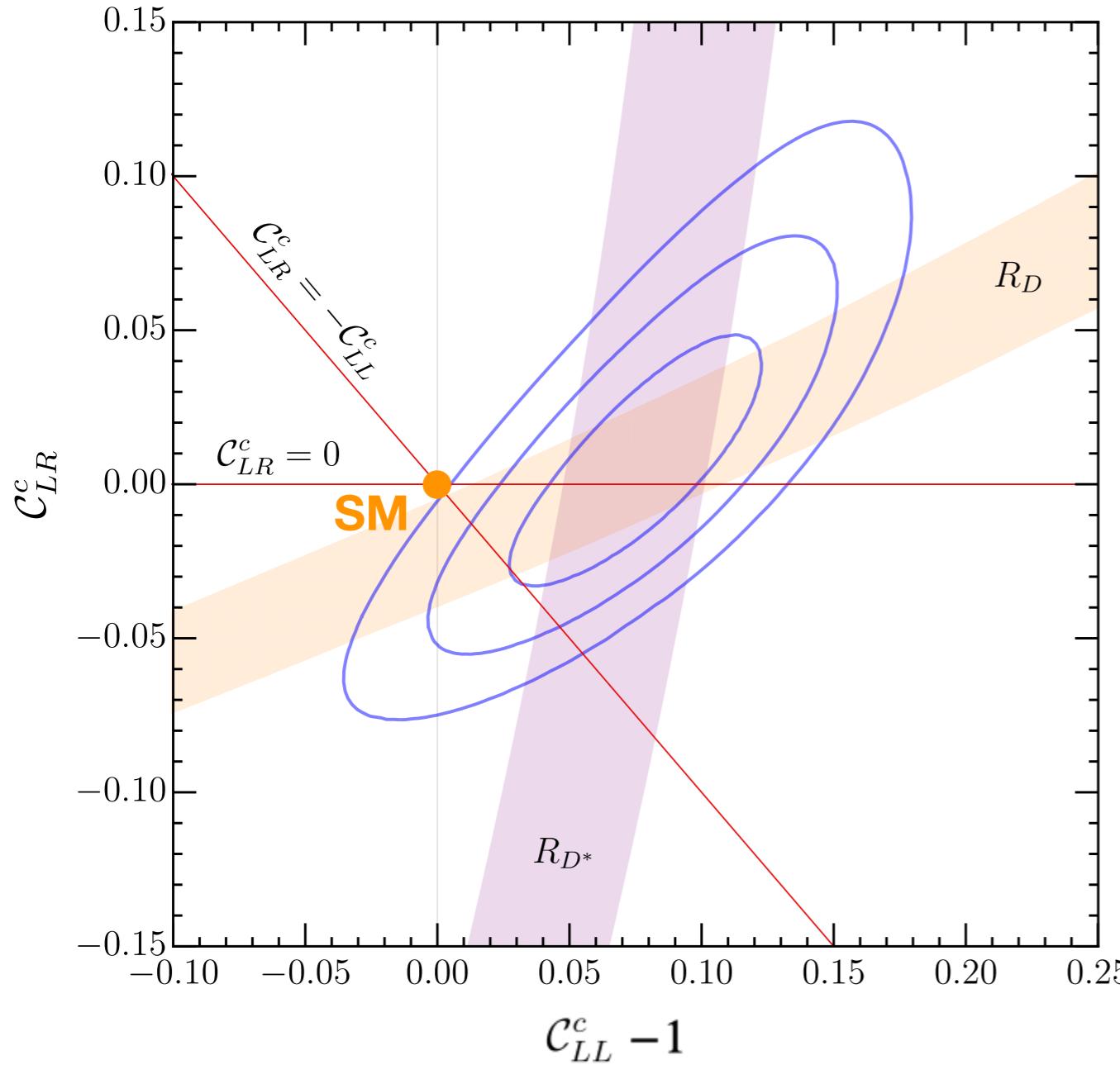


Clean observable: QCD uncertainties cancel (to a large extent) in the ratios

Consistent result by three experiments (BaBar, Belle and LHCb) show a 3.1σ deviation (R_D and R_{D^*} comb.)

N.B.: SM different from unity due to τ mass effects (phase space)

The $b \rightarrow c\tau\bar{\nu}$ anomalies



[Cornella et al., [2103.16558](#)]

$$\mathcal{L}_{\text{eff}} = -\frac{4G_F}{\sqrt{2}} V_{cb} \sum_i C_i \mathcal{O}_i$$

$$\mathcal{O}_{LL}^i = (\bar{u}_L^i \gamma_\mu \nu_L)(\bar{\tau}_L \gamma^\mu b_L) \quad C_{LL}^{\text{SM}} = 1$$

$$\mathcal{O}_{LR}^i = (\bar{u}_L^i \gamma_\mu \nu_L)(\bar{\tau}_R \gamma^\mu b_R) \quad C_{LR}^{\text{SM}} = 0$$

Preference for left-handed type new physics [analogous to the SM]

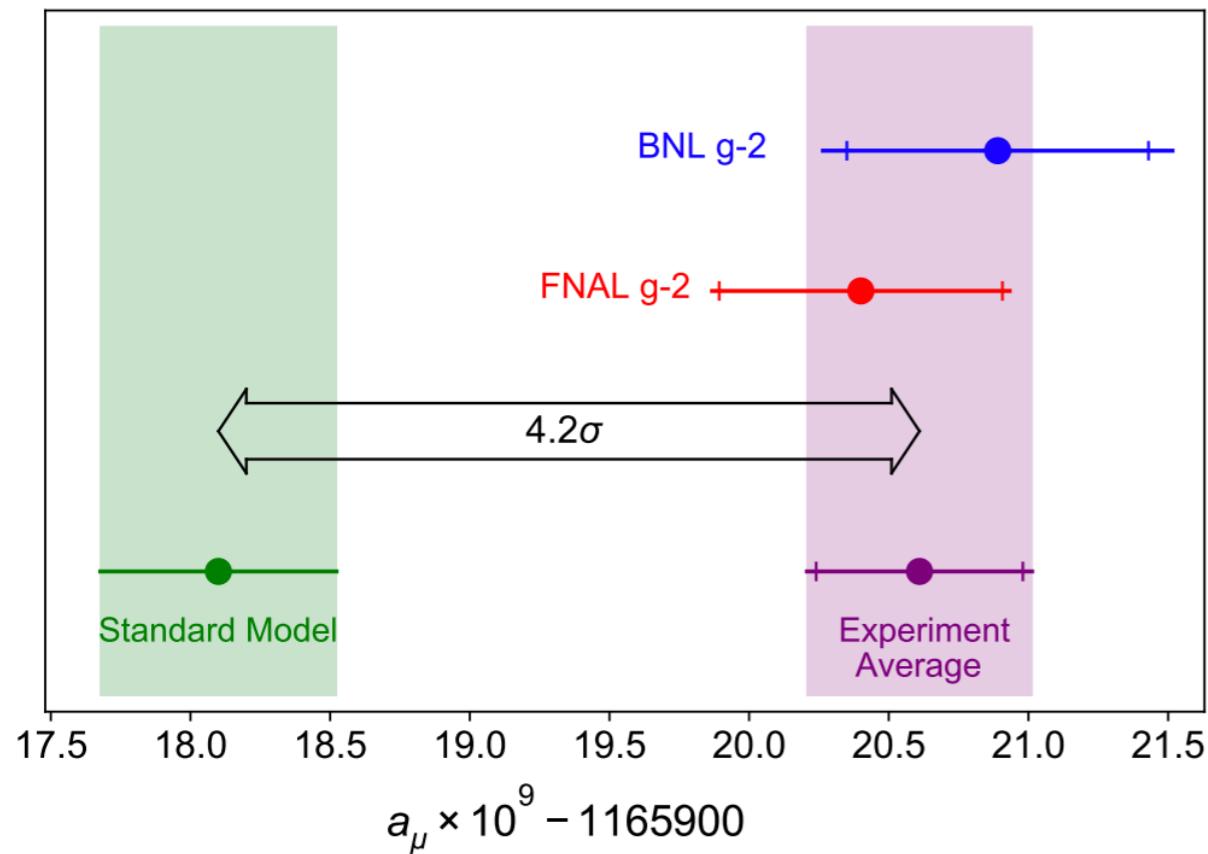
b_L τ_L
 u_L ν_L

$$\sim 10^{-2} G_F$$

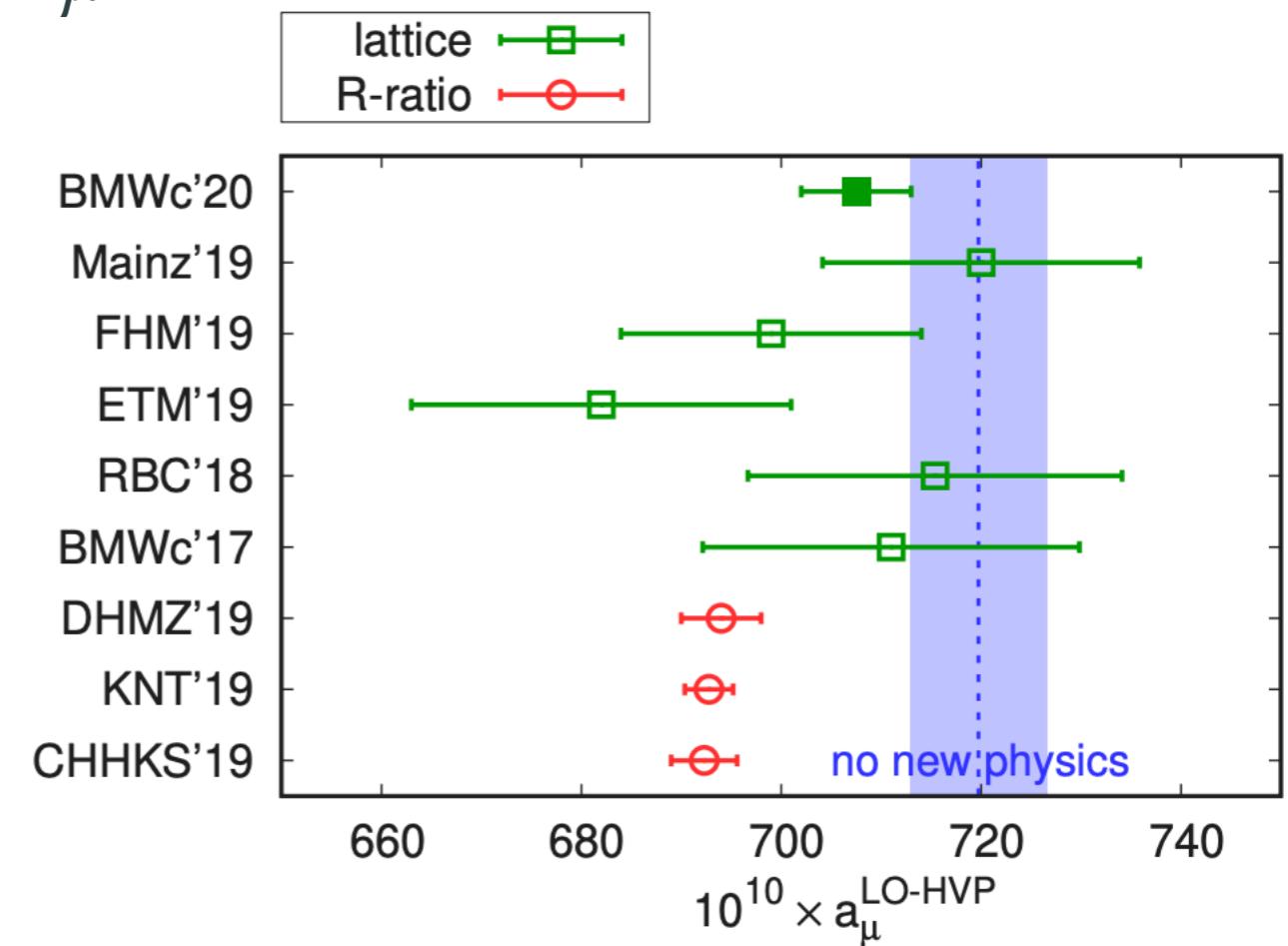
$$\implies \frac{g_{\text{NP}}^2}{M_{\text{NP}}^2} \sim \frac{1}{(2 \text{ TeV})^2}$$

Although other new physics structures are also possible

A brief comment on $(g - 2)_\mu$



[Muon g-2 collaboration, [2104.03281](#)]



[Borsanyi et al., [2002.12347](#)]

Recent confirmation by Fermilab of the Brookhaven experimental result
 [strong evidence of new physics 4.2σ (Fermilab + Brookhaven comb.)]

While the SM prediction is dominated by QED, the SM error is dominated by QCD
 [current evaluation uses data-driven methods (R-ratio)... in tension with BMW lattice]

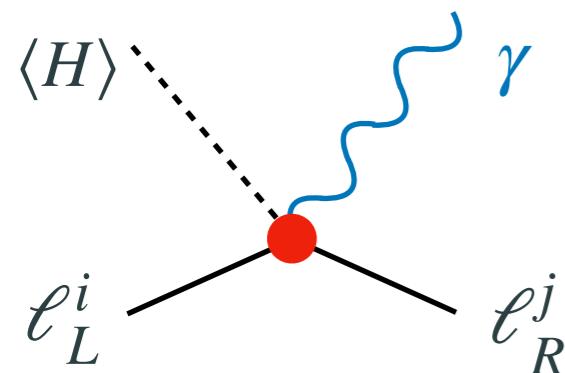
A brief comment on $(g - 2)_\mu$

The discrepancy is roughly the same size as the SM electroweak contribution

$$\Delta a_\mu \equiv a_\mu^{\text{exp}} - a_\mu^{\text{SM}} \approx (a_\mu^{\text{SM}})_{\text{EW}} \approx \frac{m_\mu^2}{16\pi^2} \times \frac{4G_F}{\sqrt{2}} \approx 3 \times 10^{-9}$$

$$a_\mu \equiv \frac{g_\mu - 2}{2}$$

→ NP is either light or is **not chirally suppressed**



$$\frac{\mathcal{B}(\mu \rightarrow e\gamma)}{3 \times 10^{-13}} \approx \left(\frac{\Delta a_\mu}{3 \times 10^{-9}} \right)^2 \left(\frac{\theta_{12}}{10^{-5}} \right)^2$$

$$\frac{\mathcal{B}(\tau \rightarrow \mu\gamma)}{4 \times 10^{-8}} \approx \left(\frac{\Delta a_\mu}{3 \times 10^{-9}} \right)^2 \left(\frac{\theta_{23}}{10^{-2}} \right)^2$$

Naive expectation:

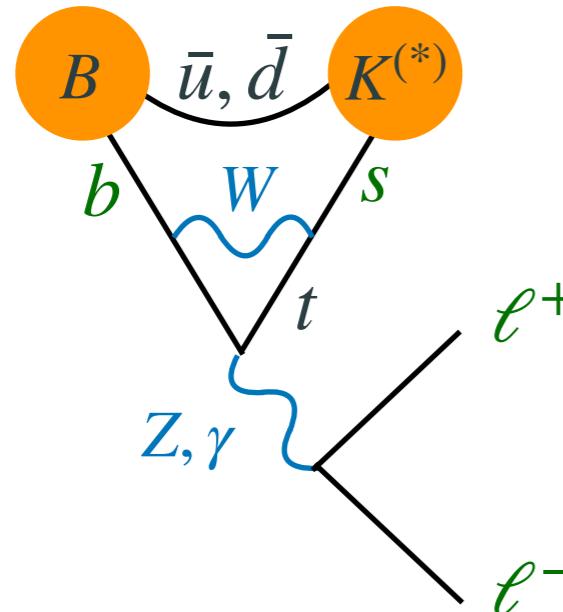
$$\theta_{12}^2 \sim m_e/m_\mu \approx 5 \times 10^{-3}$$

$$\theta_{23}^2 \sim m_\mu/m_\tau \approx 6 \times 10^{-2}$$

→ NP should be **nearly lepton flavor conserving**

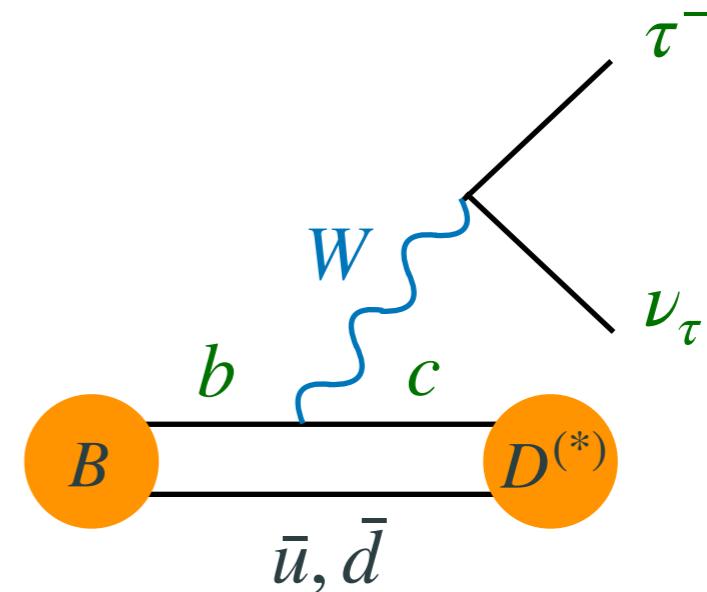
Not easy to reconcile $(g - 2)_\mu$ with both B anomalies and, more generally, with models with a “natural” flavor structure (i.e. similar to that of the SM Yukawa)

Combined explanation of B anomalies: Theoretical hint



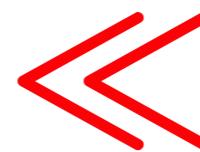
$$3_Q \rightarrow 2_Q 2_L 2_L$$

~10% of a SM loop effect



$$3_Q \rightarrow 2_Q 3_L 3_L$$

~10% of a SM tree-level effect



The only source of **lepton flavor universality violation** in the SM (Yukawas) follows a very similar trend: $y_e \ll y_\mu \ll y_\tau$

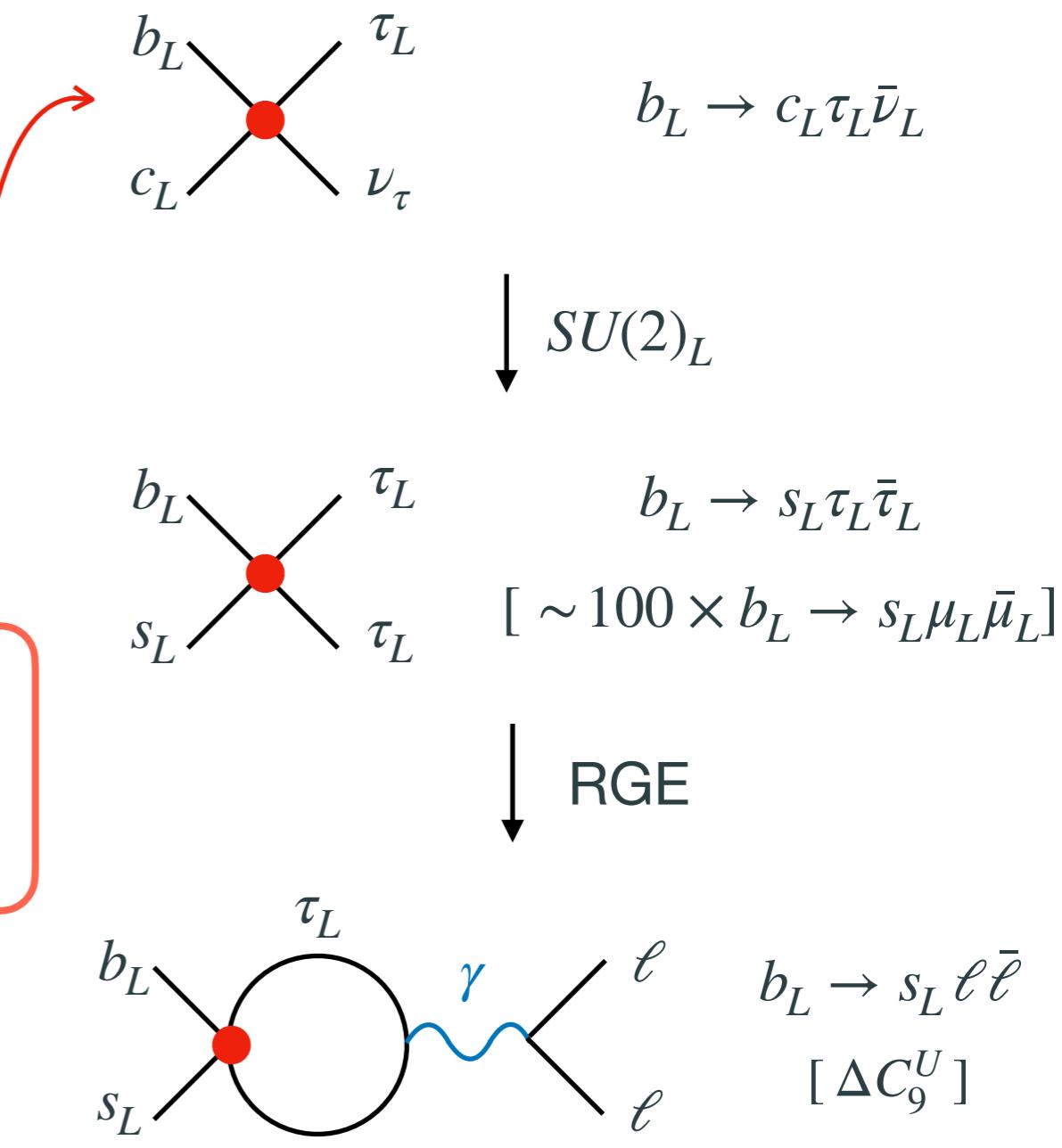
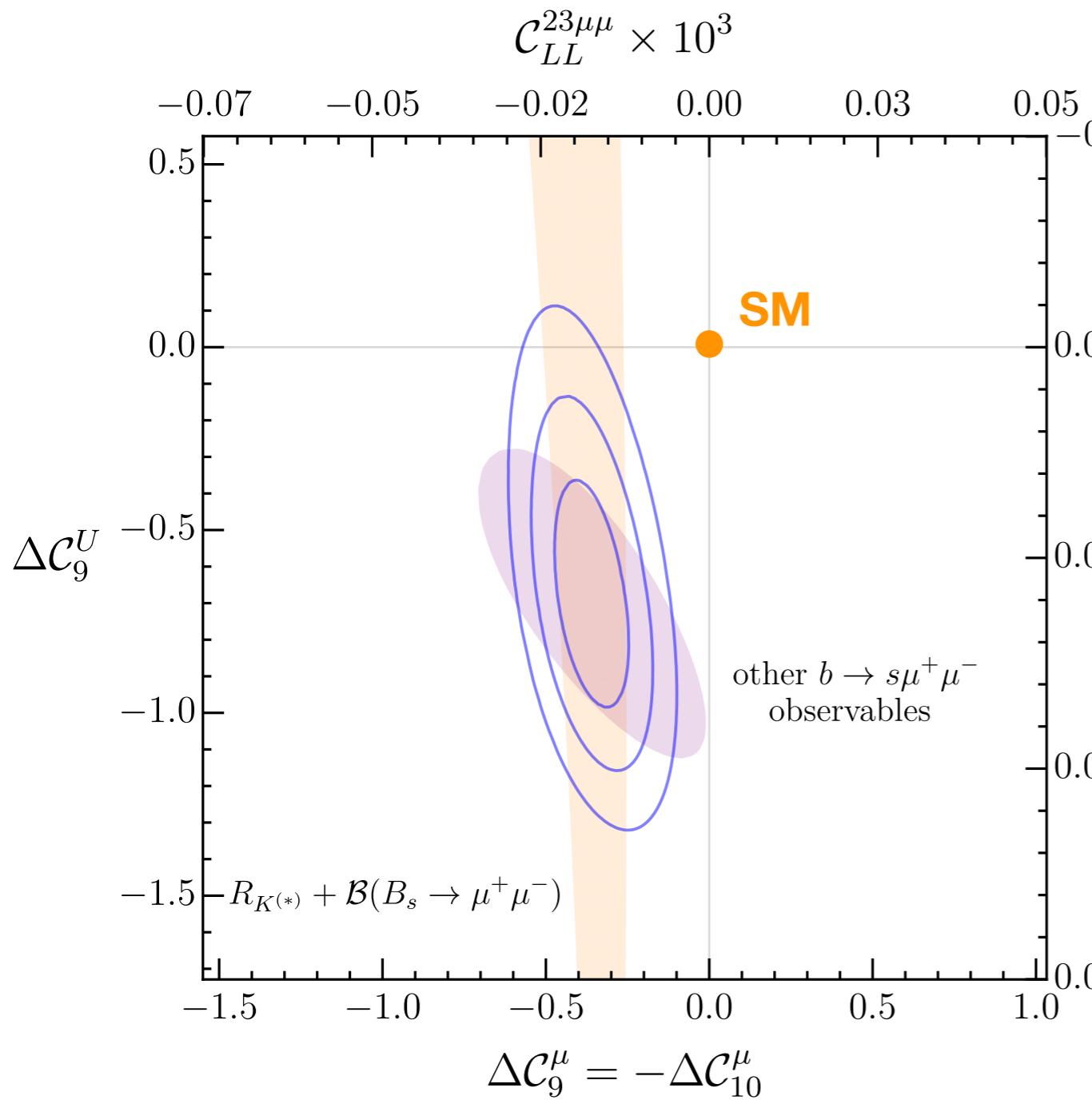
Combined explanation with TeV-scale new physics dominantly coupled to 3rd family with $\mathcal{O}(0.1)$ suppressed couplings to 2nd family

[Hints of underlying (flavor) symmetries and/or solutions to the SM flavor puzzle?]

Combined explanation of B anomalies: Pheno hint

$$\mathcal{O}_{LL}^{ij\alpha\beta} = (q_L^i \gamma_\mu \ell_L^\alpha)(\bar{\ell}_L^\beta \gamma^\mu q_L^j)$$

[Normalized to G_F]

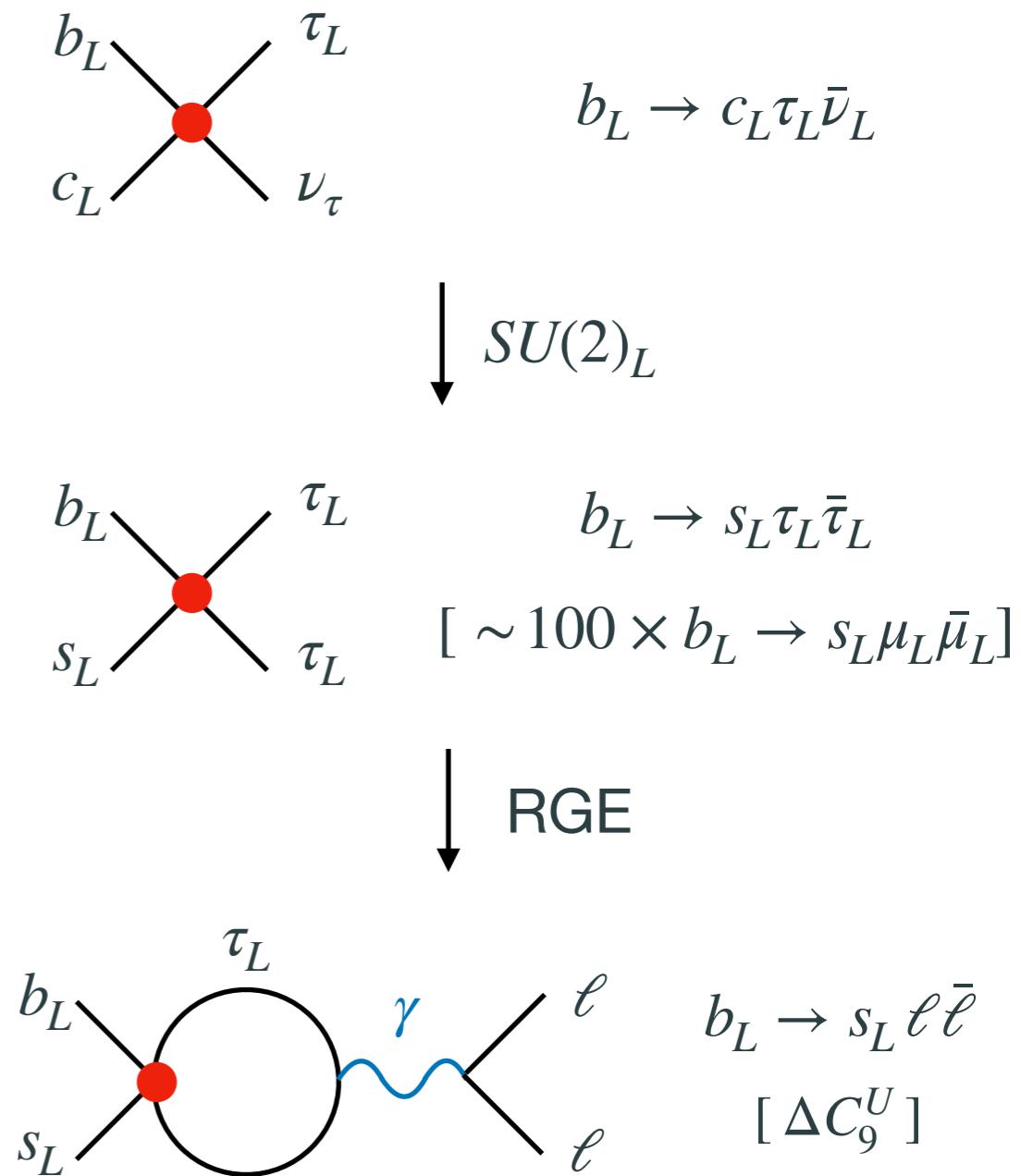
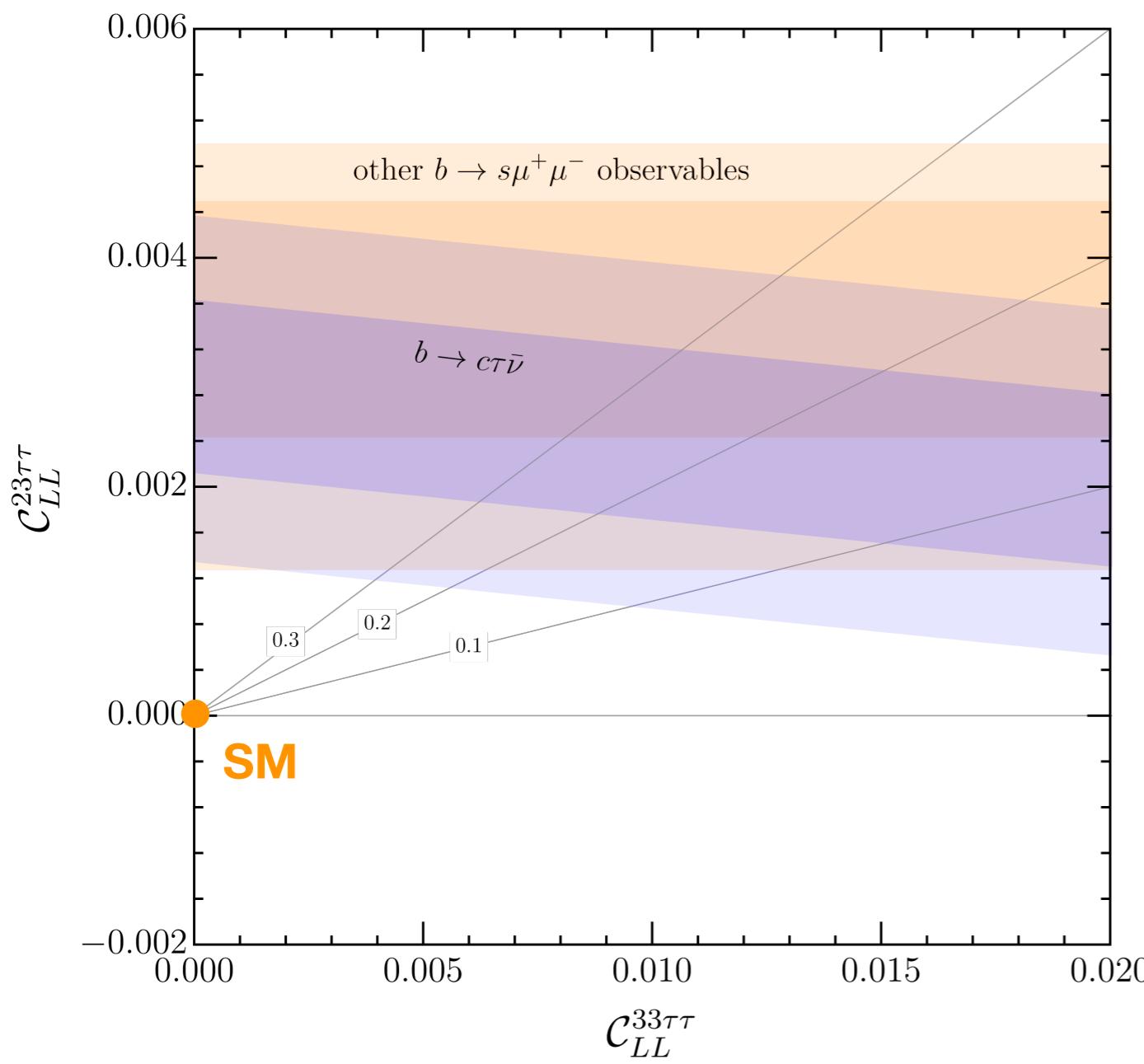


[Bobeth, Haisch, [1109.1826](#); Crivellin et al., [1807.02068](#);
Algueró et al., [1809.08447](#)]

Combined explanation of B anomalies: Pheno hint

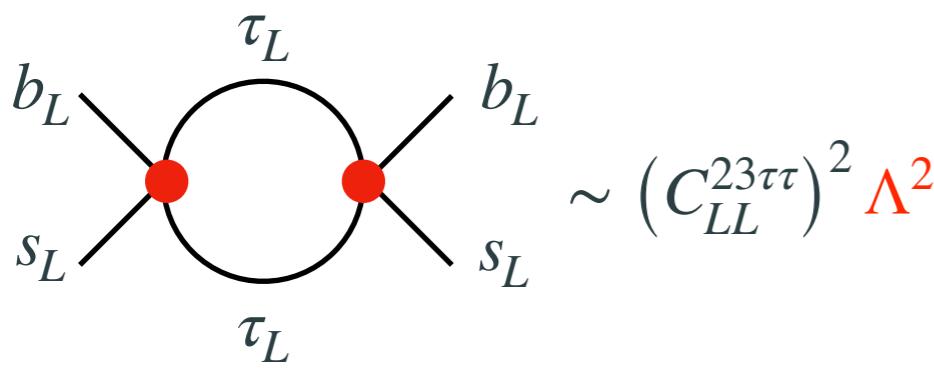
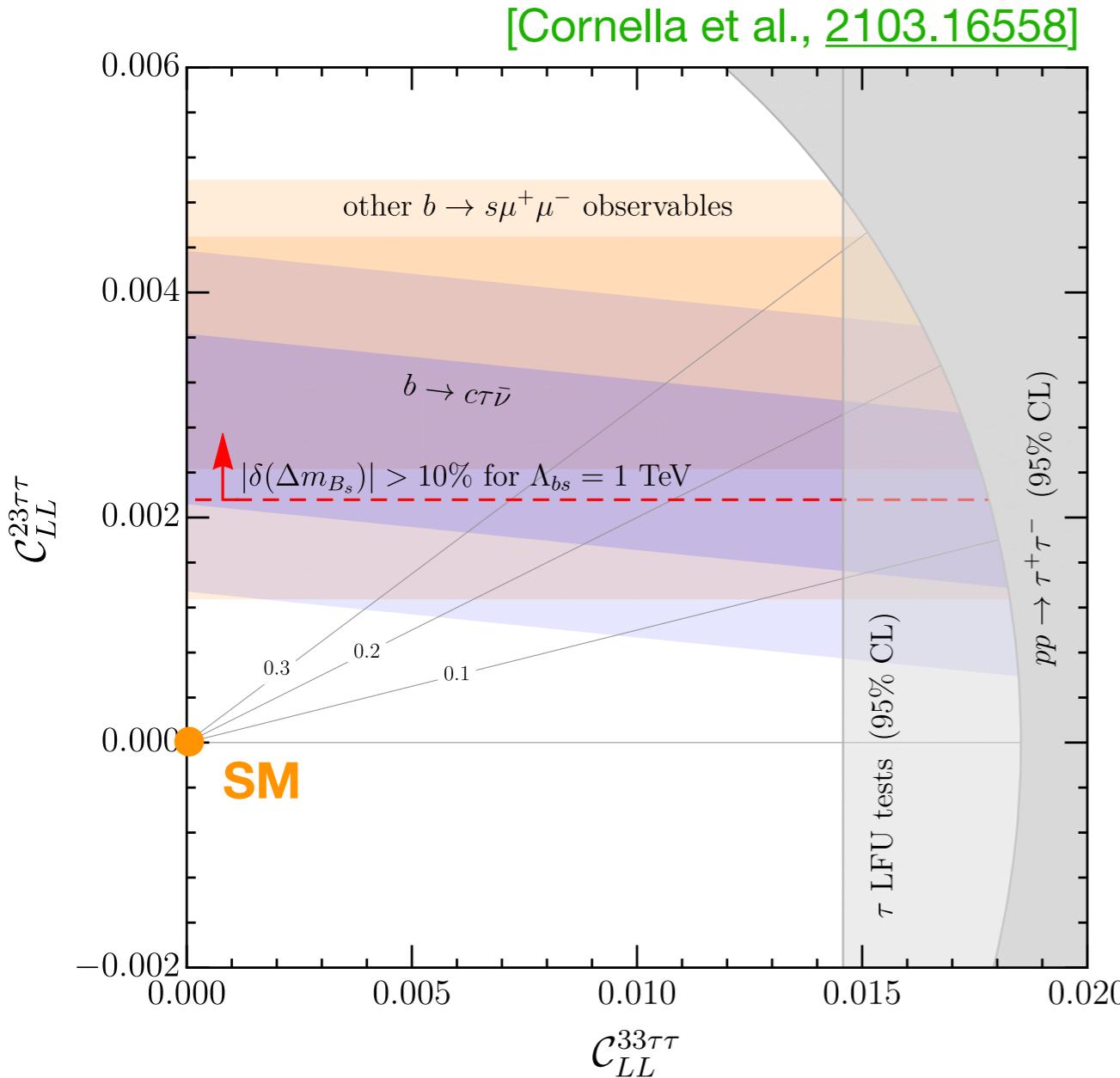
$$\mathcal{O}_{LL}^{ij\alpha\beta} = (q_L^i \gamma_\mu \ell_L^\alpha)(\bar{\ell}_L^\beta \gamma^\mu q_L^j)$$

[Normalized to G_F]



[Bobeth, Haisch, [1109.1826](#); Crivellin et al., [1807.02068](#); Algueró et al., [1809.08447](#)]

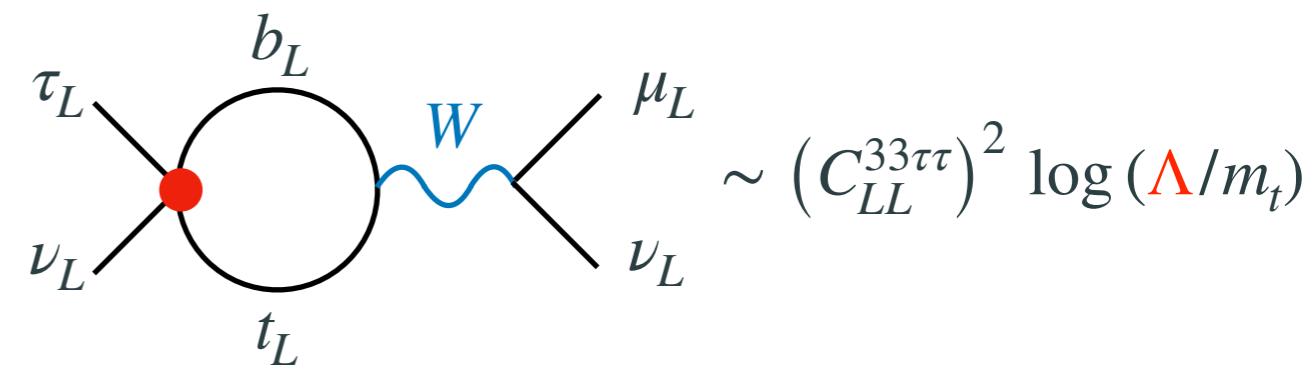
Additional constraints



$$\mathcal{O}_{LL}^{ij\alpha\beta} = (q_L^i \gamma_\mu \ell_L^\alpha)(\bar{\ell}_L^\beta \gamma^\mu q_L^j) = \frac{1}{2} [\mathcal{O}_{\ell q}^{(1)} + \mathcal{O}_{\ell q}^{(3)}]^{ij\alpha\beta}$$

There is a long list of constraints:

- * No observation of $B \rightarrow K^{(*)}\nu\bar{\nu}$
- * High- p_T searches at LHC [$pp \rightarrow \tau^+\tau^-$]
- * Lepton universality tests in τ decays
[e.g. $\tau \rightarrow \mu\nu\bar{\nu}$ vs $\tau \rightarrow e\nu\bar{\nu}$]
- * Meson-antimeson mixing ($B_s - \bar{B}_s$)



From EFT analyses to simplified models



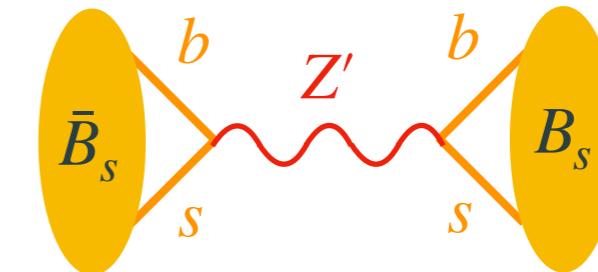
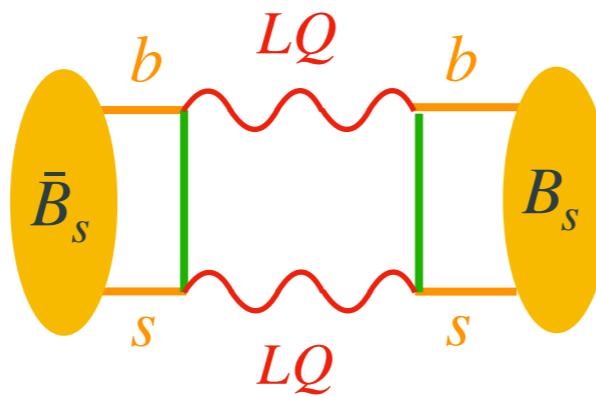
Which mediator?

If we consider only tree-level mediators [$R_{D^{(*)}}$ point to a low new physics scale]



Leptoquarks (both scalars and vectors) have two important advantages

1. $\Delta F = 2$ &
 $\tau \rightarrow \mu\nu\bar{\nu}$



2. Direct searches: t-channel versus resonant s-channel production

Which leptoquark?

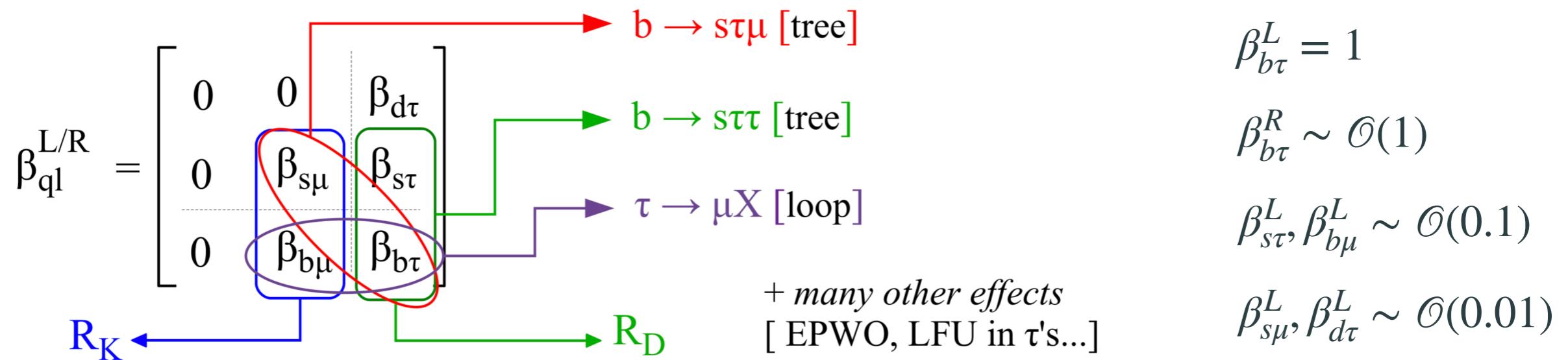
The U_1 vector leptoquark is (one of) the most promising mediator(s) to explain the B anomalies

$$\mathcal{L} \supset \frac{g_U}{\sqrt{2}} U_1^\mu \left[\beta_L^{i\alpha} (\bar{q}_L^i \gamma_\mu \ell_L^\alpha) + \beta_R^{i\alpha} (\bar{d}_R^i \gamma_\mu e_R^\alpha) \right] + \text{h.c.} \quad U_1 \sim (3, 1, 2/3)$$

It has two important advantages with respect to other leptoquarks:

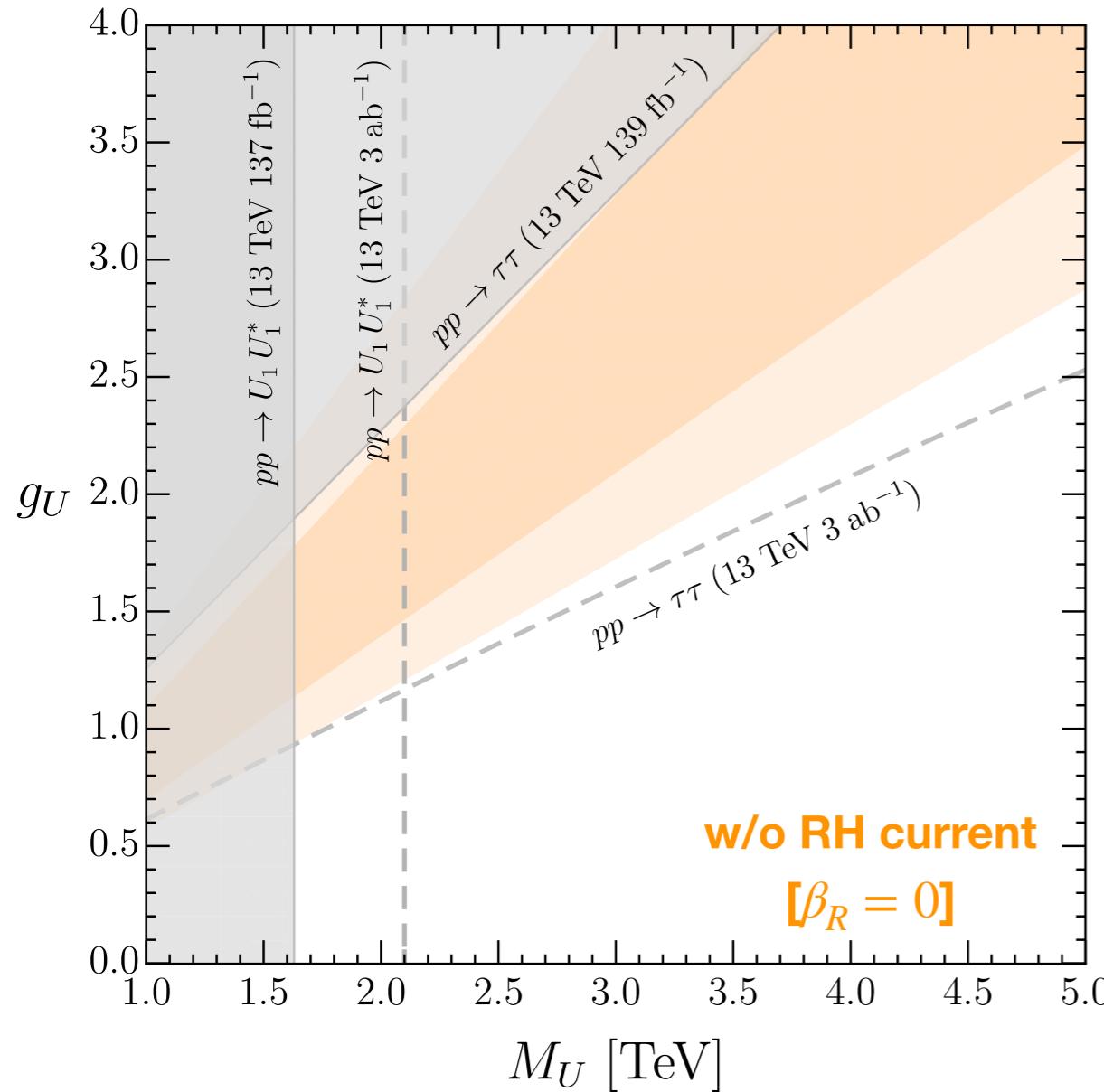
1. No tree-level contribution to $b \rightarrow s\nu\bar{\nu}$
2. Automatic protection from proton decay

It provides a good description of all low-energy data with a “natural” flavor structure

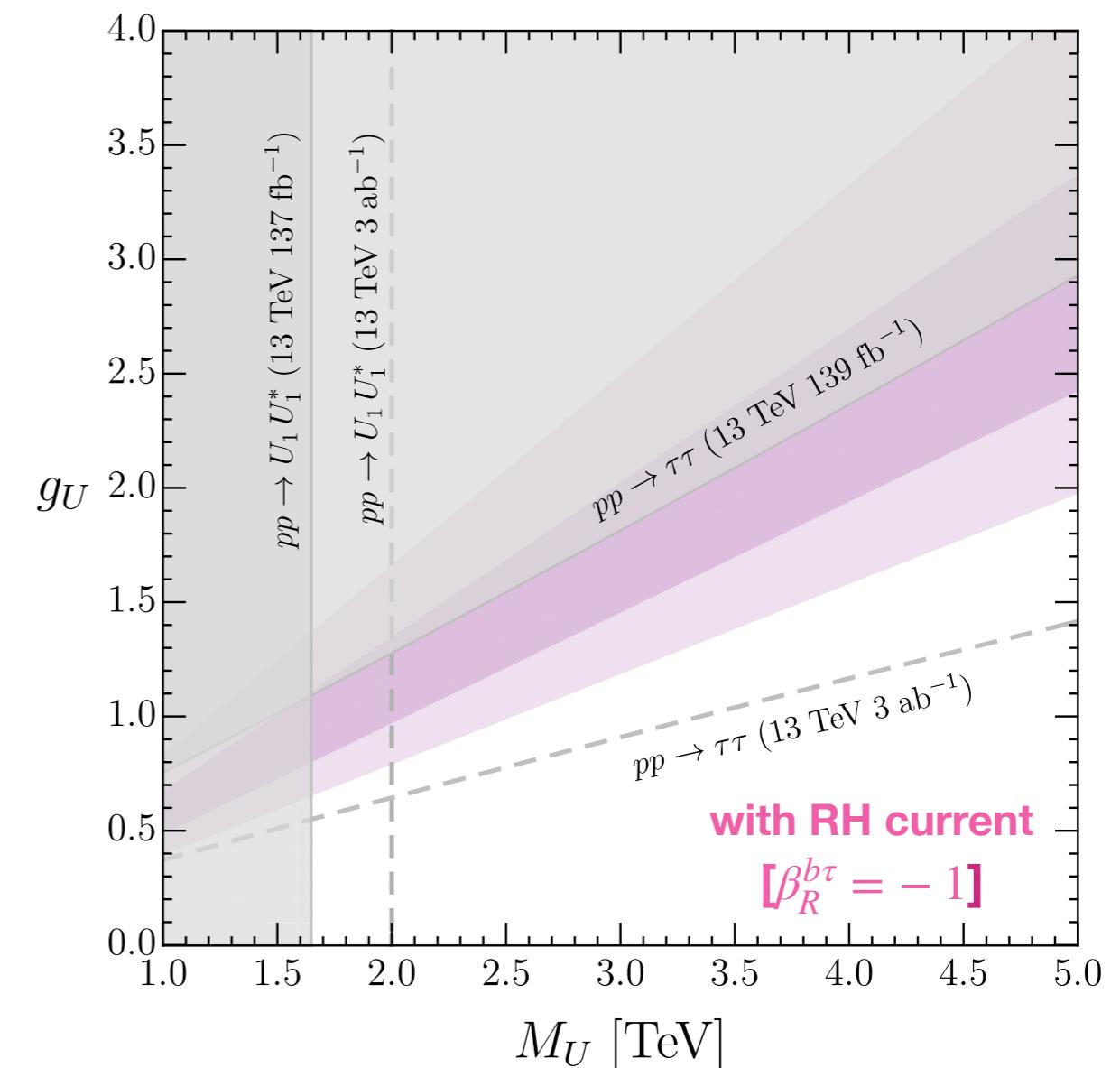


Corroborating the U_1 hypothesis: High- p_T searches

U_1 leptoquark solution also consistent with high- p_T data and [within the HL-LHC reach!](#)
[Expected **enhancement** of high- p_T $\tau^+\tau^-$ pairs in Drell-Yan data]



[Cornella et al., [2103.16558](#)]

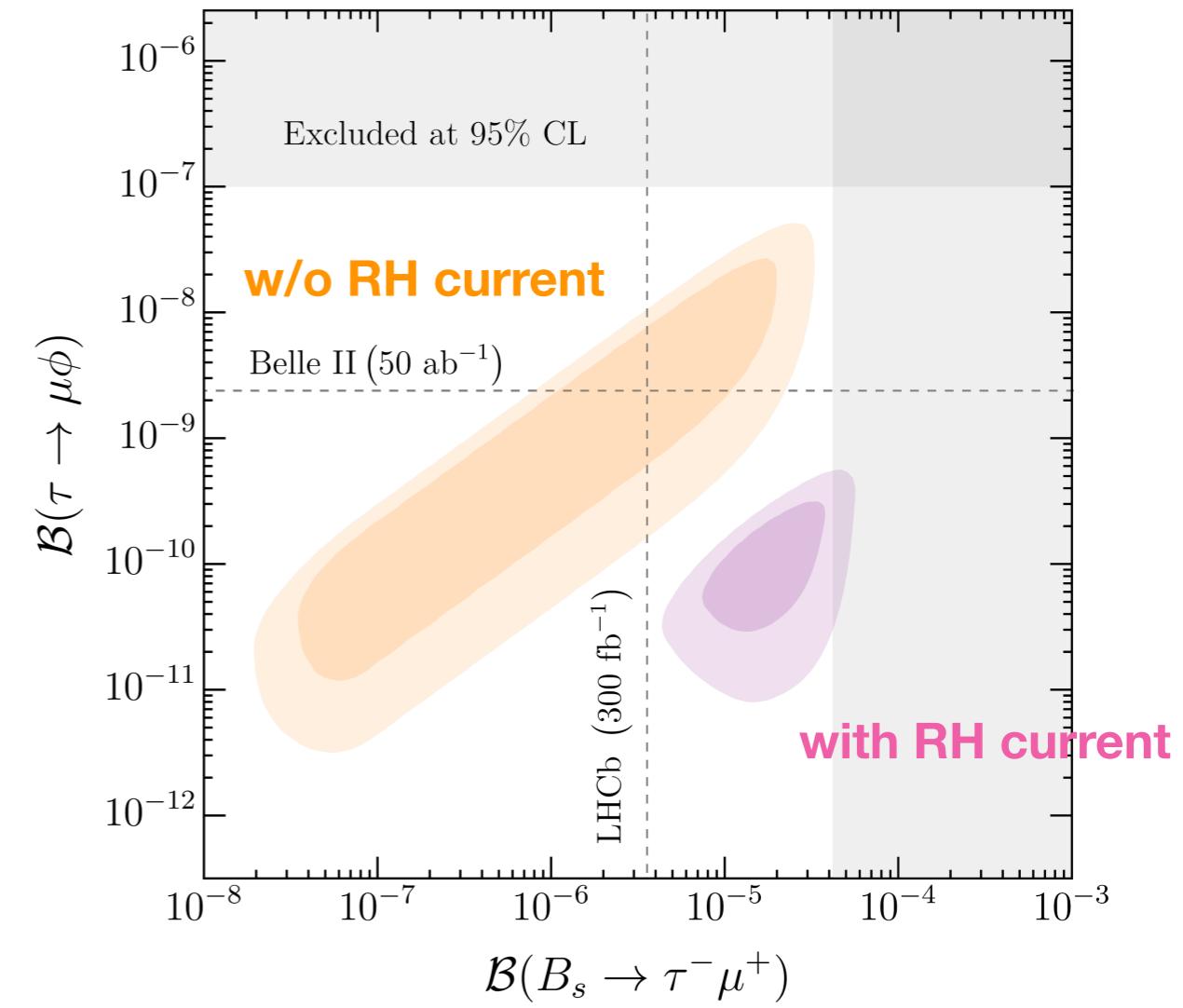
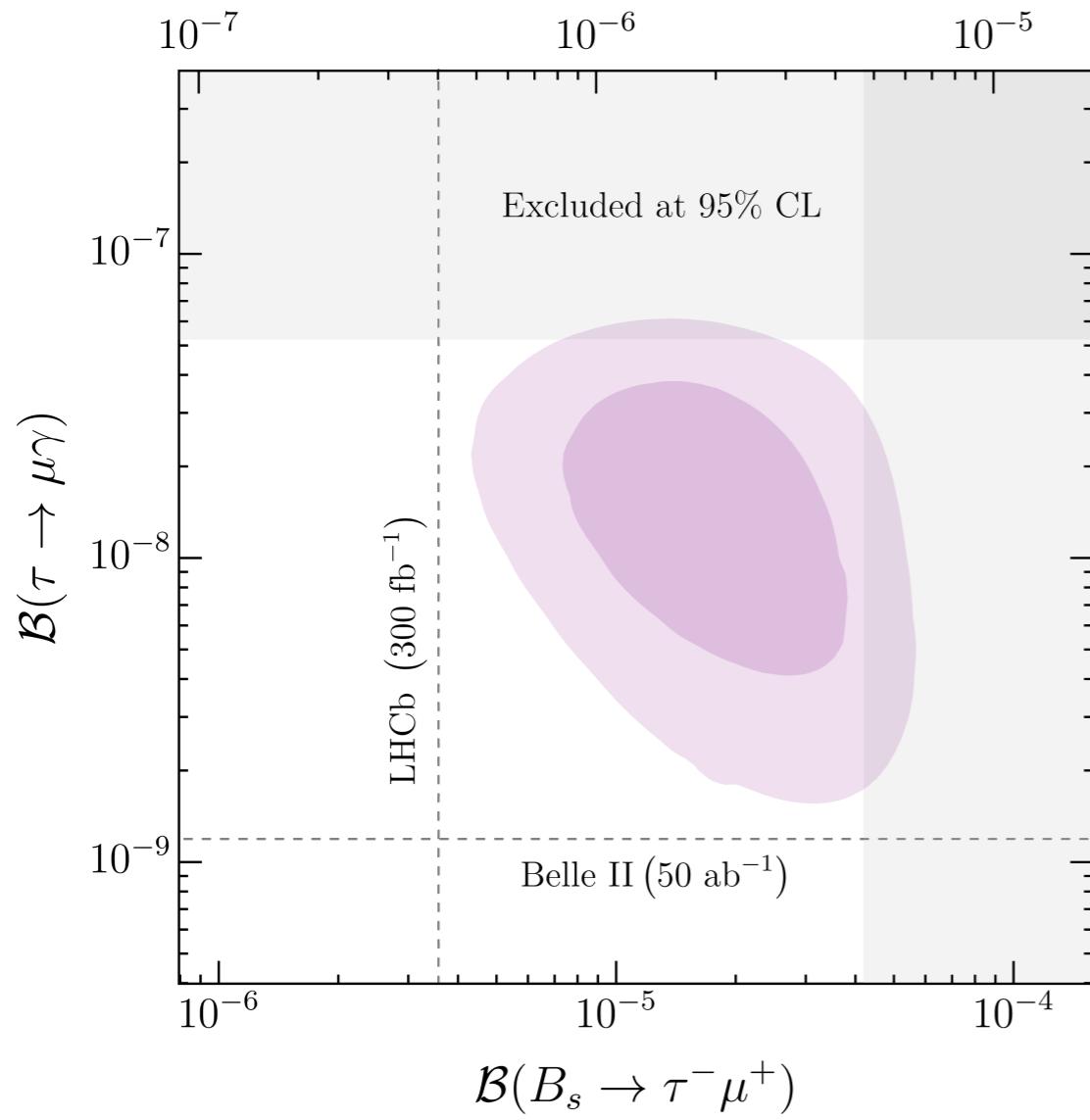


[$pp \rightarrow \tau\tau$ for U_1 originally proposed in
Faroughy, Greljo, Kamenik, [1609.07138](#)]

Corroborating the U_1 hypothesis: other B and τ decays

LHCb & Belle II essential to confirm/exclude this solution in a large class of correlated observables and, if confirmed, to determine its coupling & flavor structure

I. LFV in B and τ decays

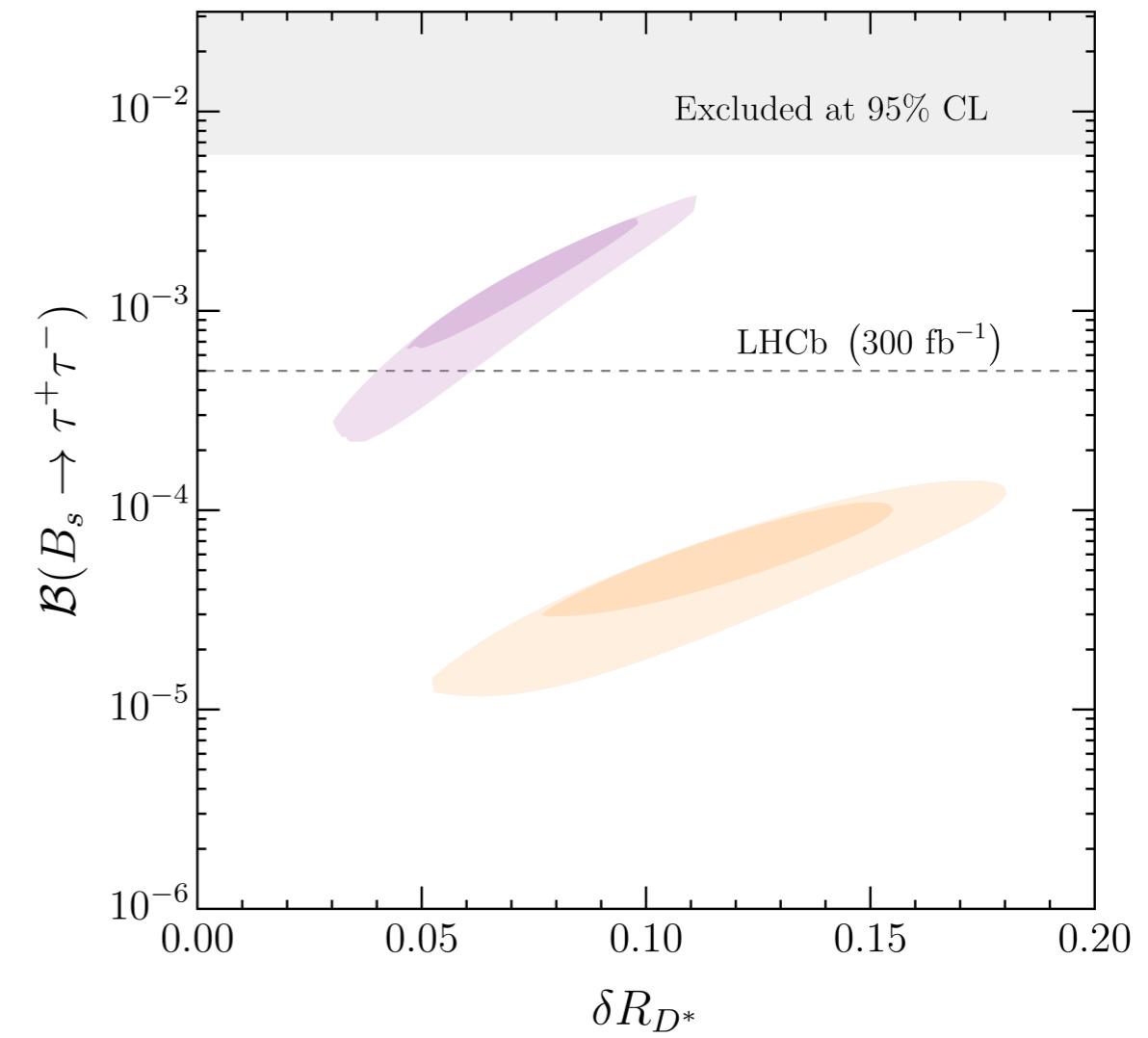
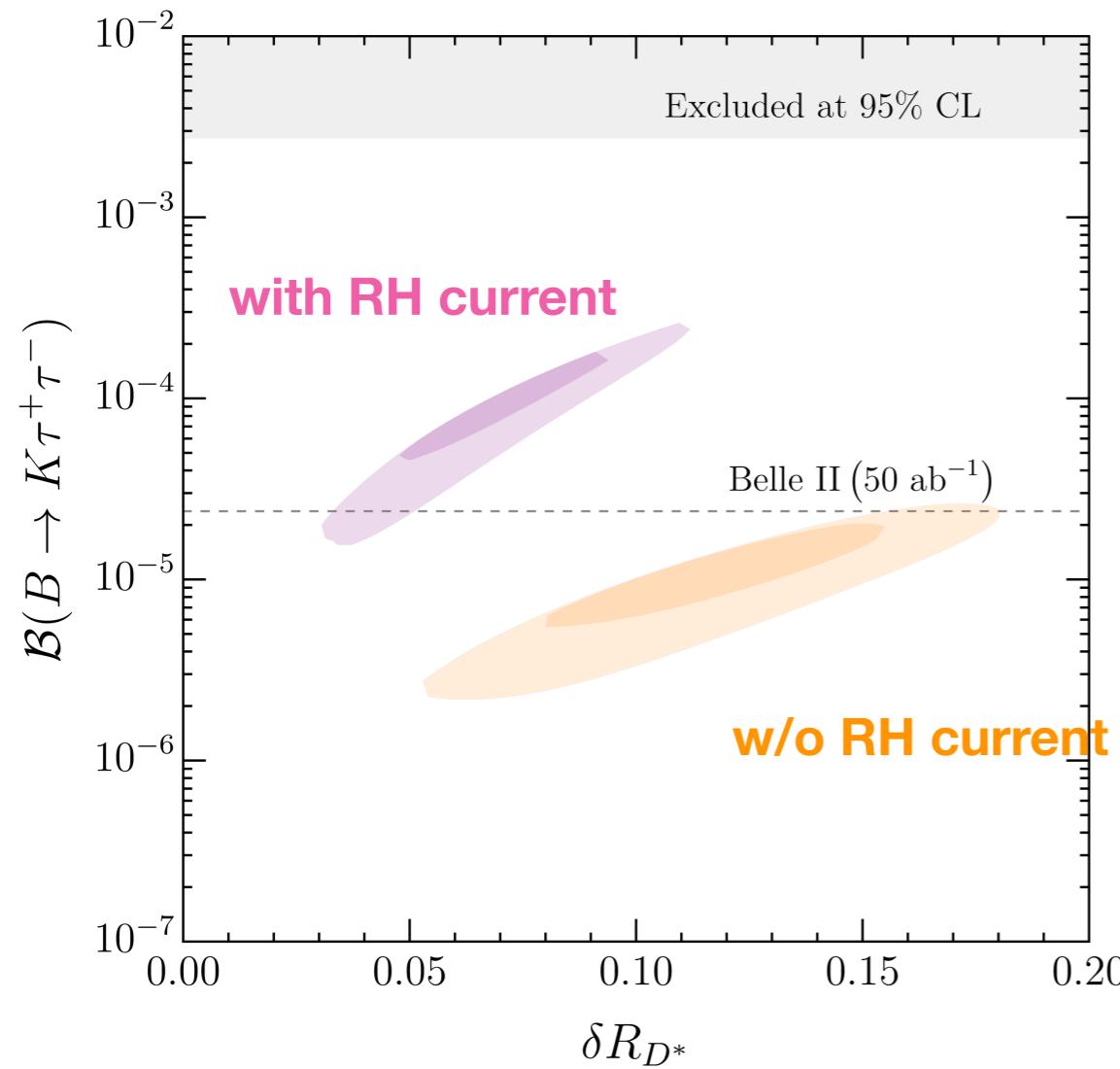


Corroborating the U_1 hypothesis: other B and τ decays

LHCb & Belle II essential to confirm/exclude this solution in a large class of correlated observables and, if confirmed, to determine its coupling & flavor structure

II. $b \rightarrow s\tau^+\tau^-$ decays

$$\text{N.B: } \delta R_{D^*} = \frac{R_{D^*} - R_{D^*}^{SM}}{R_{D^*}^{SM}}$$



[Cornella et al., 2103.16558]

Speculations about a UV complete description



Gauge UV completion for the U_1 leptoquark

The U_1 solution points to Pati-Salam unification

[Pati, Salam, [Phys. Rev. D10 \(1974\) 275](#)]

(7 years after the SM was proposed)

$$SU(4) \times SU(2)_L \times SU(2)_R$$

$$SU(4) \sim \begin{pmatrix} G^a & U^\alpha \\ (U^\alpha)^* & Z' \end{pmatrix}$$

$$\Psi_{L,R} = \begin{pmatrix} Q_{L,R}^1 \\ Q_{L,R}^2 \\ Q_{L,R}^3 \\ L_{L,R} \end{pmatrix}$$

Leptons as the fourth “color”

- ✓ $SU(4)$ is the smallest group containing the $U_1 \sim (3, 1, 2/3)$
- ✓ No proton decay (protected by symmetry)
- ✗ The (flavor universal) Pati-Salam model cannot work
 - The bounds from $K_L \rightarrow \mu e$ lift the LQ mass to 100 TeV
- ✗ The associated Z' would be excessively produced at LHC
 - $M_U \sim M_{Z'} \sim \mathcal{O}(\text{TeV})$ & $\mathcal{O}(g_s)$ Z' couplings to valence quarks

4321 model(s)

[Georgi and Y. Nakai, [1606.05865](#); Diaz, Schmaltz, Zhong, [1706.05033](#); Di Luzio, Greljo, Nardecchia, [1708.08450](#)]

We can “protect” the light families by decorrelating $SU(4)$ from the SM color group

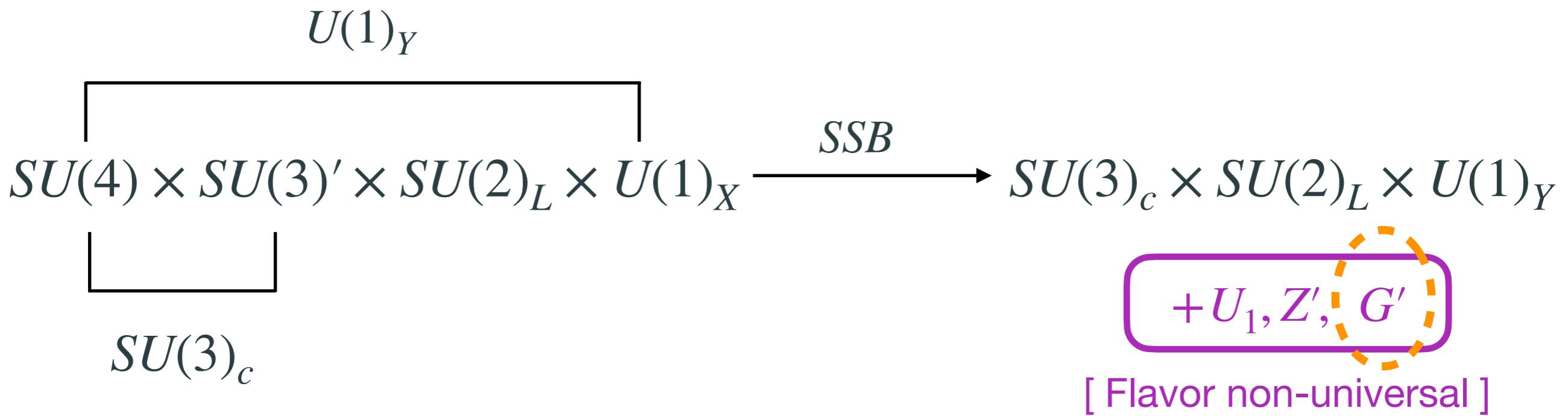
PS group:

$$\mathcal{G}_{\text{PS}} \supset \mathbf{SU}(4) \times \mathbf{SU}(2)_L \times \mathbf{U}(1)_R$$

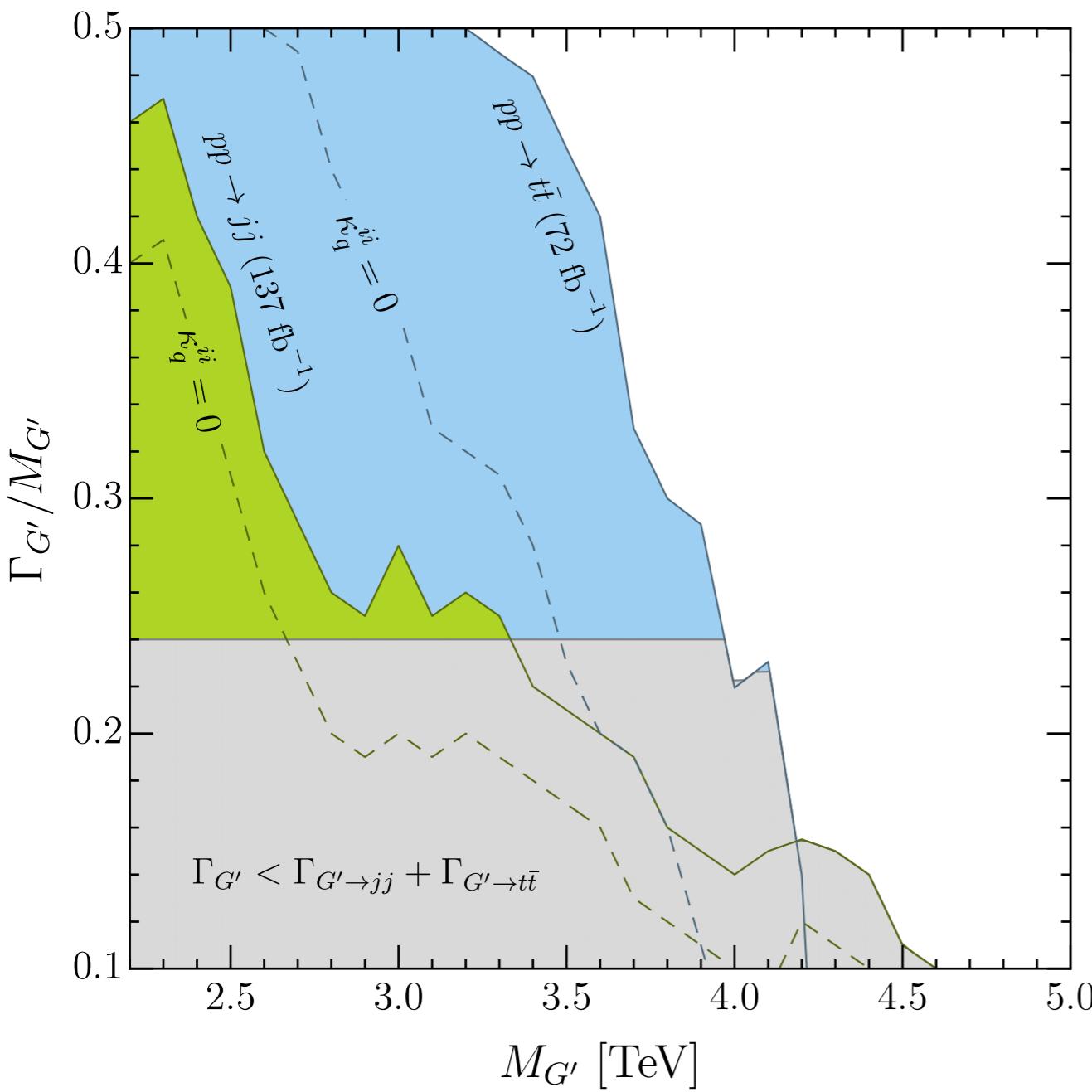
[Flavor universal]



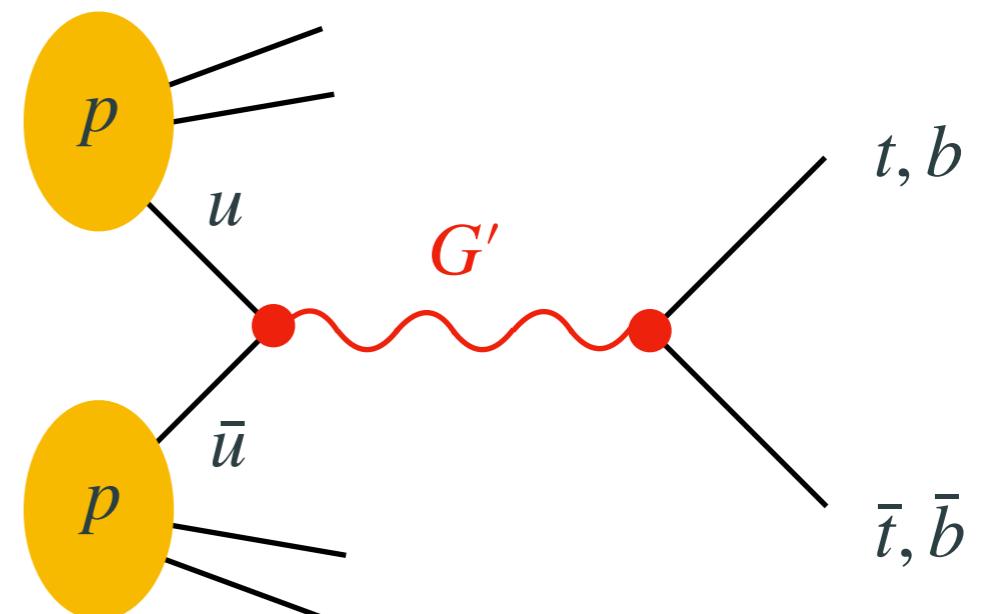
4321 group: $\mathcal{G}_{4321} \equiv \mathbf{SU}(4) \times \mathbf{SU}(3)' \times \mathbf{SU}(2)_L \times \mathbf{U}(1)_X$ [Flavor non-universal]



Coloron direct searches at the LHC



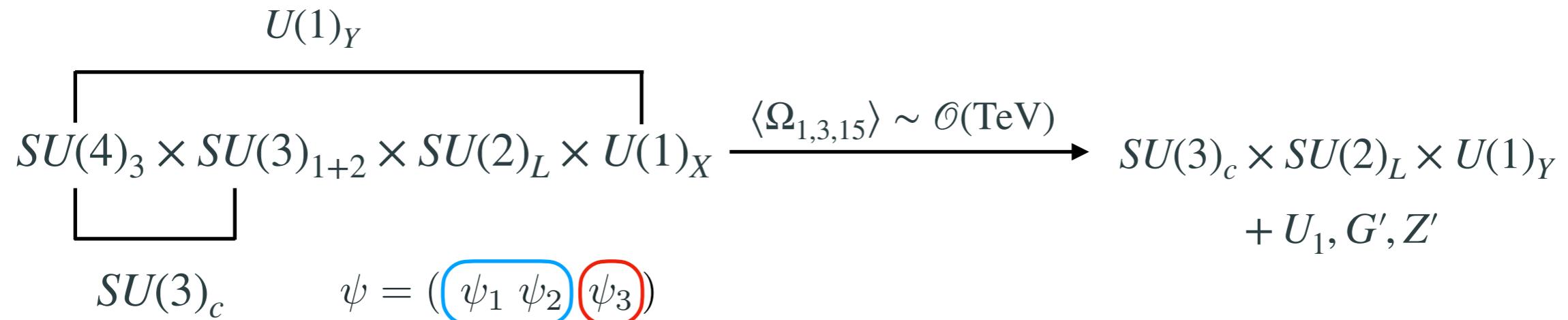
Relevant collider signatures for G'
("coloron" = heavy color-octet vector)



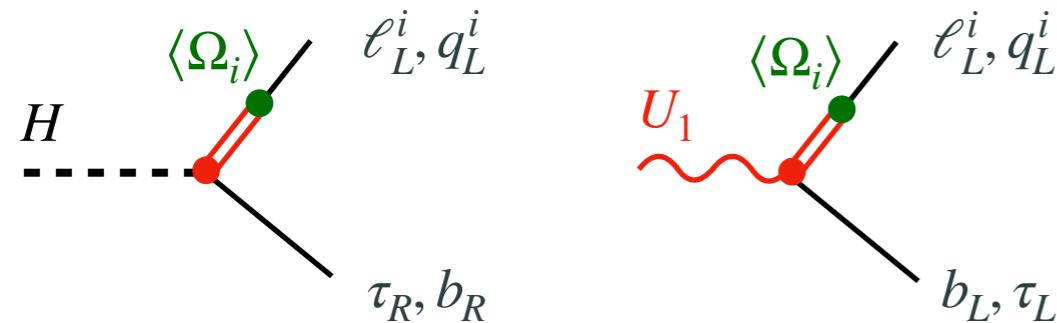
Strongest constraint on the model
scale from $pp \rightarrow \bar{t}t$

[Cornella et al., 2103.16558]

Third-family quark-lepton unification at the TeV scale



- ★ Direct new physics couplings to 3rd family only
- ★ CKM mixing and new physics couplings to lighter families via (small) mixing with χ fermions



Field	$SU(4)$	$SU(3)'$	$SU(2)_L$	$U(1)_X$
q_L^i	1	3	2	1/6
u_R^i	1	3	1	2/3
d_R^i	1	3	1	-1/3
ℓ_L^i	1	1	2	-1/2
e_R^i	1	1	1	-1
ψ_L	4	1	2	0
ψ_R^\pm	4	1	1	$\pm 1/2$
χ_L^i	4	1	2	0
χ_R^i	4	1	2	0
H	1	1	2	1/2
Ω_1	4	1	1	-1/2
Ω_3	4	3	1	1/6
Ω_{15}	15	1	1	0

1st & 2nd families
3rd family
 $n_\chi = 2$

- ★ Fully calculable loop contributions

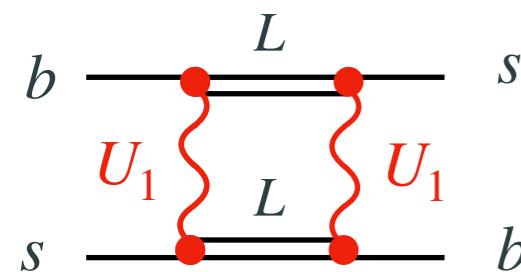
[JFM, Isidori, König, Selimovic,
[1910.13474](#), [2006.16250](#), [2009.11296](#)]

[Bordone, Cornella, JFM, Isidori [1712.01368](#), [1805.09328](#);
Greljo, Stefanek, [1802.04274](#);
Cornella, JFM, Isidori [1903.11517](#)]

Importance of loop effects: U_1 is the new W

Some (important) effects appear only at one loop

I. bound on M_L from $\Delta F = 2$
(similar to SM with charm quark)



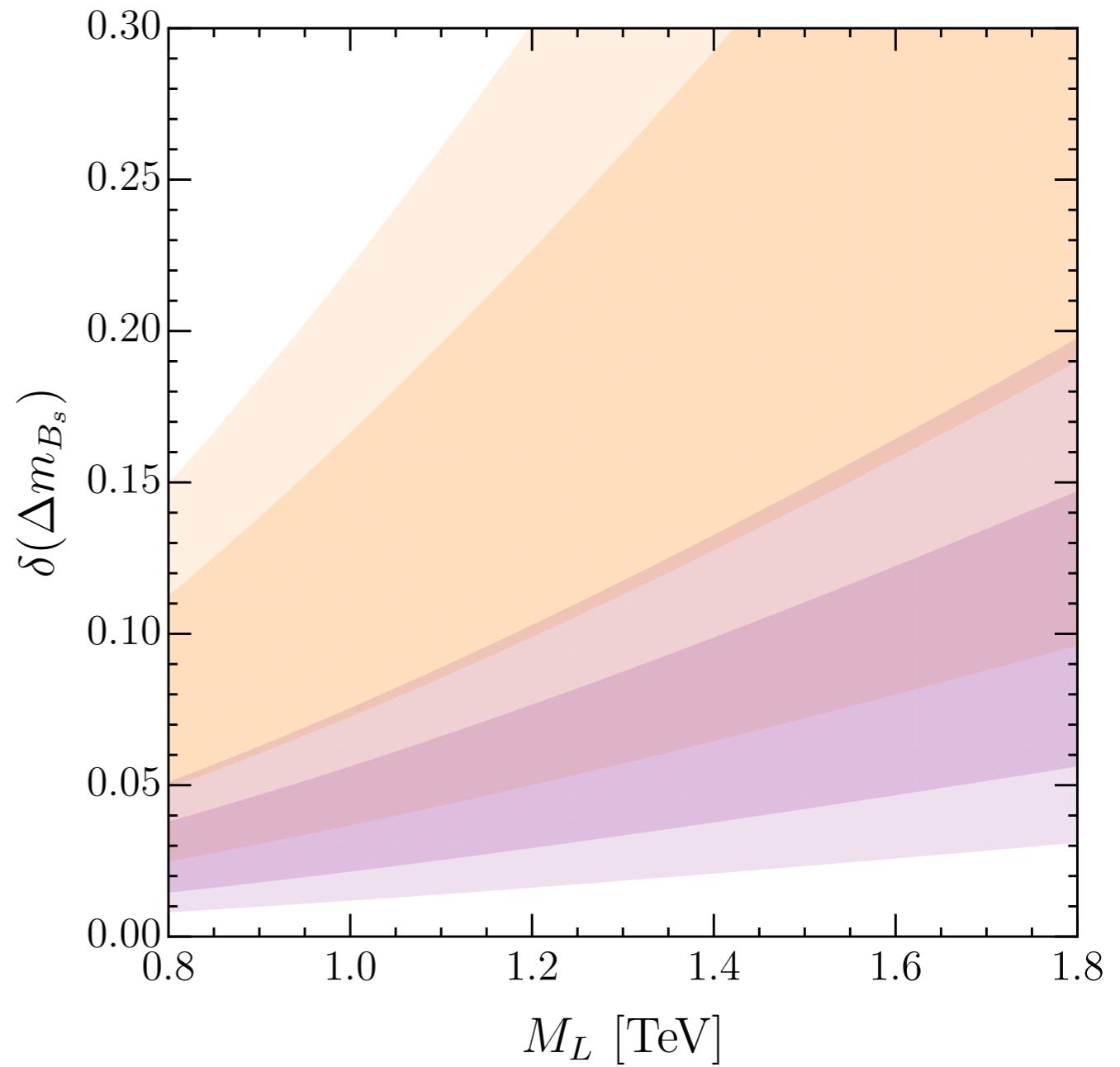
$$\chi = \binom{Q}{L}$$

$$\sim \Delta R_{D^*}^2 M_L^2 \implies M_L \sim \text{TeV}$$

[Within the LHC reach!]

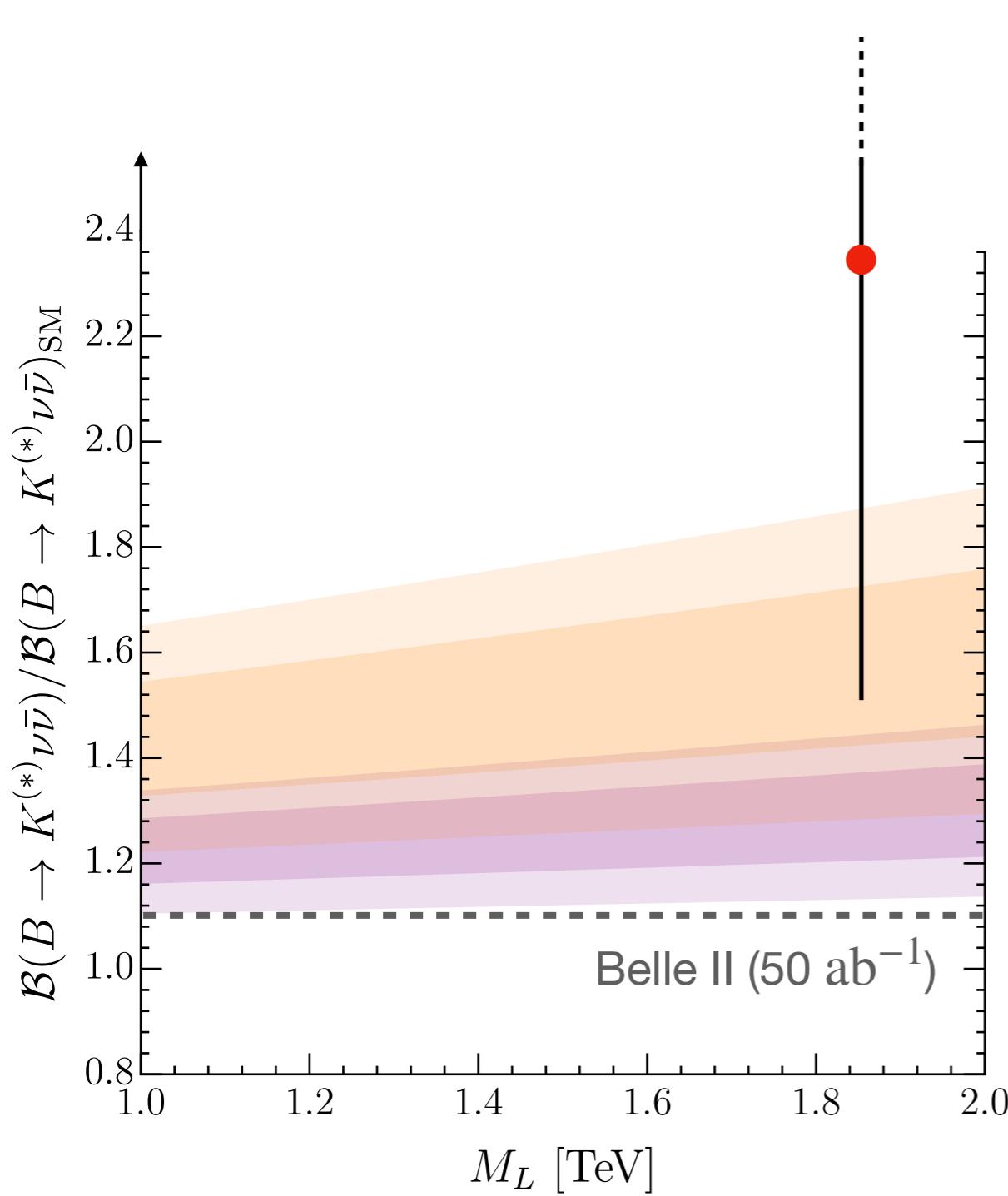
GIM-like suppression of FCNC loops

[di Luzio, JFM, Greljo, Nardecchia, Renner [1808.00942](#);
Cornella, JFM, Isidori [1903.11517](#);
JFM, Isidori, König, Selimovic, [2009.11296](#)]

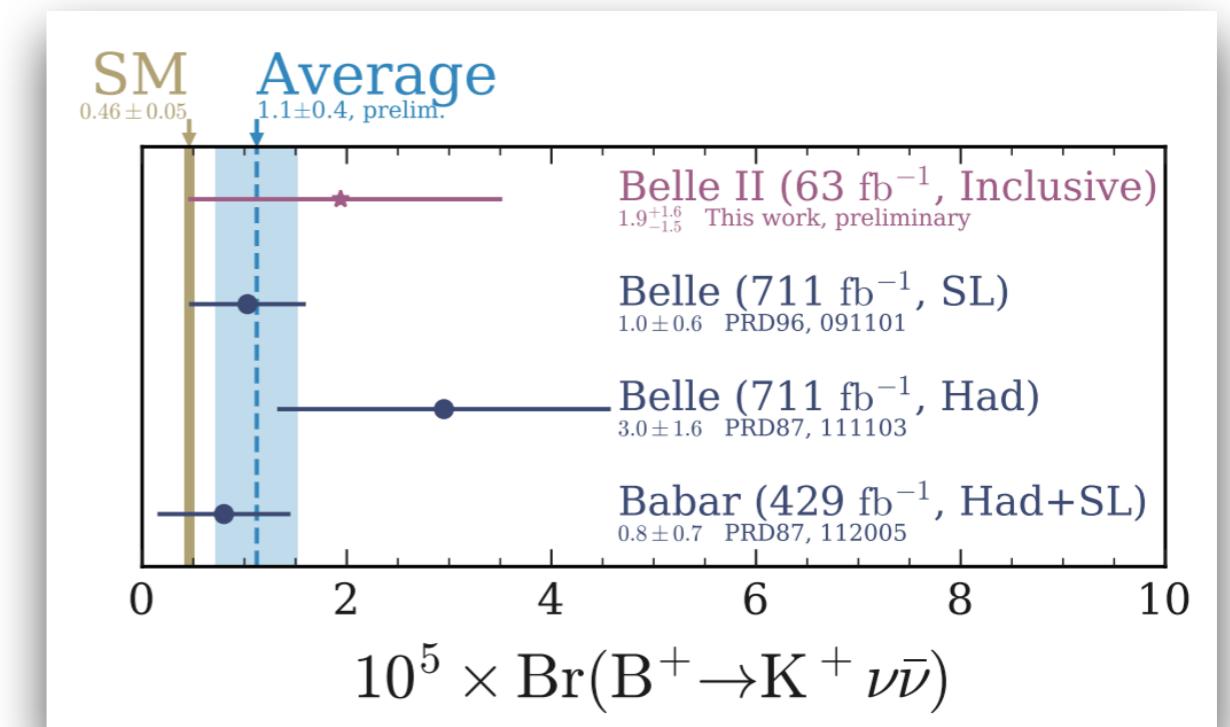
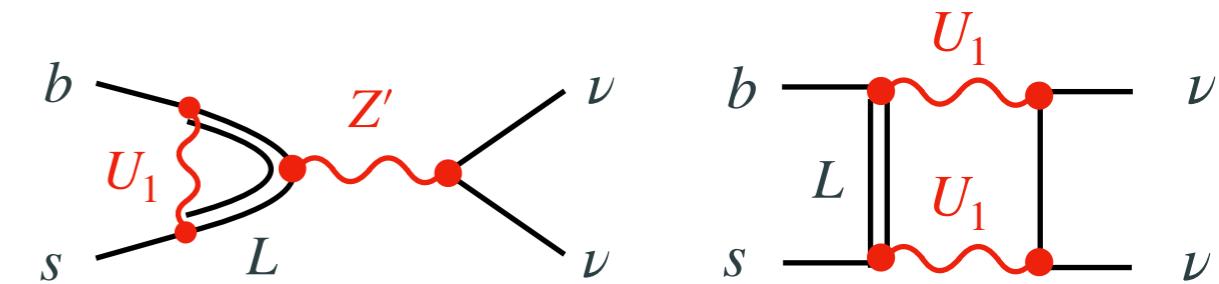


Importance of loop effects: U_1 is the new W

Some (important) effects appear only at one loop



II. $B \rightarrow K^{(*)} \nu \bar{\nu}$ [JFM, Isidori, König, Selimovic, [2009.11296](#)]



Filippo Dattola (Belle II) @ Moriond EW

Back to Pati-Salam... but in five dimensions!

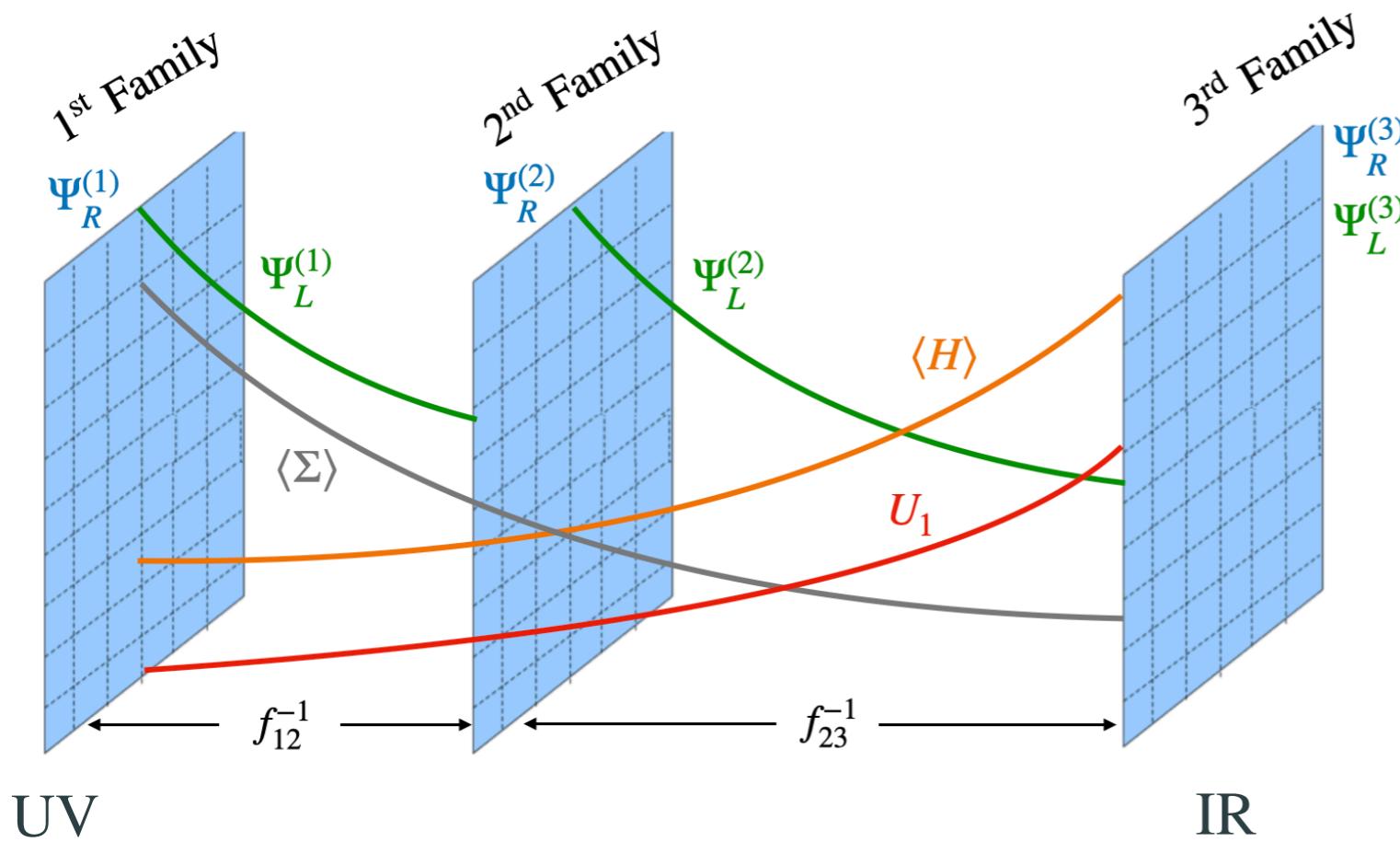
Warped compact extra dimension (AdS_5) with multiple four-dimensional branes

$$ds^2 = e^{-2ky} \eta_{\mu\nu} dx^\mu dx^\nu - dy^2$$

$$\text{PS}_{5D} \equiv SU(4) \times SU(2)_L \times SU(2)_R$$

Flavor \longleftrightarrow fermion (quasi-)localization
in each of the branes
(fermion mixing from nearest neighbor
interactions)

[Dvali, Shifman, '00]



[JFM, Isidori, Pagès, Stefanek, [2012.10492](#)

JFM, Isidori, Lizana, Selimovic, Stefanek, w.i.p]

$$f_L^i(y) \sim e^{-c_i y}$$

Higgs IR (quasi-)localization \longleftrightarrow
stabilization of the mass hierarchies
(as in the Randall-Sundrum model)

[Randall, Sundrum '99]

$SU(4)$ symmetry broken in the UV
 $\longleftrightarrow U_1$ LQ as $SU(4)$ KK mode(s)
(dominantly coupled to 3rd family)

Conclusions

The statistical significance of the **LFU anomalies in $b \rightarrow s\ell\ell$** data is growing and amounts to 4.6σ when considering only observables with a very robust theory prediction [Recent evidence of LFUV in R_K alone (3.1σ)]

LFU anomalies in $b \rightarrow c\tau\bar{\nu}$ data are also seen. Although statistically less significant, in combination with $b \rightarrow s\ell\ell$ data, they point to TeV-scale new physics with a similar flavor structure to that of the SM Yukawas

→ Possible connection to the **origin of SM flavor hierarchies** from a multi-scale picture

New physics picture consistent with low- & high-energy data, but **new physics effects should emerge soon in other observables** (closing in on the mass gap?)

[**Caveat:** Many of the predictions rely on the present value of $R_{D^{(*)}}$]

Plenty of upcoming measurements from both the energy and intensity frontiers [e.g. part of LHC run II data still to be analyzed and Belle II data coming]



Thank you!